

Bridging Nuclear and Quark Matter: Quantum van der Waals Approach to Quarkyonic Transition

Roman Poberezhnyuk

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In collaborations with: Horst Stoecker, Volodymyr Vovchenko, and Tripp Moss

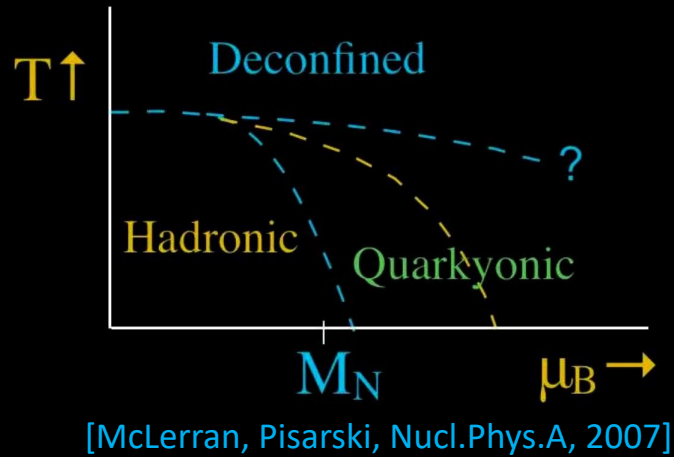


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- Quarkyonic matter (quasiparticle, quark-hadron duality, baryquark)
- Theory of baryon-baryon interactions
- Results
- Applications to neutron star phenomenology

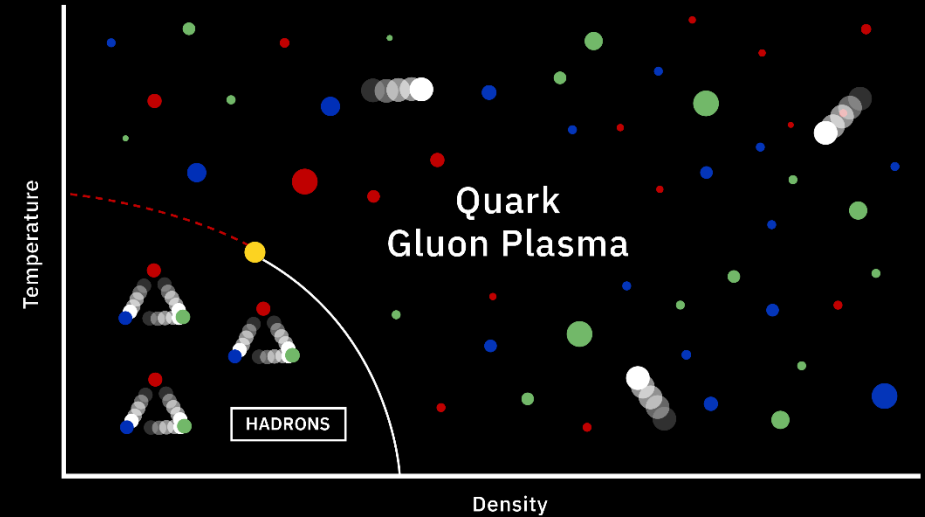
Quarkyonic phase



- Is still confined (color neutrality, quarks are not free) although very dense.
- Density of quarks exceed energy scale of confining interactions ($n_q \gg \Lambda_{QCD}^3$)
- (Confined) quarks become natural degrees of freedom.
- Pauli exclusion principle between individual quarks starts to play a significant role.
- Inspired by limit of large N_c of QCD where confinement persists.

Possible quarkyonic phase would have properties different from both hadronic matter and QGP

Deconfined phase (Quark matter)

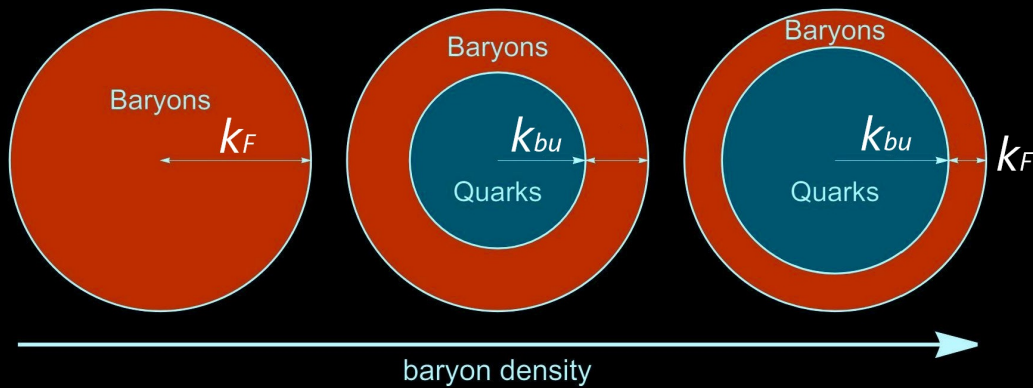


- Degrees of freedom are quarks, Screening of color charges, Polyakov loop expectation value in non-zero

The density of the onset of deconfinement is unknown

Postulates “free” quarks and triplets of “confined” quarks (nucleons) which Pauli exclude each other. Nucleons consist of N_c quarks, $m_q = m_N/N_c$, $N_c=3$.

Shell structure in momentum space



$$n_N^{\text{id}}(k_F) = \frac{g}{2\pi^2} \int_{k_{\text{bu}}}^{k_F} k^2 dk,$$

$$\varepsilon_N^{\text{id}}(k_F) = \frac{g}{2\pi^2} \int_{k_{\text{bu}}}^{k_F} k^2 \varepsilon_N(k) dk,$$

$$n_Q^{\text{id}}(q_F) = \frac{g}{2\pi^2} \int_0^{k_{\text{bu}}/N_c} q^2 dq,$$

$$\varepsilon_Q^{\text{id}}(q_F) = N_c \frac{g}{2\pi^2} \int_0^{k_{\text{bu}}/N_c} q^2 \varepsilon_Q(q) dq.$$

$$g = g_Q = g_N = 4$$

$f_q = n_q/n_B$ at each n_B is found by minimizing the energy density.

Infrared regulator limits the abruptness of quark onset:

$$\rho_Q(q) \rightarrow \rho_Q(q) \frac{\sqrt{q^2 + \Lambda^2}}{q}$$

[McLerran, Pisarsky, Nucl.Phys.A, 2007]

[McLerran, Reddy, PRL, 2019]

[Kie Sang Jeong, McLerran, Sen, PRC, 2020]

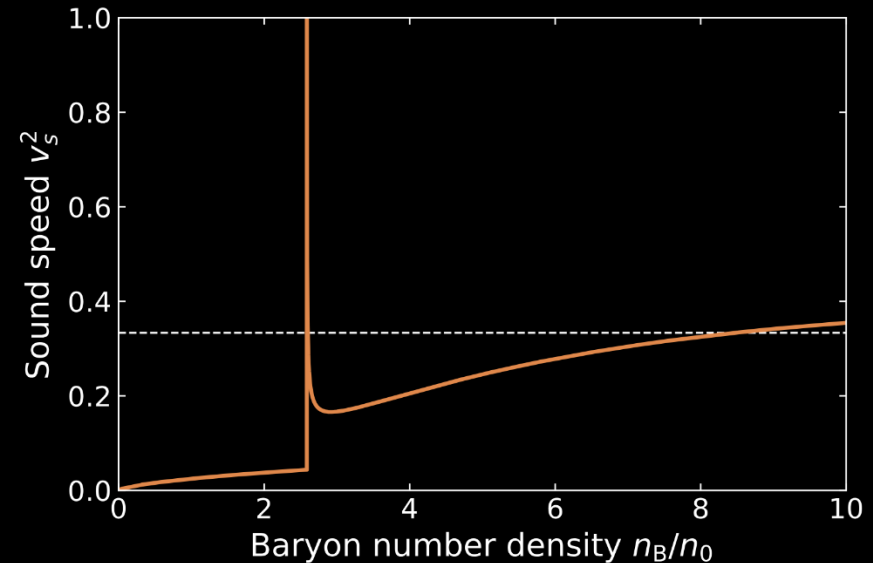
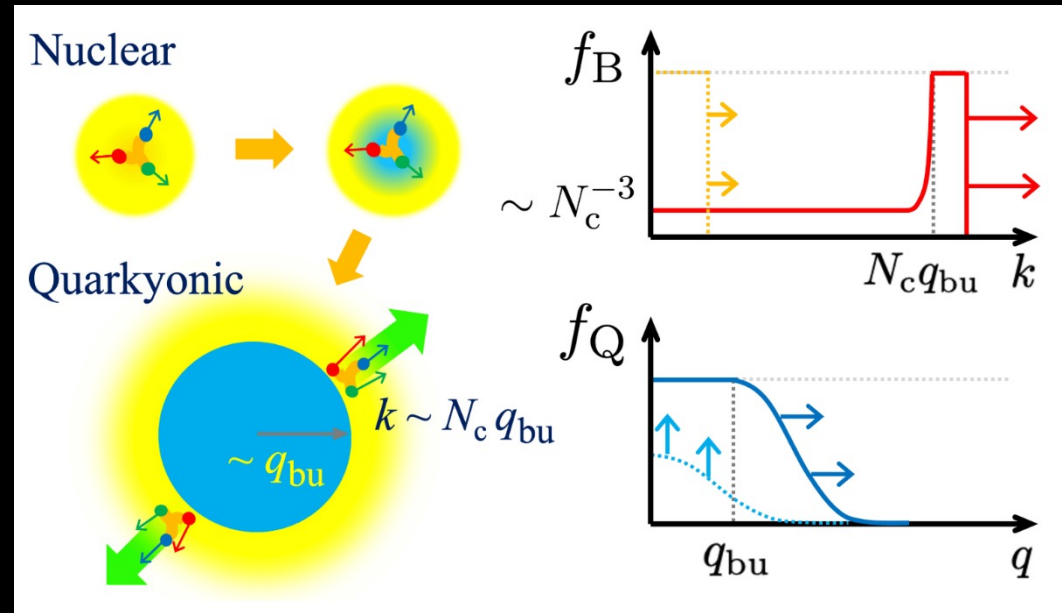
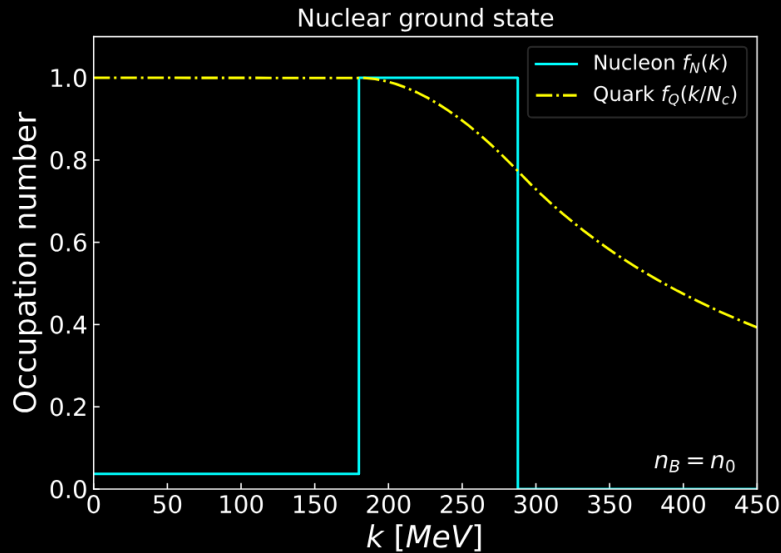
Realization of quarkyonic matter at T=0

Dual description

$$n_B = 4 \int_k f_B(k) = 4 \int_q f_Q(q)$$

$$\varepsilon = \varepsilon_B[f_B]|_{n_B} = \varepsilon_Q[f_Q]|_{n_B}$$

$$[f_Q(q)]_{f\sigma} = \sum_{i=n,p,\dots} \sum_{\sigma'=\uparrow,\downarrow} \int_k \left[\varphi\left(\mathbf{q} - \frac{\mathbf{k}}{N_c}\right) \right]_{f\sigma}^{i\sigma'} [f_B(k)]_{i\sigma'}$$

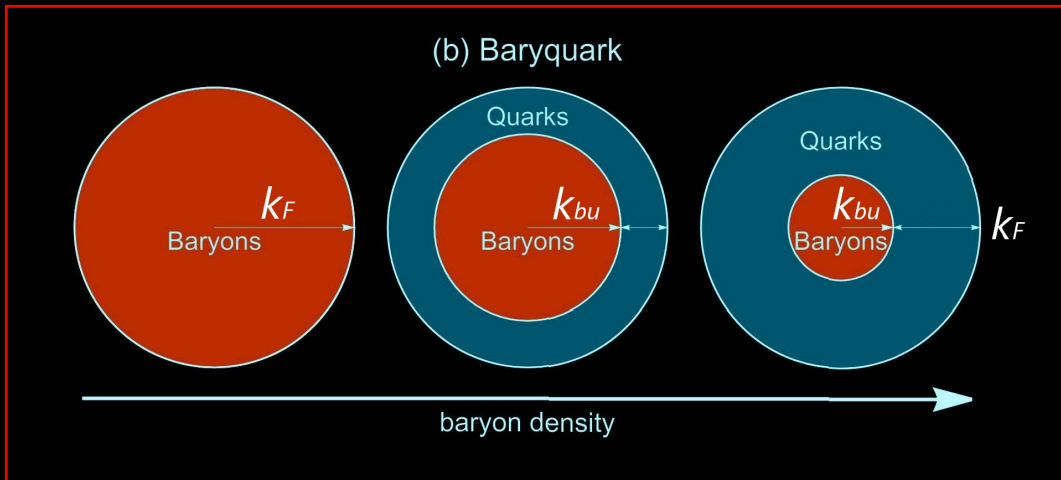
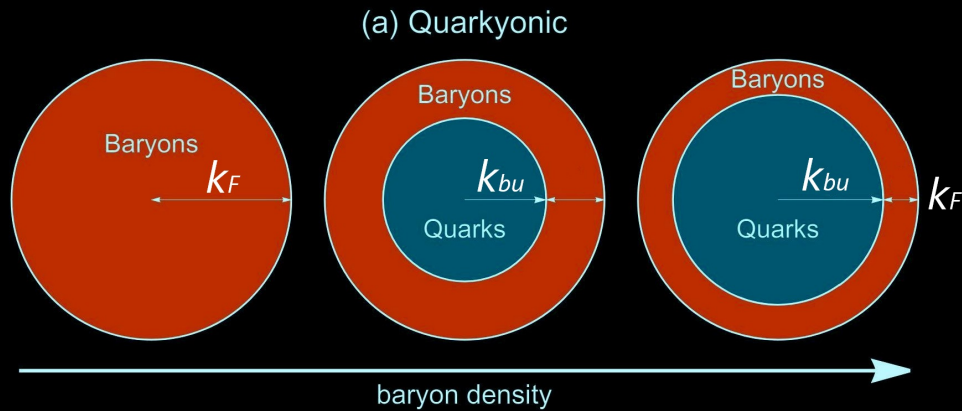


[Kojo,PRD,2021]

[Fujimoto, Kojo, McLerran, PRL, 2024]

[Koch, McLerran, Miller, Vovchenko, 2403.15375, 2024]

Shell structure in momentum space



$$n_N^{\text{id}}(k_F) = \frac{g}{2\pi^2} \int_0^{k_{\text{bu}}} k^2 dk,$$

$$\varepsilon_N^{\text{id}}(k_F) = \frac{g}{2\pi^2} \int_0^{k_{\text{bu}}} k^2 \varepsilon_N(k) dk,$$

$$n_Q^{\text{id}}(q_F) = \frac{g}{2\pi^2} \int_{k_{\text{bu}}/N_c}^{k_F/N_c} q^2 dq,$$

$$\varepsilon_Q^{\text{id}}(q_F) = N_c \frac{g}{2\pi^2} \int_{k_{\text{bu}}/N_c}^{k_F/N_c} q^2 \varepsilon_Q(q) dq.$$

Does not require infrared regulator
 The most energetically favorable configuration
 However, not yet consistent with dual approach

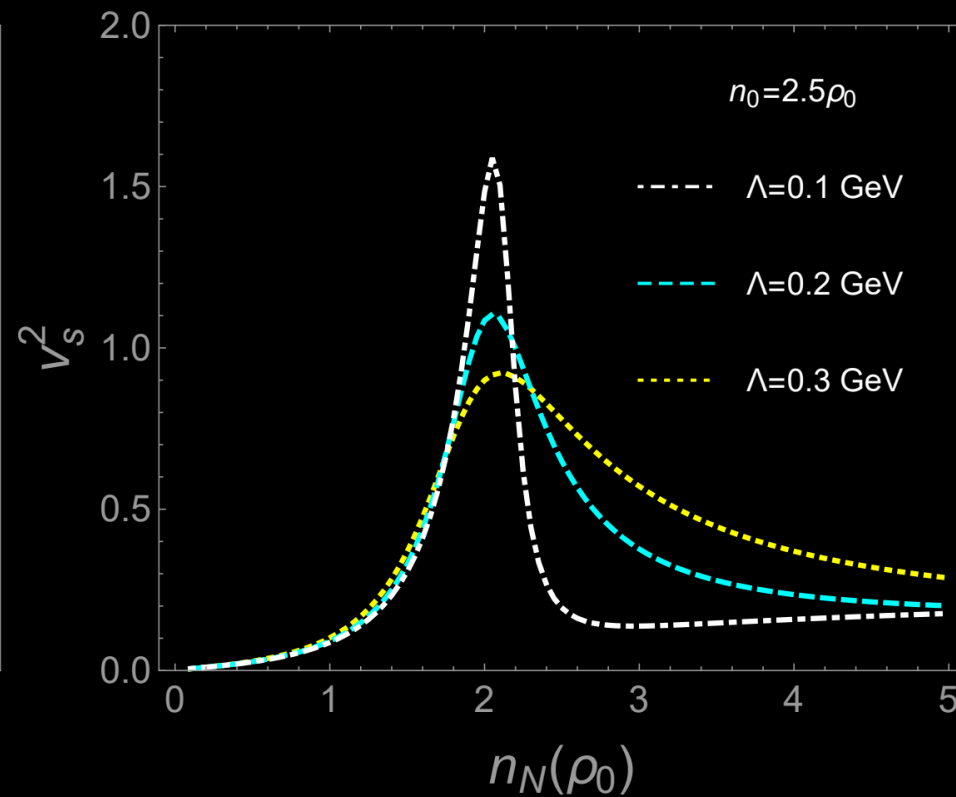
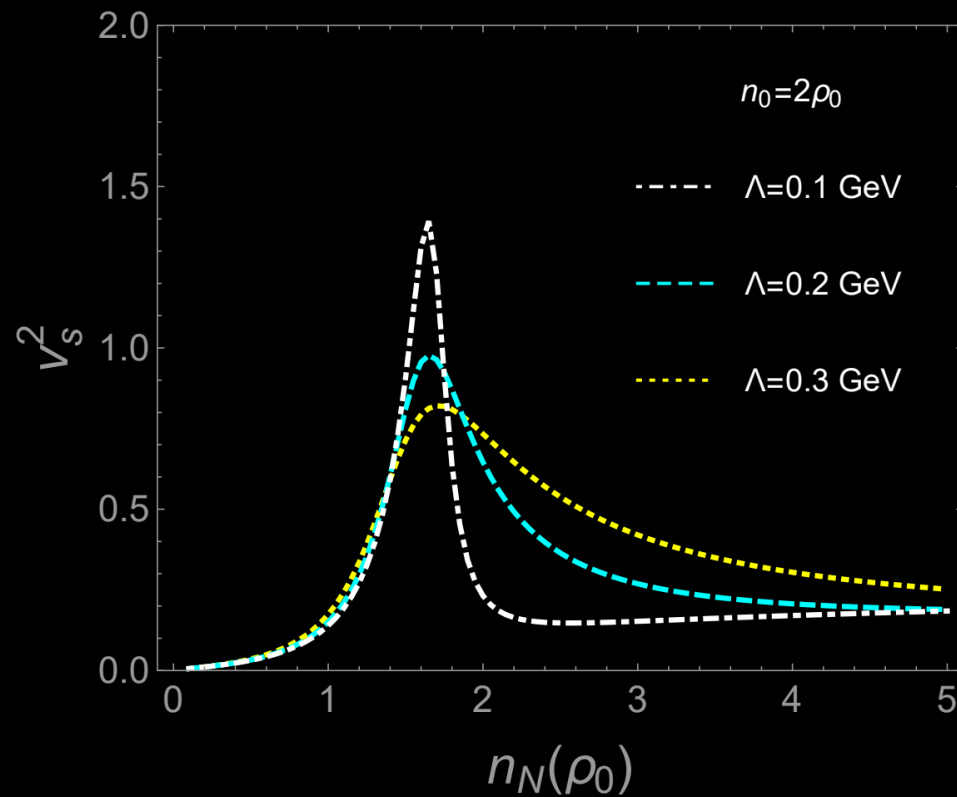
$f_q = n_q/n_B$ at each n_B is found by minimizing the energy density.

Quarkyonic matter: baryon excluded volume repulsion

$$n_{ex}^N = \frac{n_N^N}{1 - n_N^N/n_0}$$

$b \equiv 1/n_0$ is a free parameter

The role of attraction between baryons is unclear



[Kie Sang Jeong, McLerran, Sen, PRC, 2020]

Including nucleon-nucleon interactions: Quantum van der Waals theory

QvdW interactions

particle density: $n = f(n)n_{id}^*$

energy density: $\varepsilon = f(n)\varepsilon_{id}^* + nu(n)$

excl. volume factor: $f(n) = 1 - bn$

attractive mean-field: $u(n) = -an$

b – repulsion, a – attraction

Ground state of nuclear matter

$p = 0, \varepsilon/n_B \approx 922 \text{ MeV}, \Rightarrow a, b$

$n_B = \rho_0 \approx 0.16 \text{ fm}^{-3}$

[Vovchenko, Anchishkin, Gorenstein, PRC, 2015]

Modifications to the excluded volume

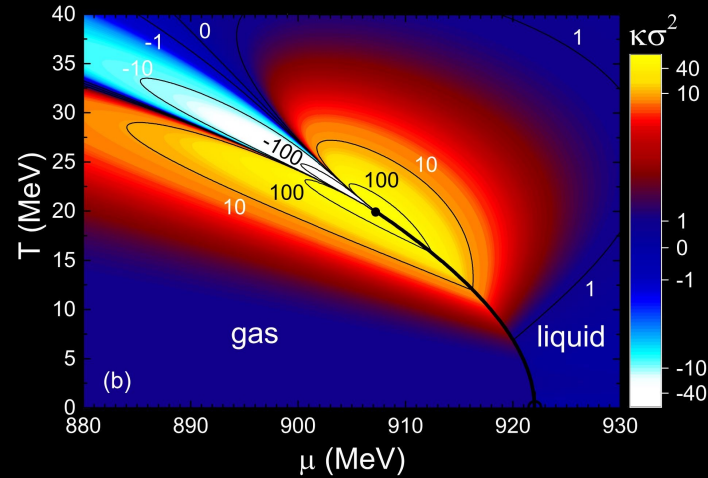
Carnahan-Starling (CS):

$$f_{CS}(n) = \exp \left[-\frac{3bn}{4 - bn} - \frac{4bn}{(4 - bn)^2} \right]$$

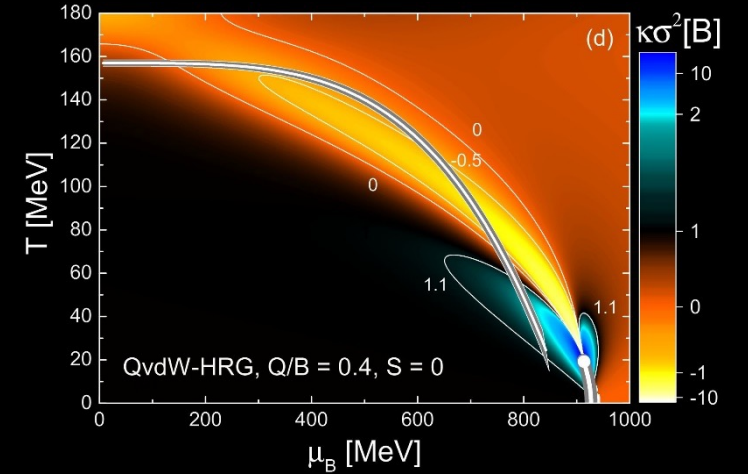
Trivirial Model (TVM):

$$f_{TVM}(n) = \exp \left[-bn - \frac{b^2 n^2}{2} \right]$$

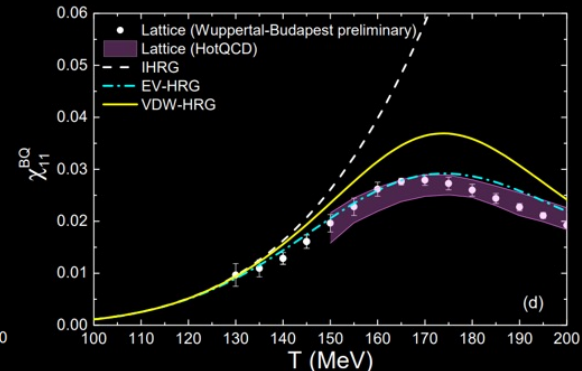
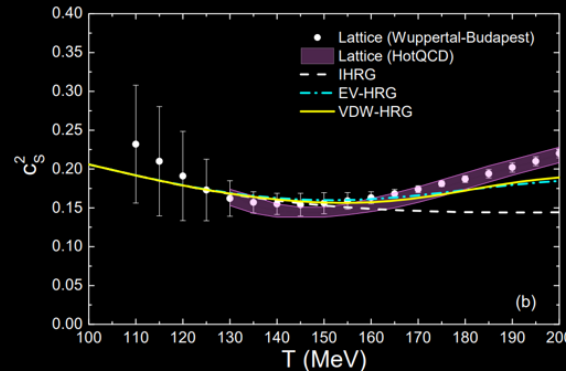
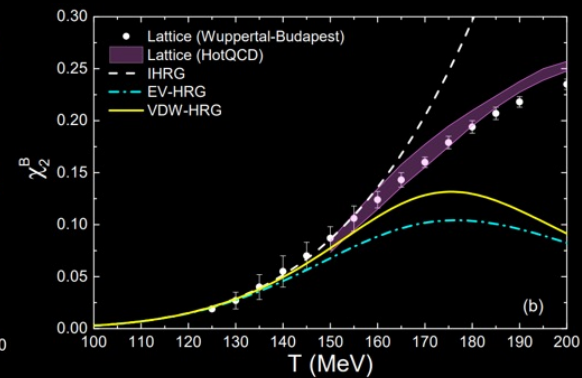
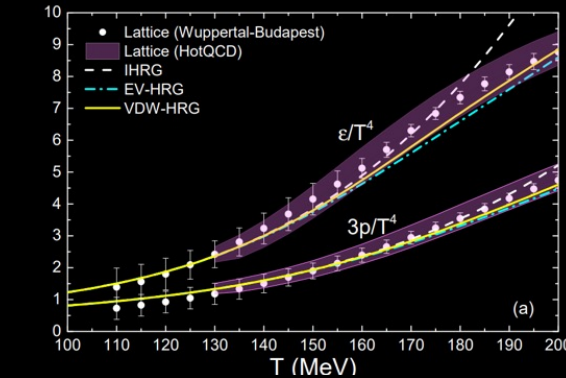
[Vovchenko, Greiner, Koch, Stoecker, PRD, 2020]



[Vovchenko, Anchishkin, Gorenstein, RP, PRC, 2015]

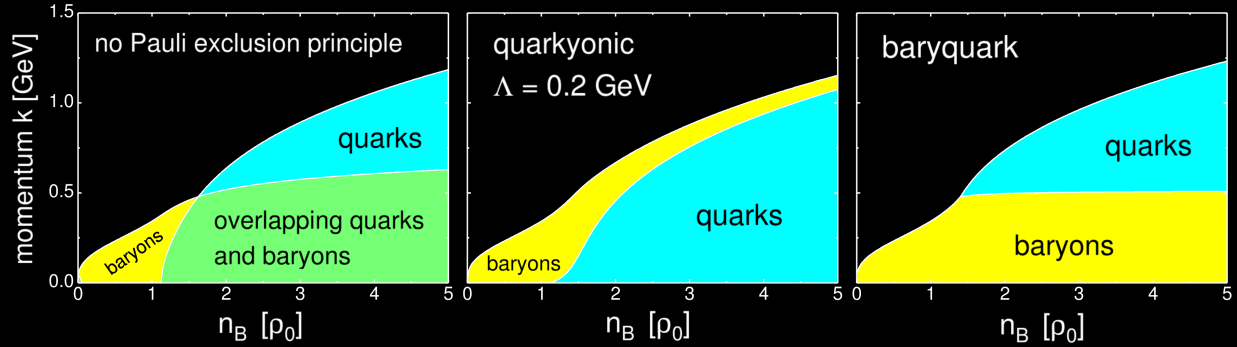


[RP, Vovchenko, Motorenko, Gorenstein, Stoecker, PRC, 2019]



[Vovchenko, Gorenstein, Stoecker, PRL, 2017]

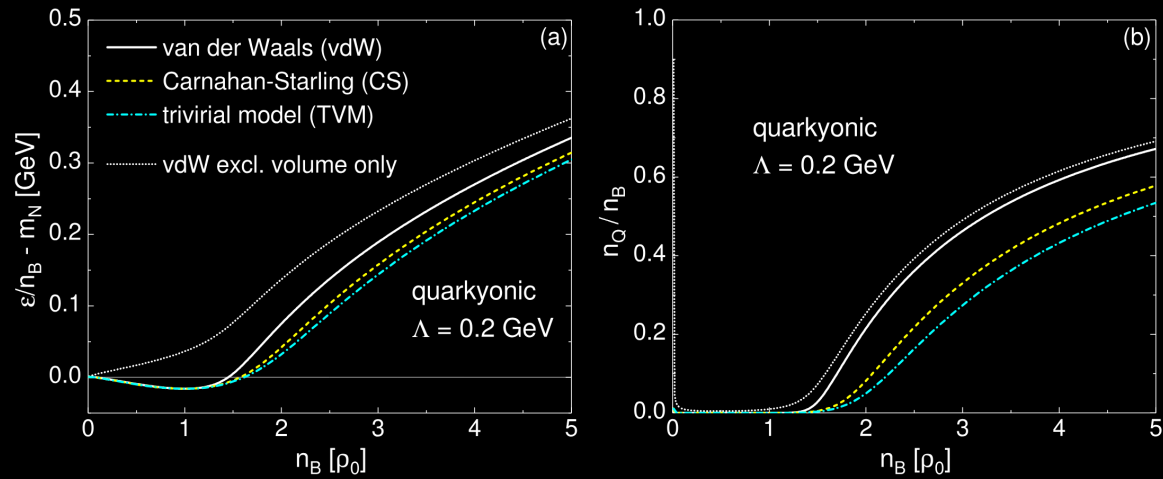
Quarkyonic matter with QvdW interactions for nucleons at $T = 0$



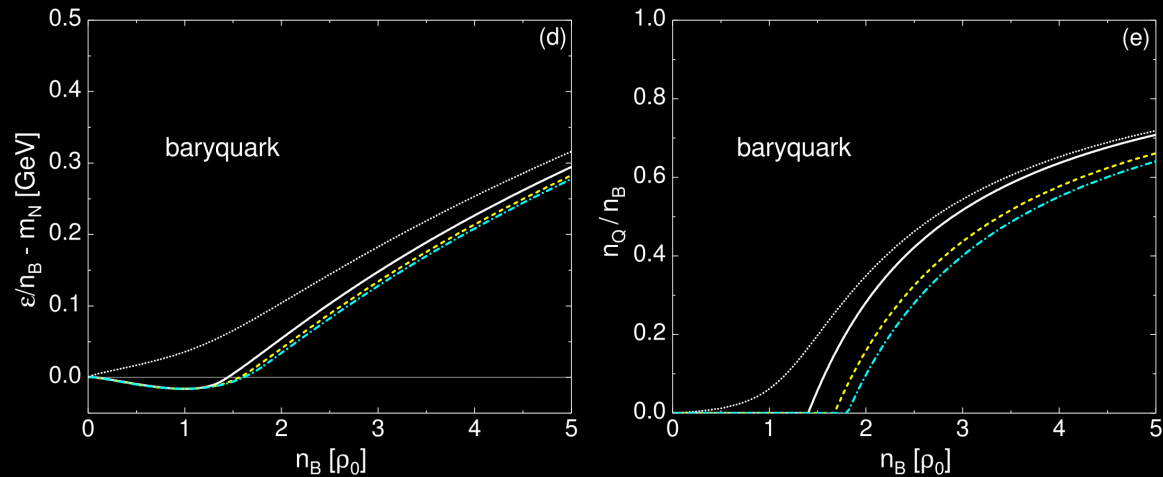
← matter composition in the momentum space

$$n_N = f(n_N) n_N^{\text{id}}(k_F),$$

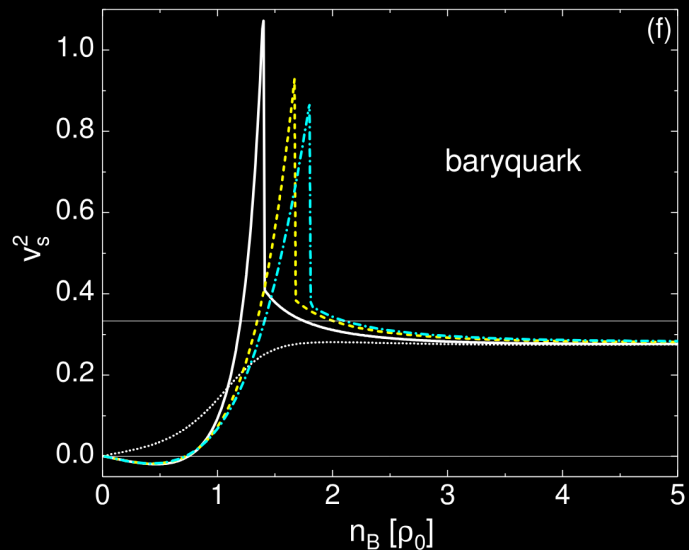
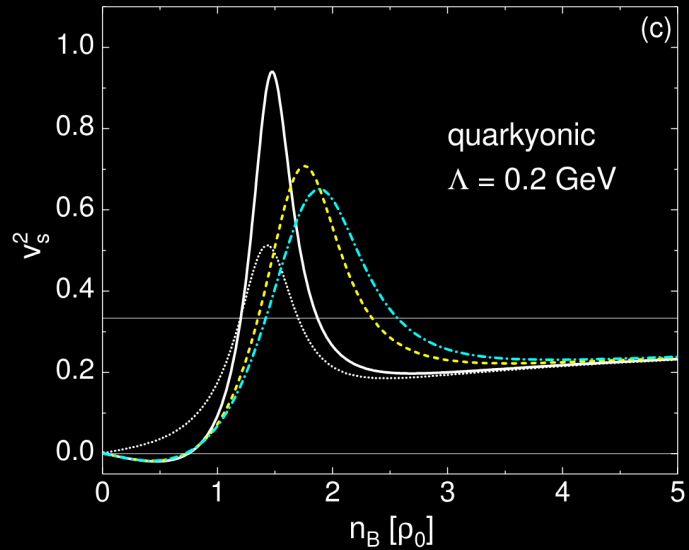
$$\varepsilon_N = f(n_N) \varepsilon_N^{\text{id}}(k_F) + n_N u(n_N)$$



- The transition to quarkyonic regime takes place at $n_B^{\text{tr}} = 1.5 - 2 \rho_0$
- In contrast to the earlier works, n_B^{tr} does not depend on any free parameter but stems from empirical properties of the nuclear ground state.
- The obtained estimates suggest that quarkyonic matter can be reachable in intermediate energy HIC.
- The transition is qualitatively similar for all considered modifications of EV regardless of whether they have limiting density.
- Baryquark configuration is the most energetically favorable in all considered scenarios. The quark onset is sudden with a kink in $f_q(n_B)$. At the same time, the resulting EoS are very similar.



Quarkyonic matter with QvdW interactions for nucleons: sound velocity



➤ In all considered cases the transition to quarkyonic regime results in the pronounced peak in sound velocity v_s^2 which rises well above the conformal limit.

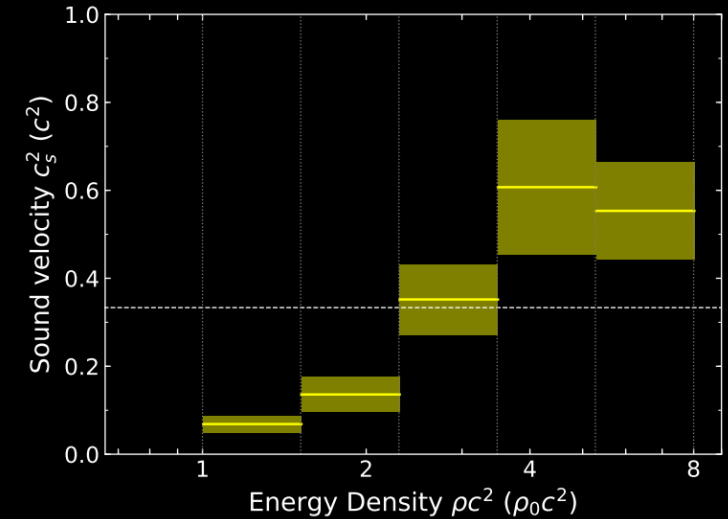
➤ The presence of nuclear attraction and the associated GS makes the peak in v_s^2 considerably sharper.

➤ We confirm that quarkyonic matter with the excluded volume mechanism requires the introduction of a regulator (Λ) to avoid the singular behavior in v_s^2 .

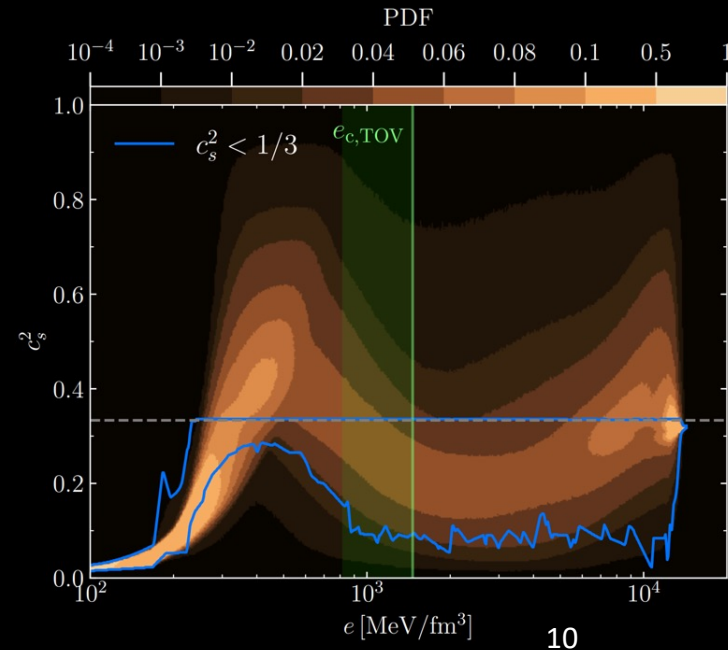
➤ For baryquark matter

- no singular behavior is observed even in the absence of the regulator.
- v_s^2 exhibits a discontinuous drop at the quark onset as a result of a kink.
- v_s^2 is causal for all considered cases except of vdW-EV where it slightly overshoots unity.

[RP, Stoecker, Vovchenko, PRC, 2023]



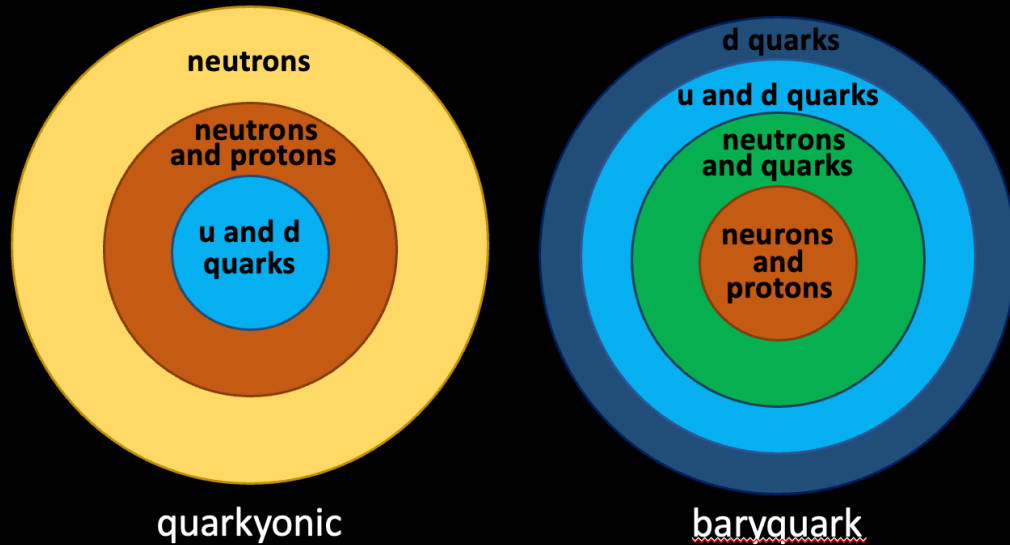
[Fujimoto, Fukushima, Murase, PRD, 2020]



[Altıparmak, Ecker, Rezzolla, Astrophys.J.Lett., 2022]

- Application to neutron stars [Tripp, RP, Vovchenko, in progress]
- Quark-baryon interactions (Kie Sang Jeong)
- Dual description of quarkyonic matter
- Inclusion of heavier baryons and quark flavors
- Chiral symmetry restoration
- Extension to finite temperatures, see e.g. [Sen, Warrington, Nucl.Phys.A,2021]

Application to neutron stars: asymmetric matter (in progress)



Assumptions:

- In the case of QY matter u and d quarks share the same Fermi surface, $q_{bu}^u = q_{bu}^d = q_{bu}$
- Asymmetry parameter:

$$y = \frac{\rho_Q}{\rho_B} = \frac{n_p}{n_p + n_n} = \frac{\frac{2}{3}n_u - \frac{1}{3}n_d}{\frac{1}{3}n_u + \frac{1}{3}n_d}$$

$$\varepsilon_N = \sum_{i \in p, n} f_{vdW}(x_i) \varepsilon_i^{id} - a_n(n_p^2 + n_n^2) - 2a_{pn}n_p n_n$$

$$n_{p,n} = f_{vdW}(x_{p,n}) n_{p,n}^{id}$$

$$x_p = b_n n_p + b_{pn} n_n$$

$$x_n = b_{pn} n_p + b_n n_n$$

$$f_{vdW}(x) = 1 - x$$

$$f_{CS}(x) = \exp \left[-\frac{3x}{4-x} - \frac{4x}{(4-x)^2} \right]$$

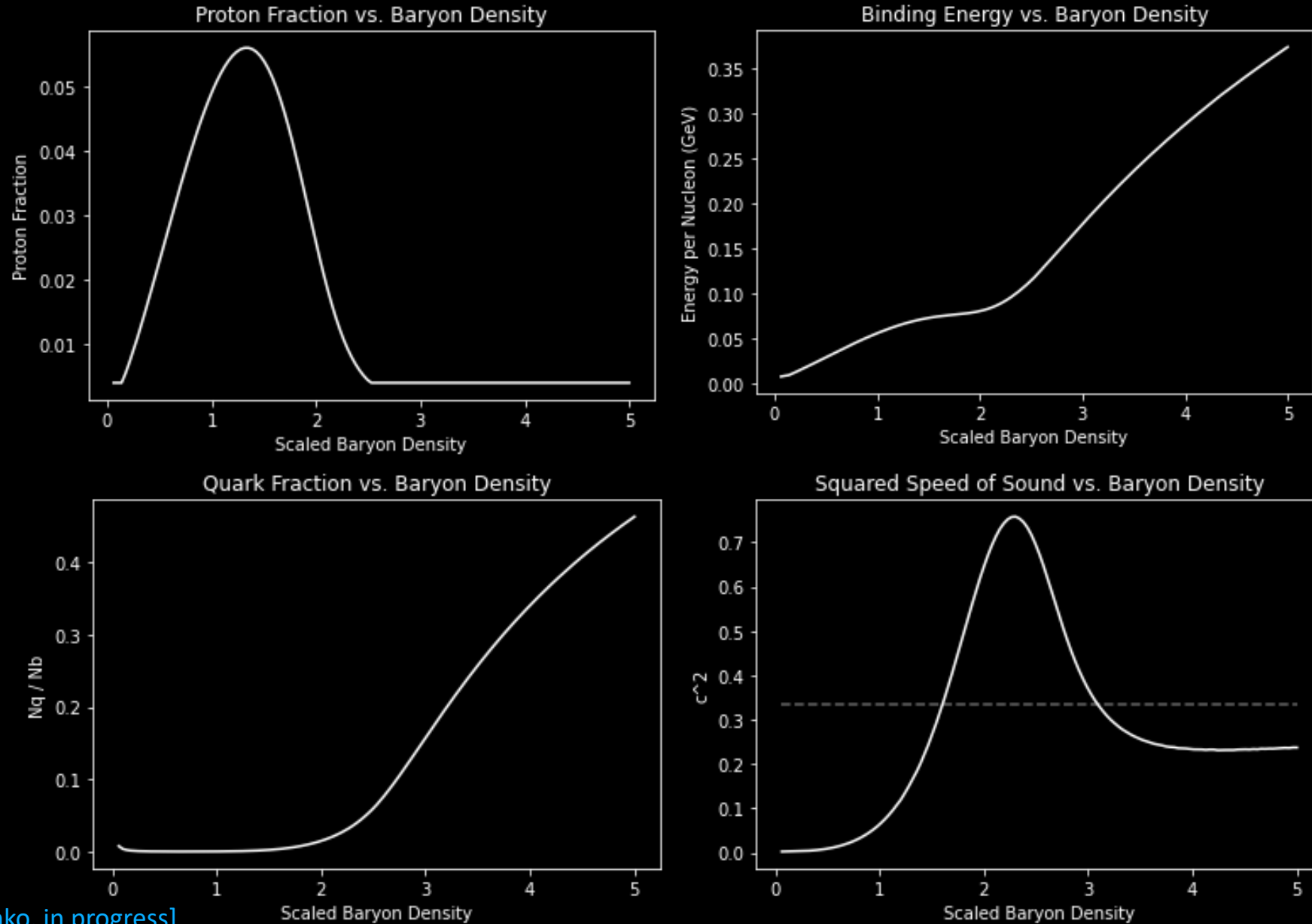
$$f_{TVM}(x) = \exp \left[-x - \frac{x^2}{2} \right]$$

$$a_{pn}, a_n, b_{pn}, b_n \leftarrow$$

$$\frac{a_{pn} + a_n}{2} = a, \quad \frac{b_{pn} + b_n}{2} = b$$

$$J \approx 32 \text{ MeV}, \quad L \approx 59 \text{ MeV}$$

Application to neutron stars (preliminary)



- Quarkyonic matter theory has been extended to include QvdW nucleon-nucleon interactions
- In contrast to previous works we incorporate the empirical properties of the nuclear matter
- The model predicts a transition to quarkyonic regime at densities around $1.5-2\rho_0$ with speed of sound rising significantly above the conformal limit. It suggests that QY matter is reachable in HICs.
- Baryquark matter is more energetically favorable than quarkyonic matter for all considered scenarios and does not need to rely on free parameters
- There is a number of topics that can be studied in future works based on the present study (e.g. application to neutron star phenomenology)

	vdW	CS	TVM	Experiment
a [MeV fm ³]	329	347	349	—
b [fm ³]	3.42	4.43	4.28	—
ρ_0 [fm ⁻³]	0.16	0.16	0.16	0.15 ± 0.01
K_0 [MeV]	763	601	564	250 - 315
T_c [MeV]	19.7	18.6	18.3	17.9 ± 0.4
n_c [fm ⁻³]	0.072	0.070	0.069	0.06 ± 0.01