A Pseudo-Dirac Bino and Neutrino Masses at the LHC

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Pilar Coloma, SI, PRL. 117 no.11 111803, arXiv:1606.06372 Julia Gehrlein, Patrick Fox, SI, in preparation

Neutrinos have mass

PDG review, <u>www.nu-fit.org</u>, ...

Mass eigenstates are different than flavor eigenstates

Neutrinos are massive particles



Massless in the SM — no right-handed neutrinos

SM singlet fermions

1- Can have Dirac masses

$$y_{\nu}\bar{\ell}\,\tilde{\phi}\,\nu_R \to m_{\nu}\overline{\nu_L}\nu_R \quad \checkmark \quad y_{\nu} \sim 10^{-12}$$

Fermion masses are technically natural...

'We' don't like very small numbers

2- Or Majorana masses — Seesaw mechanism

$$y_{\nu}\bar{\ell}\,\tilde{\phi}\,\nu_{R} + \frac{1}{2}M_{R}\nu_{R}^{c}\nu_{R} \longrightarrow m_{\nu} \sim \frac{y_{\nu}^{2}v^{2}}{M_{R}}$$
$$y_{\nu} \sim \mathcal{O}(1)$$
$$M_{R} \sim \mathcal{O}(10^{14} \text{ GeV}) \qquad \mathcal{V}R$$

Majorana particles are their own antiparticles

Lepton number is violated

Particle numbers are accidental symmetries of the SM...

3- Or pseudo-Dirac masses — Inverse Seesaw Mechanism

Let's say lepton number is (approximately) conserved SM singlets: N, N' $\downarrow \qquad \downarrow$ $L=1 \quad L=-1$ $Y_N \,\bar{\ell} \,\tilde{\phi} \,N + \epsilon Y_{N'} \,\bar{\ell} \,\tilde{\phi} \,N' + M_D \,\bar{N}N'^c + \mu \,\bar{N}N^c + \mu' \,\bar{N}'N'^c$ lepton number violation $\psi = \begin{pmatrix} N \\ N'^{\dagger} \end{pmatrix}$: pseudo-Dirac fermion

3- Or pseudo-Dirac masses — Inverse Seesaw Mechanism

Light neutrino masses:

$$m_{\nu} \sim \frac{\epsilon Y_N^{\mathrm{T}} Y_{N'} v^2}{M_D} + \mathcal{O}\left(\frac{Y_N^{\mathrm{T}} \mu Y_N v^2}{M_D^2}\right)$$

For $\begin{array}{c} Y_N \sim Y_{N'} \sim 1, \\ M_D \sim {\rm TeV} \end{array}$

need:

$$\epsilon \sim 10^{-12}$$

Why so small?

What we need

(Usually) SM singlet fermions as right-handed neutrinos

High mass scale and/or <u>Small</u> (lepton-number-violating) parameters makes LHC searches very hard (impossible?) Hall, Randall, *Nuc.Phys.B-352.2* 1991 Kribs, Poppitz, Weiner, arXiv: 0712.2039 *Frugiuele, Gregoire, arXiv:1107.4634*

SI, McKeen, Nelson, arXiv: 1407.8193 SI, John March-Russell, arXiv:1604.00009

My favorite model for everything! $U(1)_R$ -symmetric SUSY

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Has Dirac gauginos

Dirac gauginos are awesome - less tuning for heavier stops

Solves SUSY CP and flavor problems, ...

$U(1)_{R-L}$ - symmetric SUSY

SUSY particles are charged under $U(1)_{R-L}$

Superfields	$SU_c(3)$	$SU_L(2)$	$U_Y(1)$	$U(1)_R$	$U(1)_{R-L}$
L_i	1	2	-1/2	1	0
E_i^c	1	1	1	1	2
H_u	1	2	1/2	0	0
$W^{lpha}_{ ilde{B}}$	1	1	0	1	1
$\Phi_S = \phi_S^D + \theta S$	1	1	0	0	0
$W'_{\alpha} = \theta D$	1	1	0	1	1

Spurion D-term

SUSY is broken in a hidden sector

Dirac bino mass

No Majorana gaugino masses due to the R-charges

Dirac masses come from the spurion D-term

$$\int d^2\theta \, c \frac{W'_{\alpha}}{\Lambda_M} W^{\alpha}_{\tilde{B}} \Phi_S \rightarrow \underbrace{\frac{cD}{\Lambda_M}}_{\text{Dirac bino mass: } M_D} \tilde{B}S$$

$$\tilde{B} \equiv (1, 1, 0)_{+1} \qquad \longrightarrow \Psi = \begin{pmatrix} \tilde{B} \\ S^{\dagger} \end{pmatrix} \quad : \text{Dirac bino}$$
$$S \equiv (1, 1, 0)_{-1} \qquad \longrightarrow \Psi = \begin{pmatrix} S^{\dagger} \\ S^{\dagger} \end{pmatrix}$$

$U(1)_{R-L}$ must be broken

... because (anomaly mediation)

(Small) Majorana mass for the bino $m_{\tilde{B}} = \frac{\beta(g_Y)}{g_Y} m_{3/2}$

 $m_{3/2}$: gravitino mass

Can also have a singlino Majorana mass

 $m_{\tilde{B}} \sim m_S \ll M_D$ $W(1)_{R-L}$ is only approximately broken $\Psi = \begin{pmatrix} \tilde{B} \\ S^{\dagger} \end{pmatrix}$: pseudo-Dirac bino

All the mass terms



This is an Inverse SeeSaw scenario!

 $U(1)_{R-L} \text{ violation } \propto m_{3/2} \longrightarrow \qquad \text{important for} \\ \text{small neutrino masses} \\ \Psi = \begin{pmatrix} \tilde{B} \\ S^{\dagger} \end{pmatrix} \quad \text{:We call this "bivo"} \quad \text{(pronounced exactly like 'bino')} \\ \text{(like 'too' and 'two')} \end{cases}$

Neutrino masses



 $\rho\simeq 0.7\,$ from mass splittings

No dependence on M_D

Parameter space



No constraints from neutrinoless double-beta decay

LHC Phenomenology of Bivo

Julia Gehrlein, Patrick Fox, SI, in preparation

LHC pheno: production

RH neutrinos are usually produced via mixing with the SM neutrinos



Pay a mixing price on top of EW $\theta^2 \sim 10^{-5}$ interactions:

There are better ways to produce a bino!

LHC pheno: production

Assumptions for the phenomenology: Lightest neutralino is a pure bino Gravitino is the LSP - a few keV Bino is the NLSP Degenerate squarks

LHC pheno: production

Can produce bino via squark decays:



Gluon fusion is the main production channel at 14 TeV LHC

LHC pheno: decays

1) Decays to a gravitino and a photon



2-4) Decays to $W\ell, Z\nu, h\nu$



total width:

 $\Gamma_{tot} \sim \frac{M_D^3}{\Lambda_M^2} \sim O(\text{MeV})$

LHC pheno: signals

Bino decays promptly, no displaced vertices A combination of jets, leptons and missing energy is expected



LHC pheno: signals



Huge thanks to Angelo Monteux for the analysis codes!





Smoking gun signal



electron : muon : tau ratios are fully determined by the neutrino sector: 1:2:1

