

Track projection results and scattered electron reconstruction/kinematics

Barak Schmookler

Outline

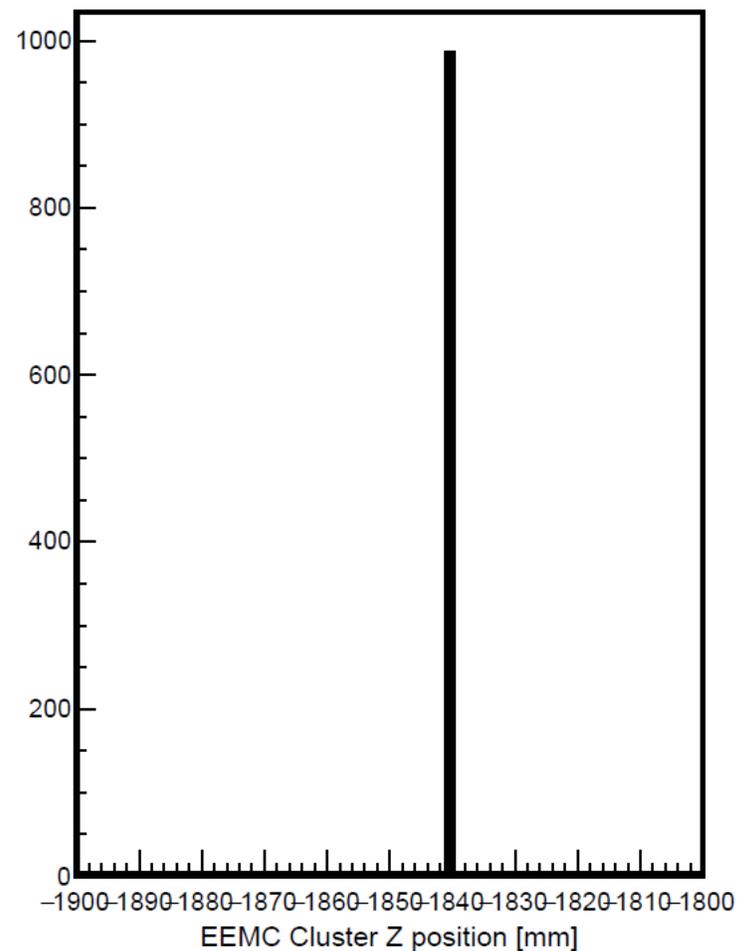
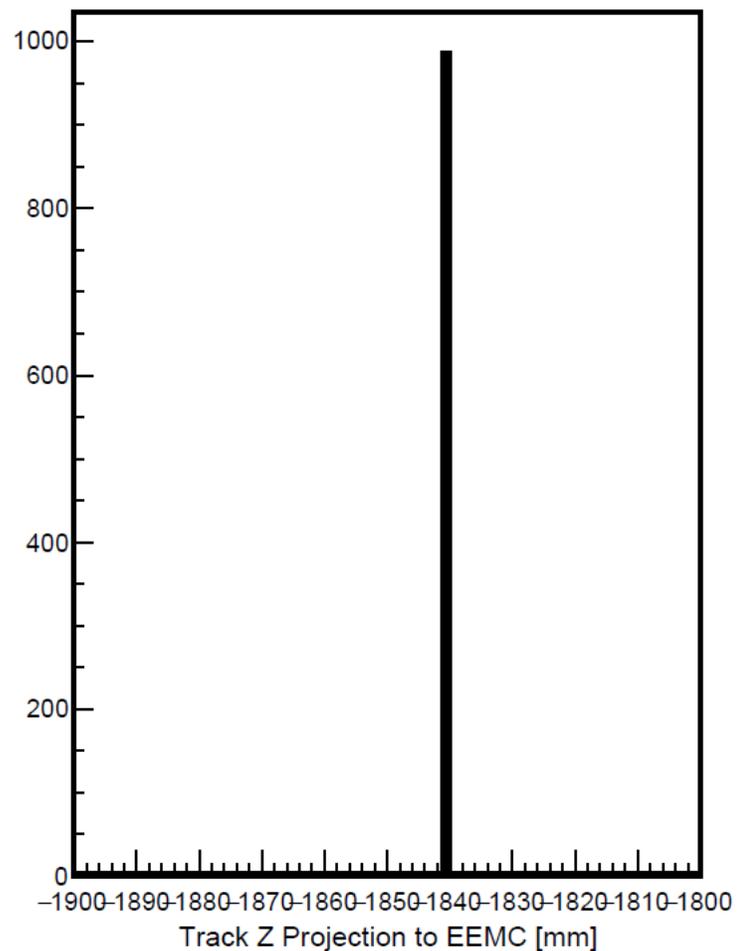
- ❑ Track projection implementation and testing in *Juggler*
- ❑ Comparison with S3 calorimeter results
- ❑ Next steps
- ❑ Scattered electron kinematics

Track projection (propagation) implementation

- We need to project the reconstructed tracks to other sub-detectors.
- A standalone code using *Juggler* output the [ACTS::Propagator](#) class was written by Wenqing Fan.
- This was then implemented into a new *Juggler* [class](#) by Barak Schmookler. Results shown on the next pages use this class.
- The class has now been ported to an *ElCrecon* [class](#) by Dmitry Romanov, and is being generalized for projections to other detectors. This will hopefully go into the next simulation campaign.
- Additional work to associate projection to track (trajectory) in output ROOT is ongoing.

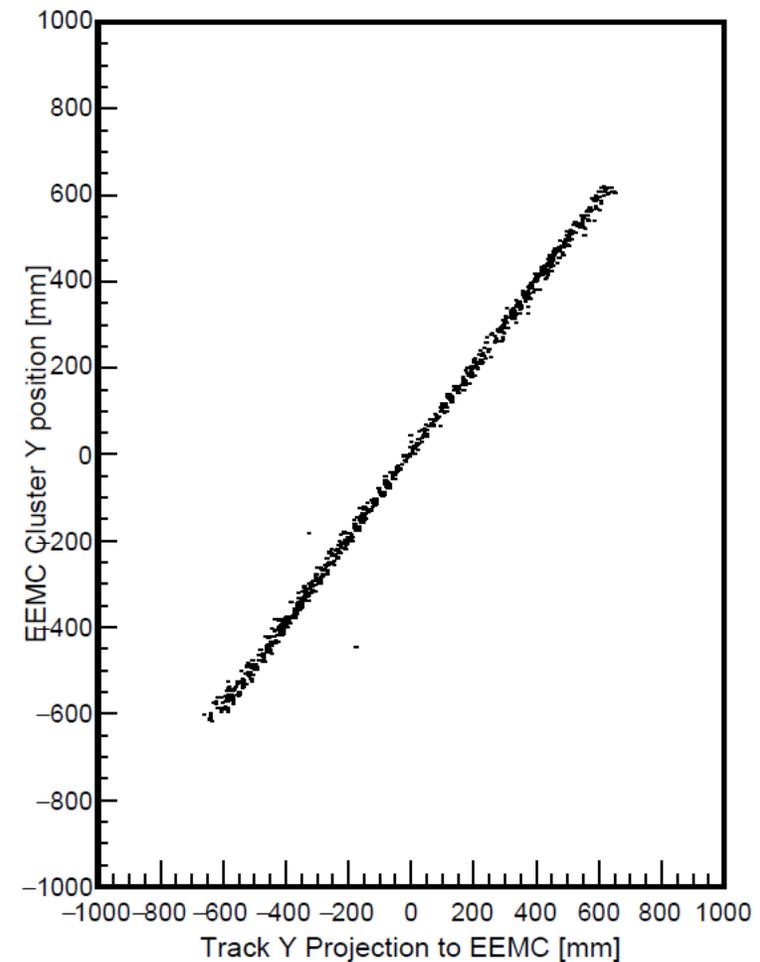
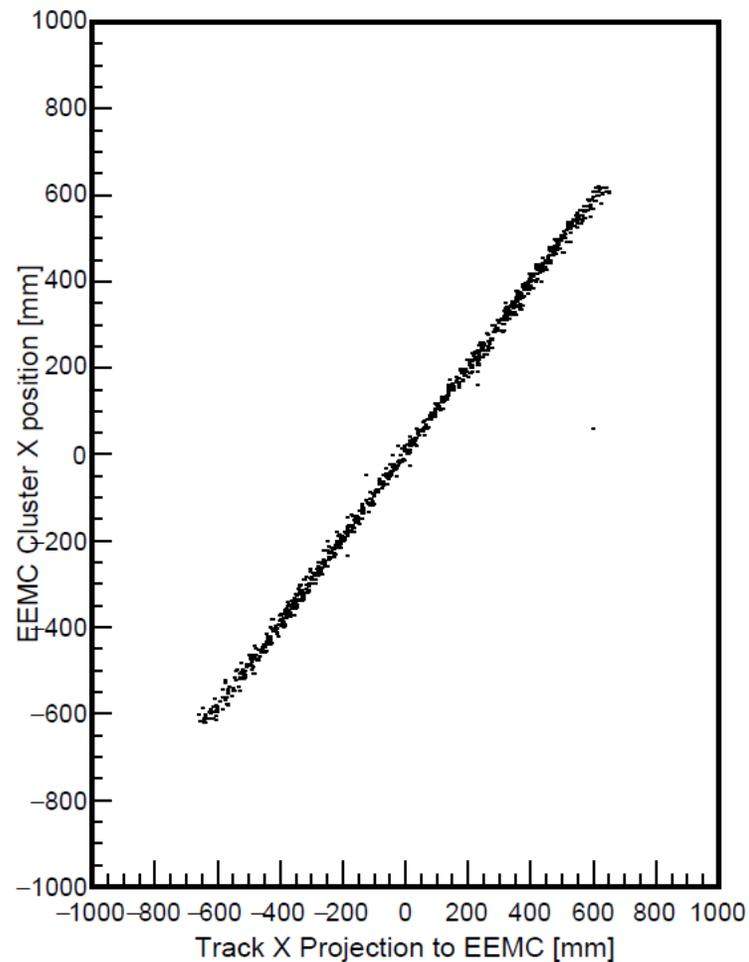
Track Projection test: EEMC

Single Electrons generated:
 $1 \text{ GeV} < E < 20 \text{ GeV}$
 $160^\circ < \theta < 170^\circ, 0^\circ < \phi < 360^\circ$



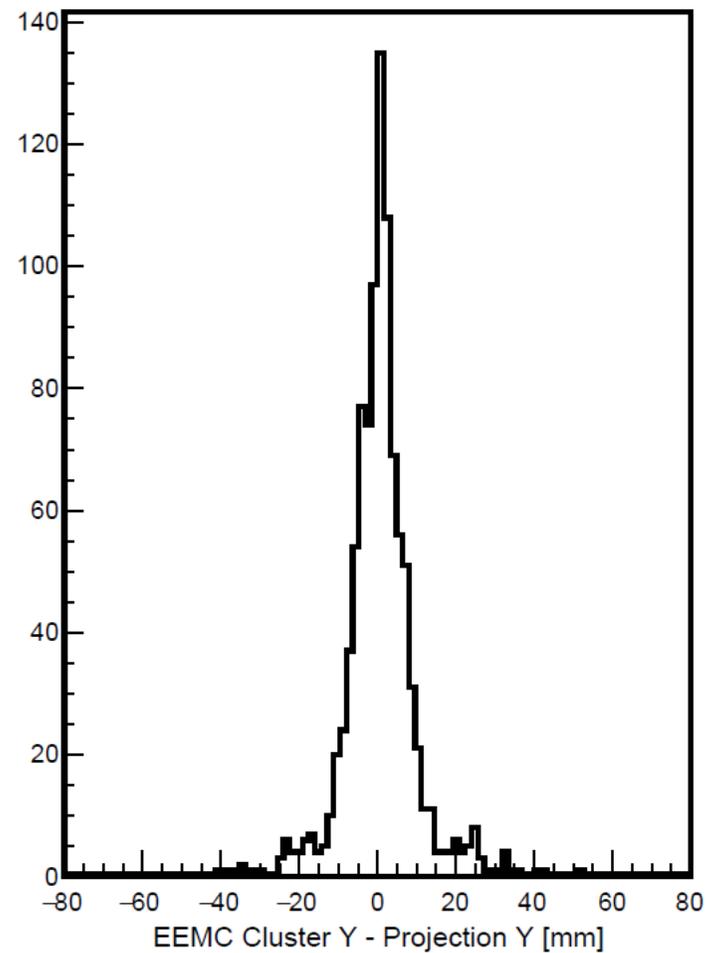
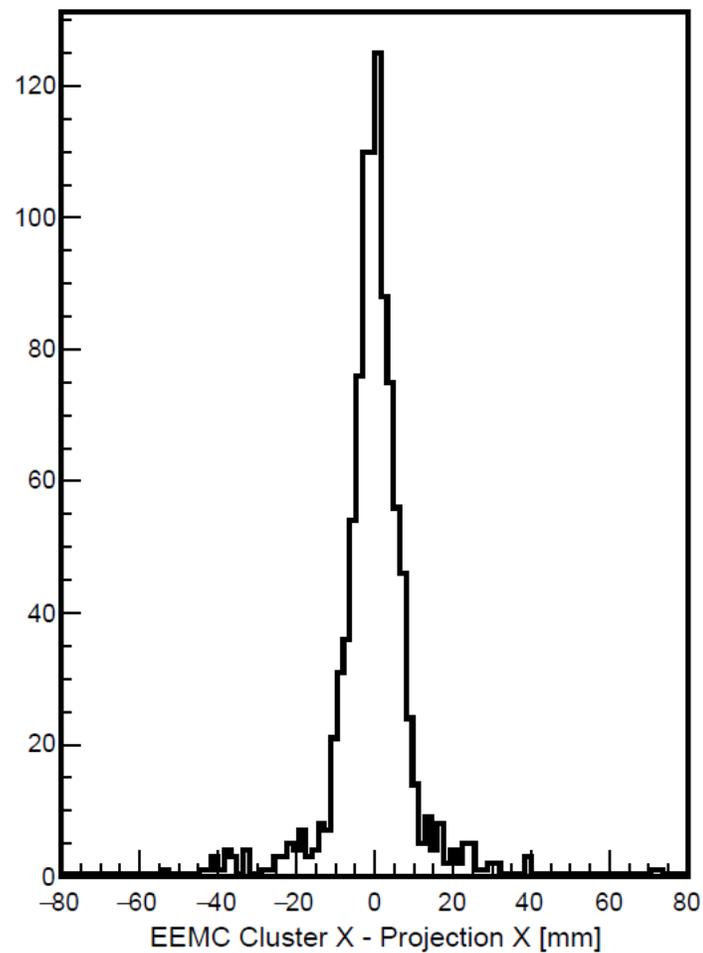
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Track Projection test: EEMC

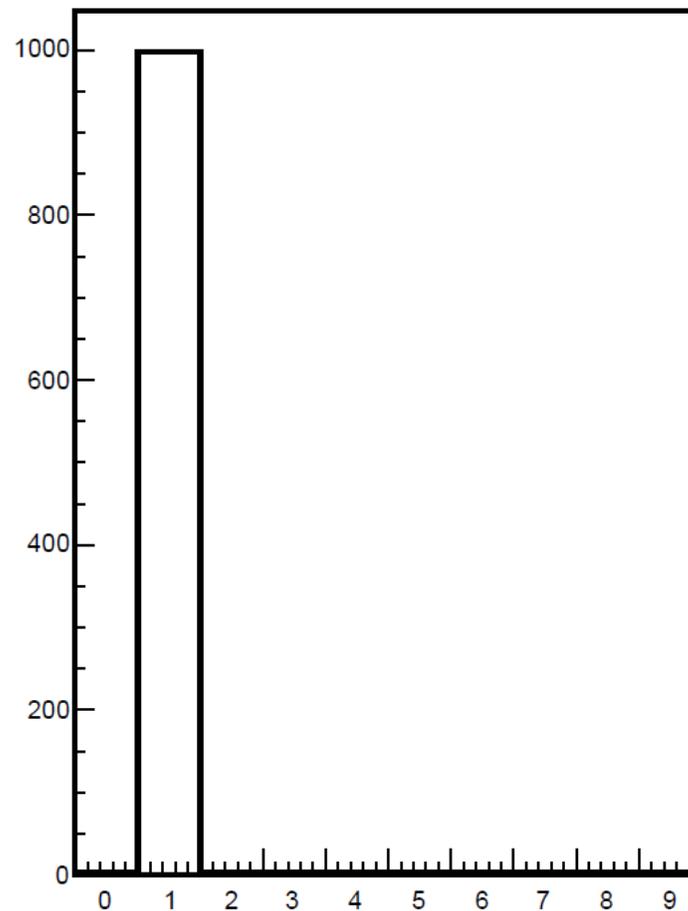
Single Electrons generated:
 $1 \text{ GeV} < E < 20 \text{ GeV}$
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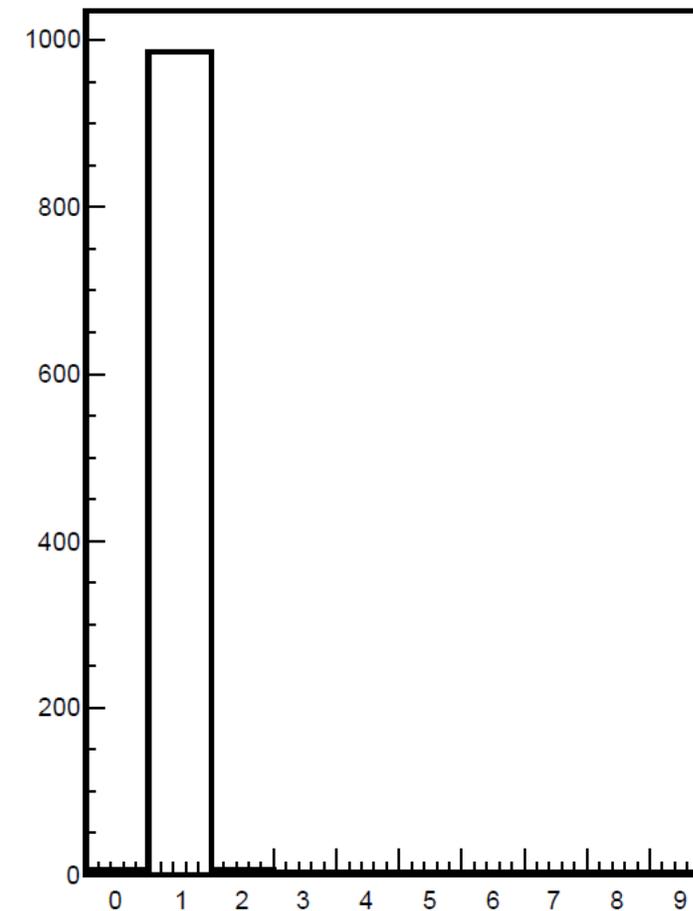
Track Projection test: EEMC

Single Electrons generated:
1 GeV < E < 20 GeV
160° < θ < 170°, 0° < ϕ < 360°

Number of reconstructed tracks

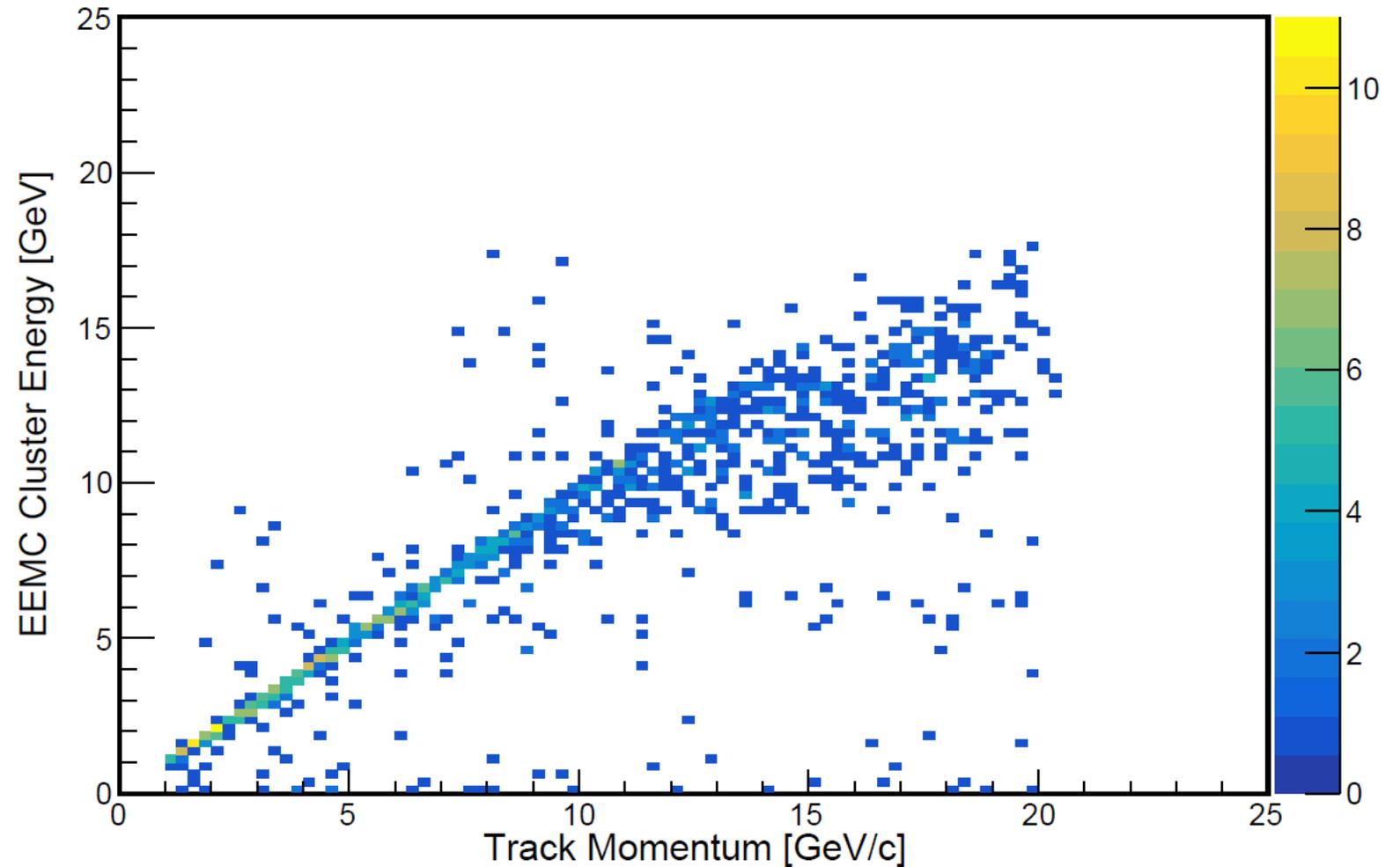


Number of EEMC clusters



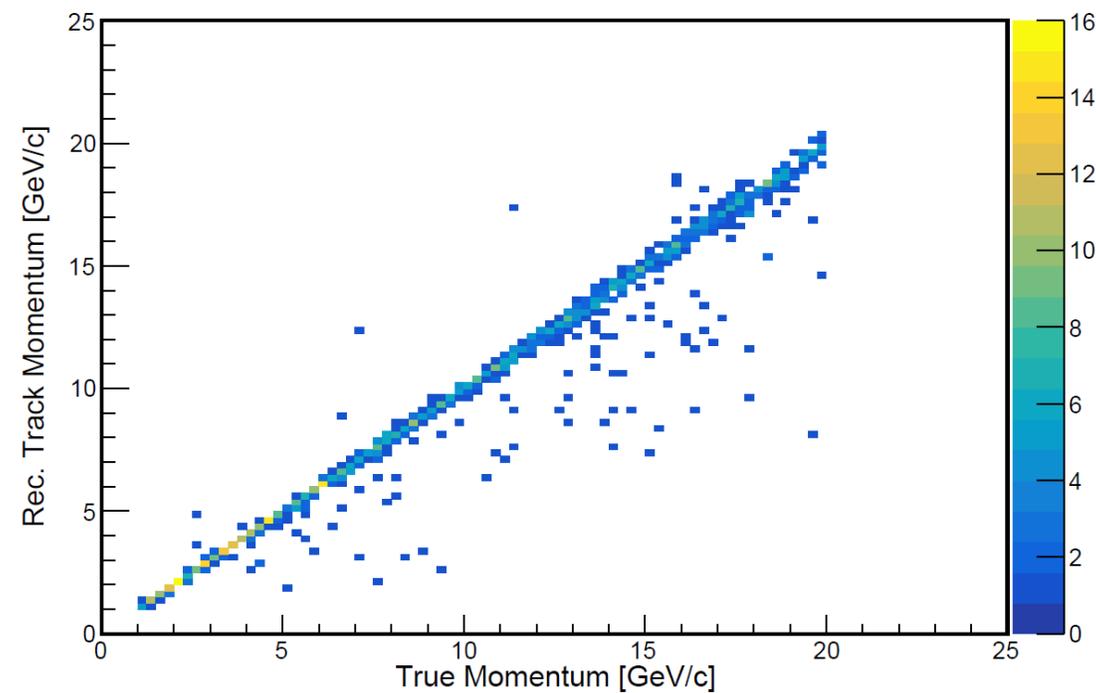
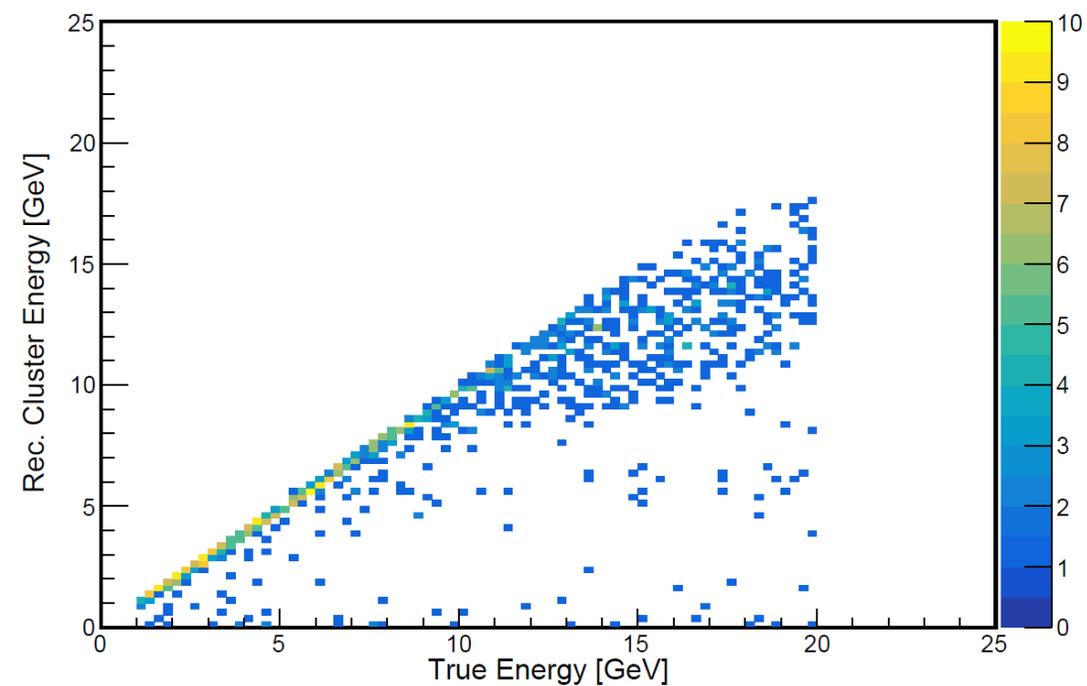
Track Projection test: EEMC

Single Electrons generated:
 $1 \text{ GeV} < E < 20 \text{ GeV}$
 $160^\circ < \theta < 170^\circ, 0^\circ < \phi < 360^\circ$



Track Projection test: EEMC

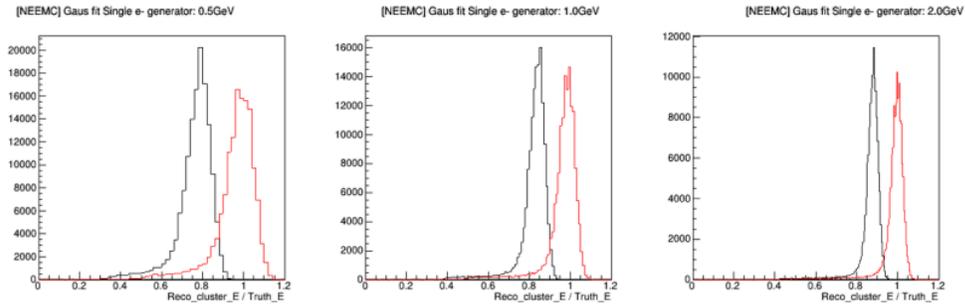
Single Electrons generated:
 $1 \text{ GeV} < E < 20 \text{ GeV}$
 $160^\circ < \theta < 170^\circ, 0^\circ < \phi < 360^\circ$



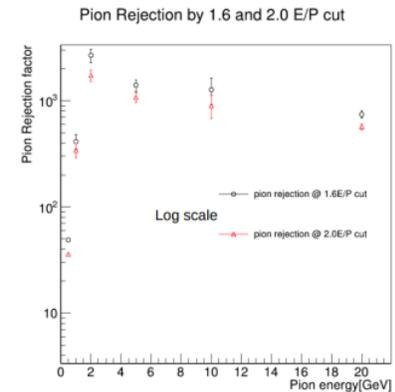
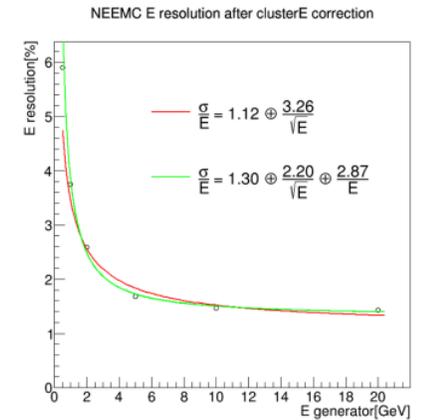
Comments

Slide from Calorimeter group

Backward ECal



- First look at single particle simulations
- Energy resolution and pion rejection values as expected
- No issue identified so far



- The track projection itself looks very good!
- I see poor energy reconstruction and resolution for the EEMC for >10 GeV electrons if I use both the ‘nightly’ and the latest tagged geometry. This seems to disagree with the results from the calorimeter group.
- I used *Juggler* for all the above plots – using the calorimeter reconstruction benchmarks ‘options’ file.
- Need to test with *EICRecon* to see if effect persists.

Next steps

1. Repeat the above studies using the *EICRecon* framework. This will take some effort, since we will need to write a 'factory' to save the track projections to our output ROOT files. Dmitry Romanov said he will help me do this next week.
2. After getting the *EICRecon* results, work with the calorimeter group to understand the resolution discrepancy.
3. Write a simple algorithm for scattered electron reconstruction (electron finder) based on track-to-cluster matching and an E/p cut. Inclusion of PID detectors can come later.

Scattered Electron kinematics

- For the beam energies that will be used at the EIC and considering scattered electrons in the pseudo-rapidity range of $-4 < \eta < 4$ (where $Q^2 \gg m_e^2$), we can relate the inclusive kinematic variables to the scattered electron angles and energies as

$$y_e = 1 - \frac{E'_e}{2E_e}(1 - \cos \theta_e) ,$$

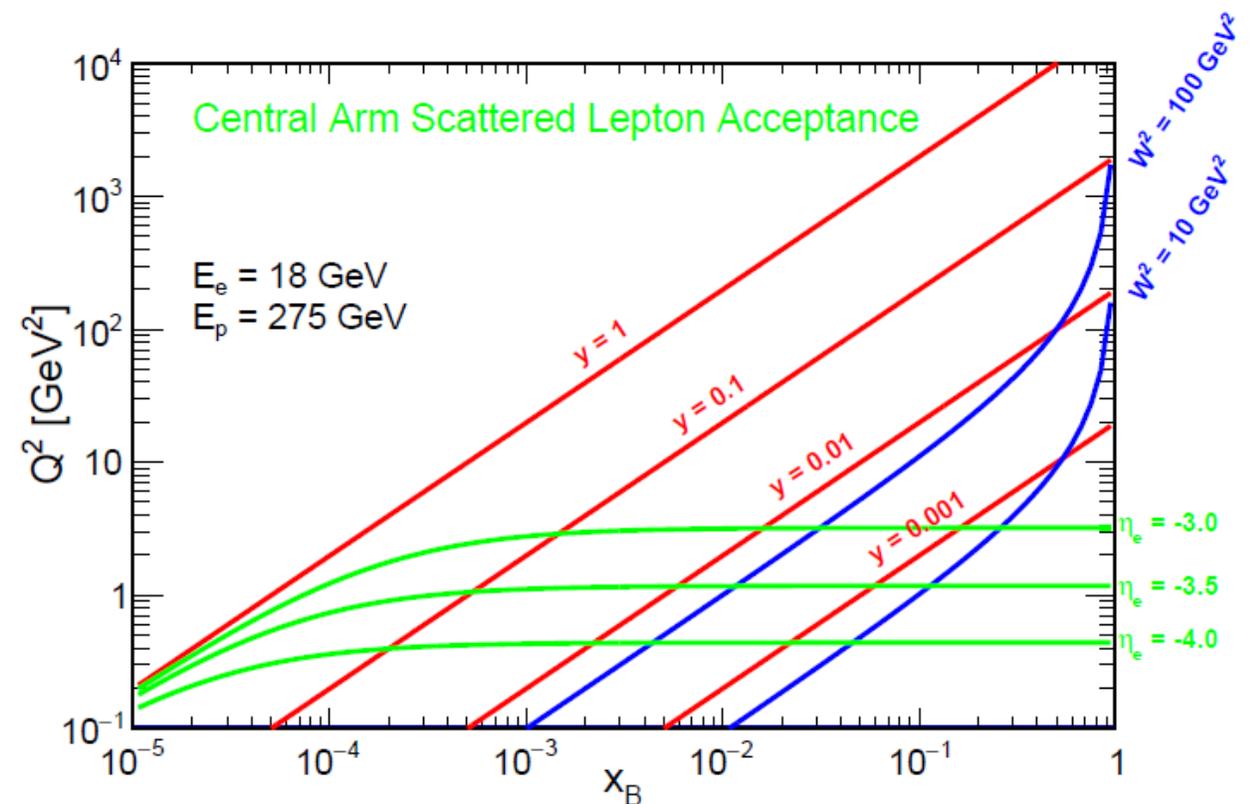
$$Q_e^2 = 4E_e E'_e \cos^2(\theta_e/2) ,$$

$$x_e = \frac{Q_e^2}{s y_e} .$$

- Note how neither Q^2 nor y depend explicitly on the proton beam energy.

Electron Angular Acceptance

- For many EIC physics processes which have a requirement that $Q^2 > 1 \text{ GeV}^2$, an angular acceptance of $\eta \gtrsim -3.6$ will allow full coverage at the highest EIC beam energy setting. At lower energies, this same acceptance coverage would allow access to lower values of Q^2 (see next slide).
- Any processes which require $Q^2 < 1 \text{ GeV}^2$ at the highest energy setting will need an extended acceptance below $\eta \approx -3.6$.
- For inclusive physics, coverage below $Q^2 = 1 \text{ GeV}^2$ has strong motivations: **to study the perturbative to non-perturbative transition; to give access to lowest possible x , which is well-aligned with the central EIC physics aims to study mass generation and dense systems of gluons; and to minimize the 'gap' in Q^2 coverage between the central detector and the far-backwards low- Q^2 tagger.**

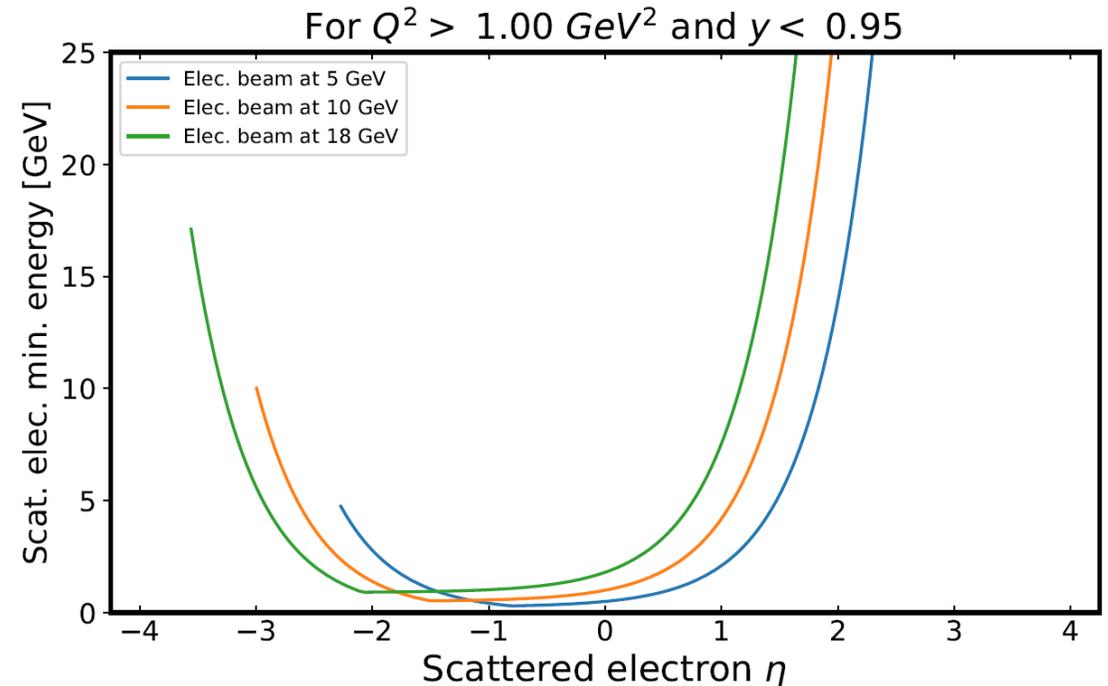


Electron Minimum momentum (energy)

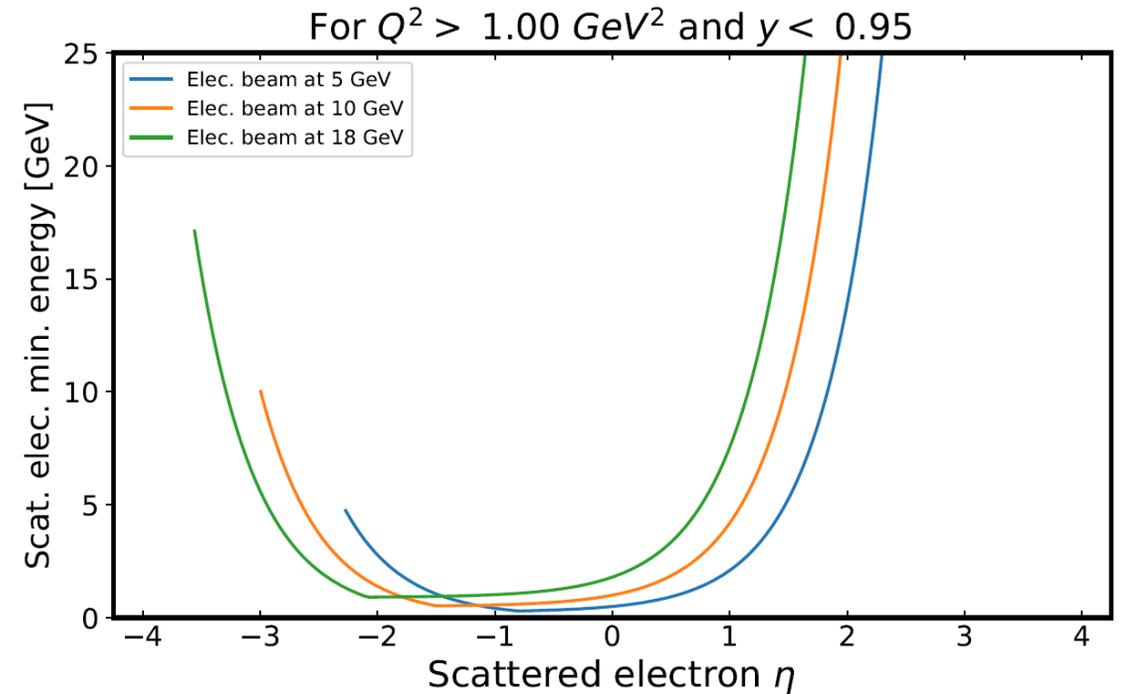
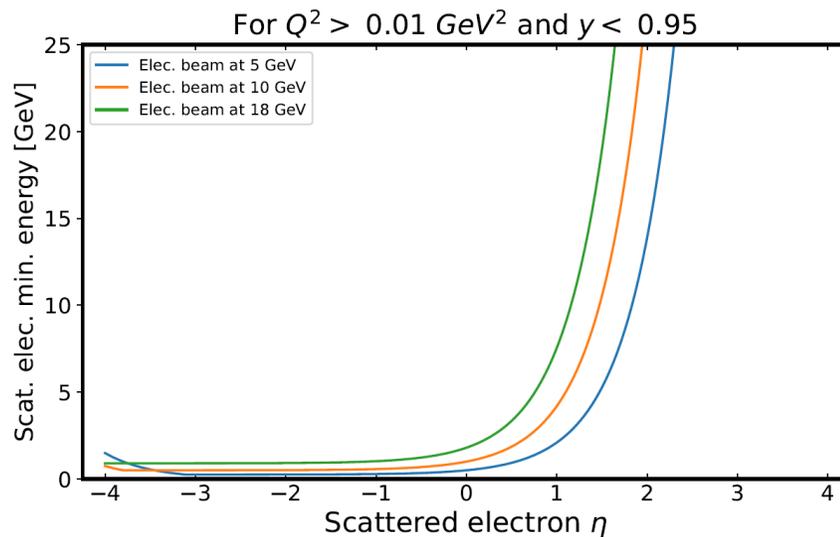
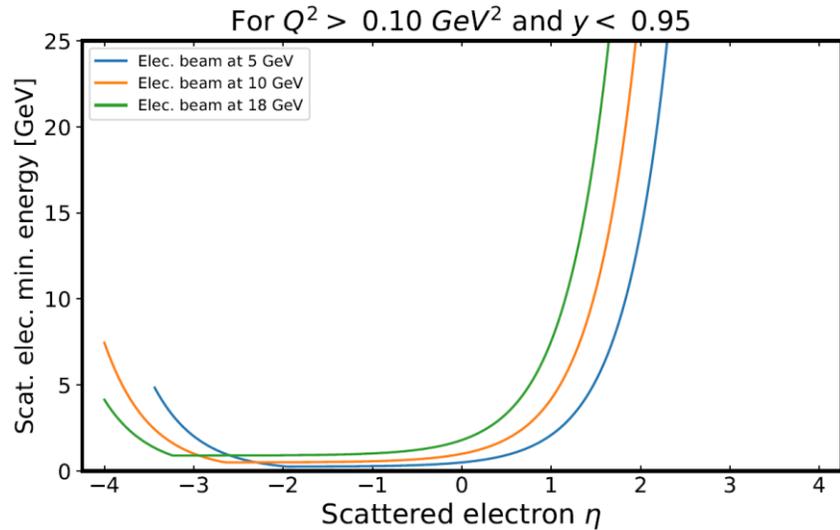
➤ Consider the case where the physics requires $Q^2 > 1 \text{ GeV}^2$. The plot on the right shows as a function of η the minimum electron energy that satisfies both the $Q^2 > 1 \text{ GeV}^2$ requirement and the $y < 0.95$ requirement.

➤ There are a few important features of this plot:

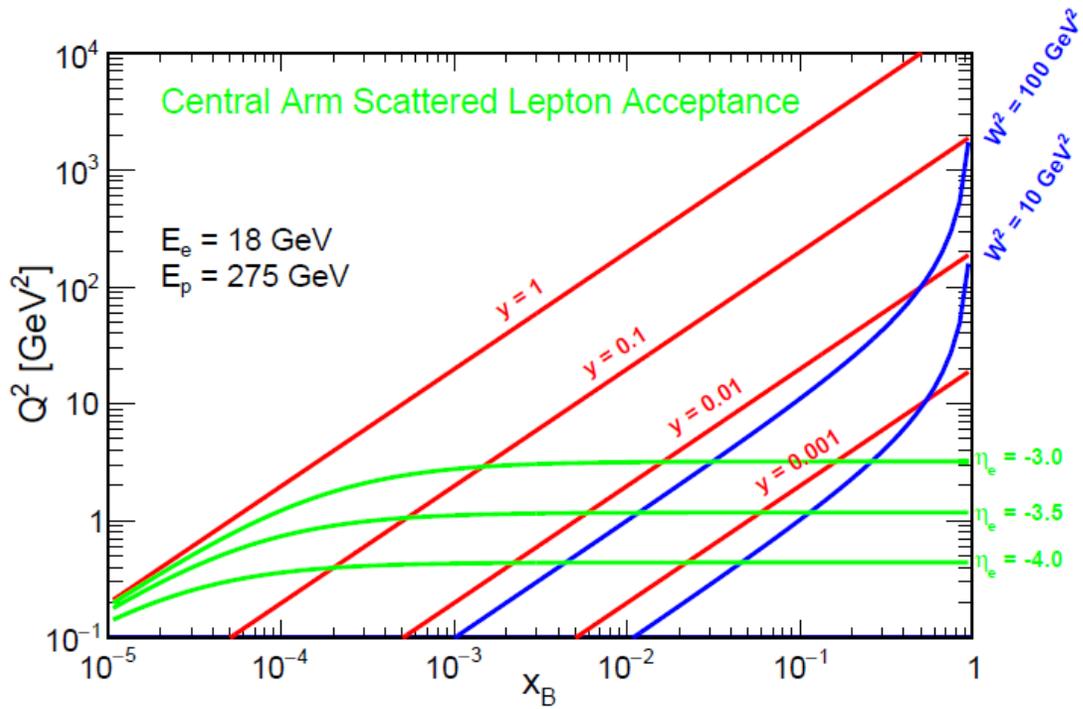
1. The curves do not extend to the lowest possible values of pseudo-rapidities. This is because at the most negative values of η , the scattered electron can not be created at $Q^2 = 1 \text{ GeV}^2$ (y would be negative), only at lower values of Q^2 .
2. Starting at the most negative η value that is allowed, each minimum energy curve decreases towards more positive values of η . For this left part of the curve, the minimum energy is exactly at the $Q^2 = 1 \text{ GeV}^2$ limit (while still satisfying the $y < 0.95$ requirement).
3. Moving towards more positive values of η , each minimum energy curve reaches a global minimum value and then begins to grow. Once the curve begins to increase towards more positive values of η , the minimum energy of the scattered electron is at the $y = 0.95$ limit (while still satisfying the $Q^2 > 1 \text{ GeV}^2$ requirement).



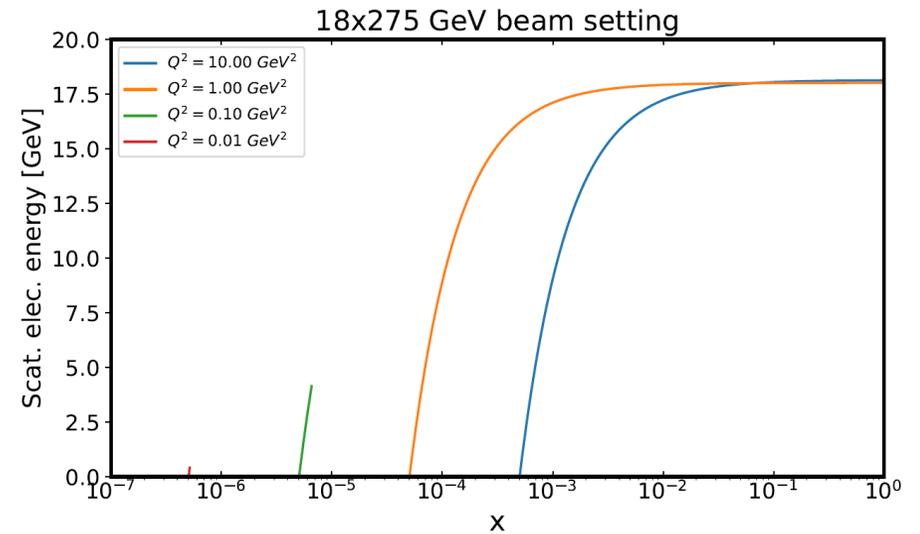
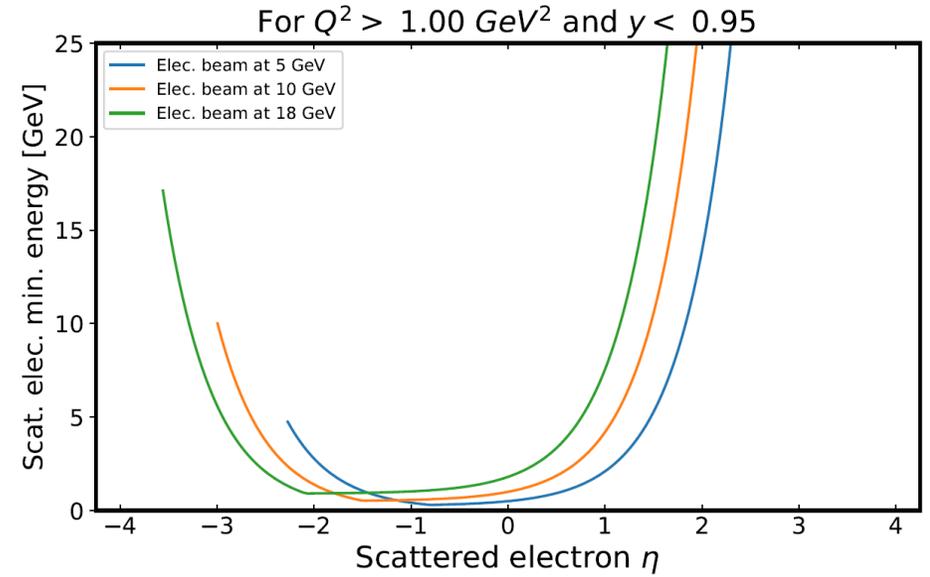
Scattered electron energy in backwards direction is quite large unless at very low Q^2



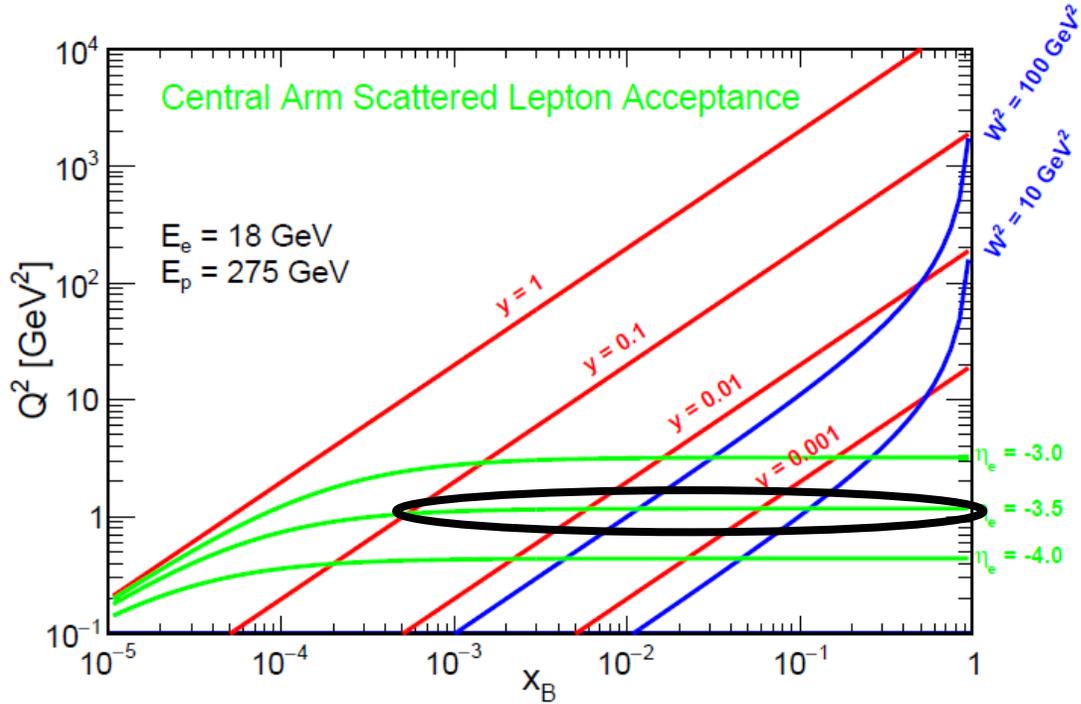
Comments



11/22/22

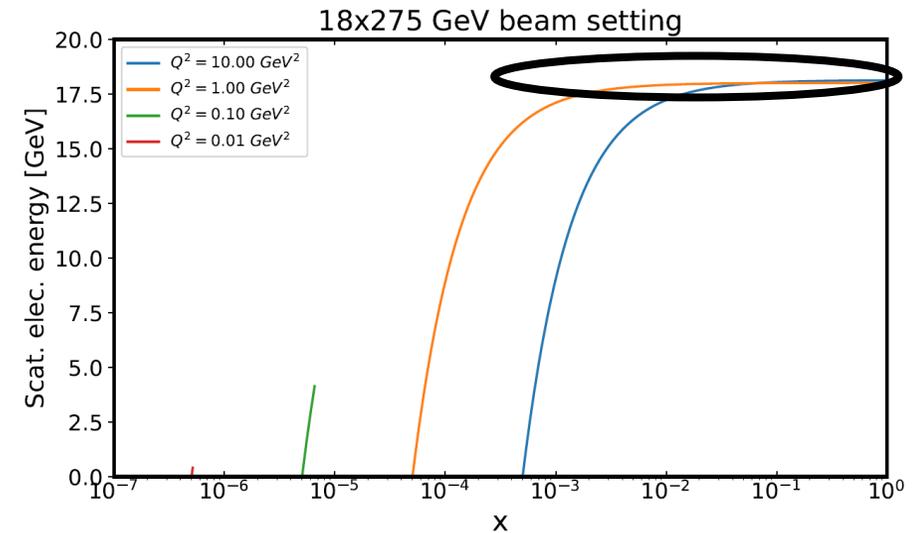
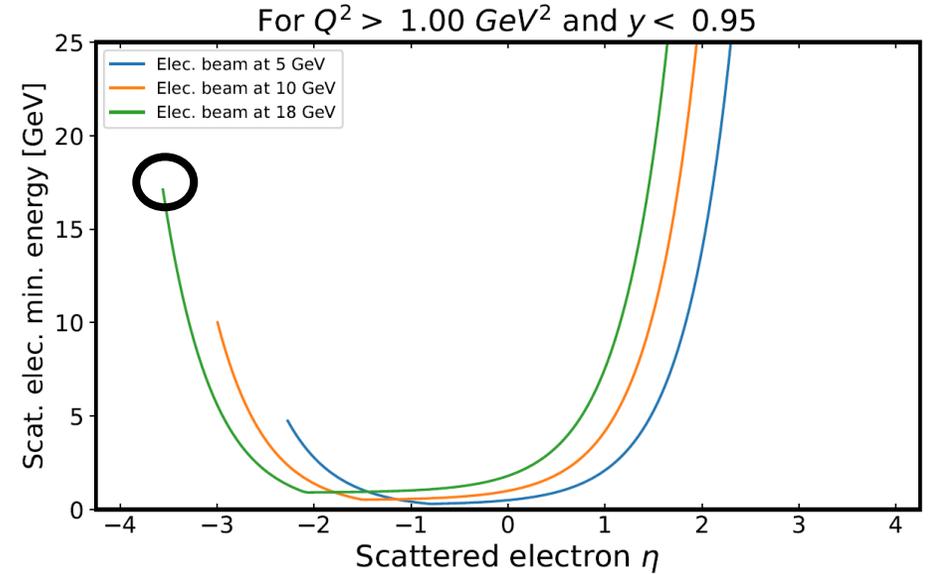


Comments



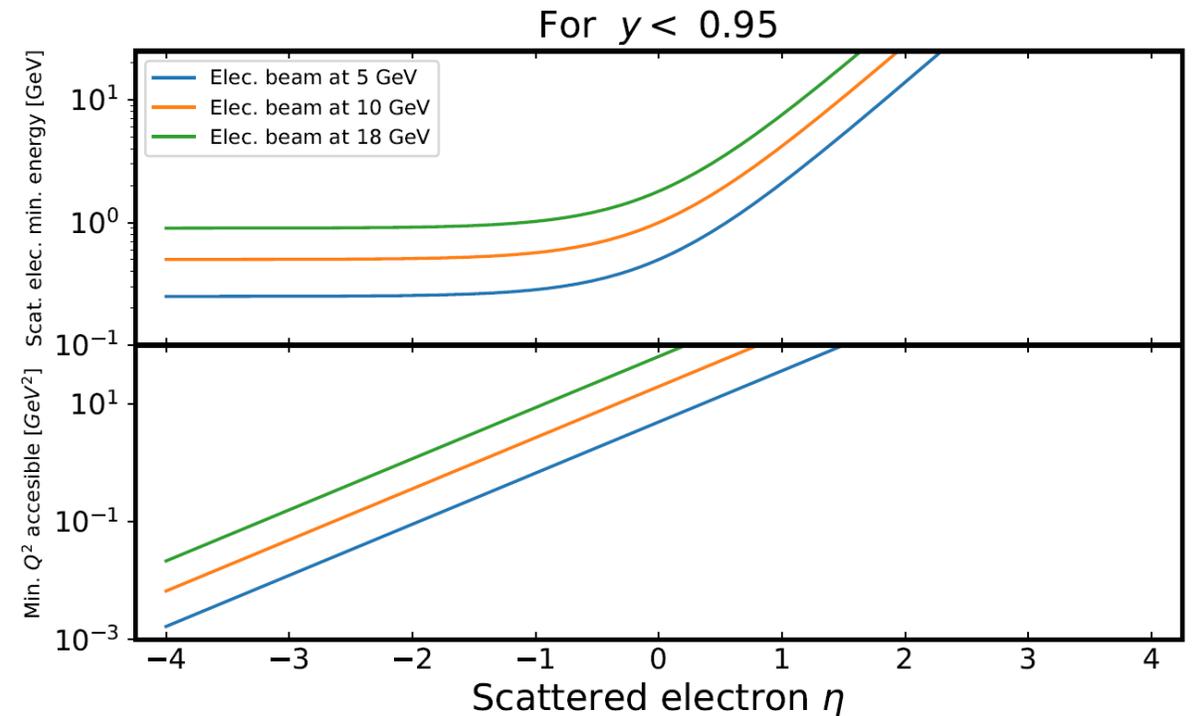
- For $Q^2 = 1 \text{ GeV}^2$, the scattered electron energy is at almost a fixed energy for a large range of x . Reconstruction of x using the scattered electron is impossible here.
- Note how more central rapidities correspond to lower x at a fixed Q^2 .

11/22/22



Comments on backwards tracking

- For the reconstruction of the scattered electron in the backwards part of the negative endcap, the track momentum will always be used for E/p cuts, and the reconstructed angle for Q^2 reconstruction.
- The scattered electron momentum (energy) reconstruction will probably rely on the tracking detector only for lower- Q^2 ($Q^2 \ll 1 \text{ GeV}^2$) events. The question of how low in Q^2 we really need to measure is being studied by the Inclusive PWG now. Questions of pion background and beam divergence effects also need to be considered.



Comments on backwards tracking

$$t \simeq -(\vec{p}_{t,\psi} + \vec{p}_{t,e})^2$$

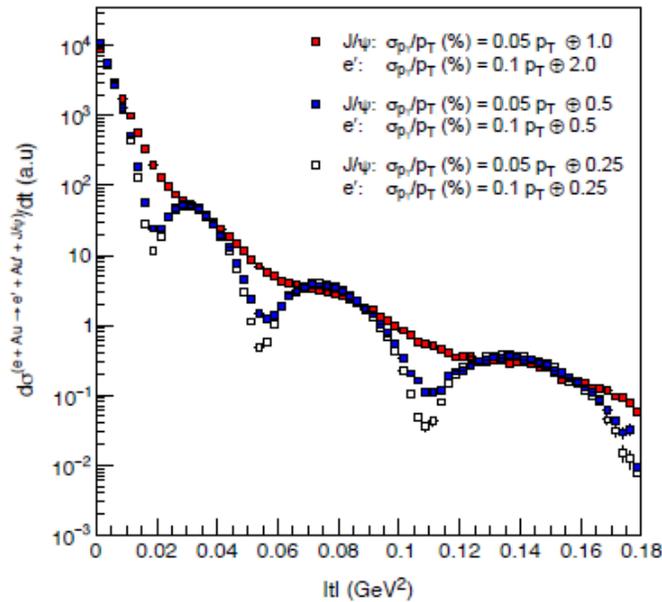


Figure 8.85: Illustration of the impact of different p_T resolutions on the coherent J/ψ production cross section, $d\sigma/dt$, for $1 < Q^2 < 10 \text{ GeV}^2$.

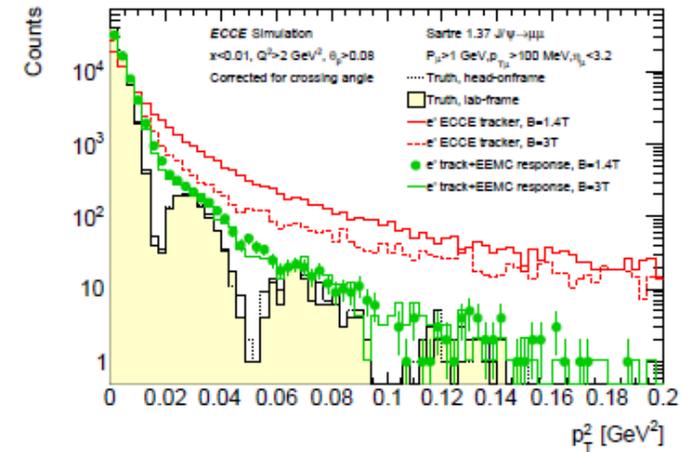


Figure 3.20: Simulation of the t dependence of the elastic diffractive J/ψ production in eA for 1.4 T and 3.0 T field strengths. Results are shown for analyses using information from the tracking (red lines) or tracing + calorimetry (green line and points) systems for the lepton momentum reconstruction. As can be seen, the t reconstruction resolution is dominated by the calorimetry resolution the field strength has small impact on this measurement.

Comments on backwards tracking

- Besides the scattered electron, other particles in the negative endcap need to be reconstructed for various physics processes.
- In addition, measurements of total $E-p_z$ are important for reducing the sizes of photoproduction background and radiative corrections.

