Neutron Star Matter: EOS & Phase Structure



NSTITUTE for NUCLEAR THEORY



CPOD 2024 - 15th Workshop on Critical Point and Onset of Deconfinement

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S Network for Neutrinos, Nuclear Astrophysics, and Symmetries



QCD under extreme conditions



quark gluon plasma

dense quark matter

color superconductor

1015 1016 1014 1017 ρ (g/cm³)

QCD under extreme conditions



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 0^{13} 10¹⁴ 10¹⁵ 10¹⁶ 10¹⁷ $\rho(g/cm^3)$

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Low Temperature Phases of Matter

Coulomb Solid + **Relativistic Electrons:**

Coulomb Solid + **Relativistic Electrons +** Superfluid Neutrons

Thermal and Transport **Properties? Vortex Dynamics? Entrainment? Disorder?** Topology?

Nuclear Pas Non-Spheric Nuclei + Elect + Superflui **Neutrons**

Thermal an Transport **Properties** Superfluid a **Vortex Dynam** Entrainmen **Disorder**? Topology

1010

1011

1012

1013

939

 $\approx \mu_B(\text{MeV})$ 1000

ata: cal rons d d d d ad ? and aics? t?	<section-header><text></text></section-header>	Strongly coupled relative matter: Composition? Chiral Symmetry? Deconfinement? Quarkyonic? Superfluidity? Superconductivity?	vistic	Color Locked Ma Trans insu
?				р (g
10 ¹²	n_{sat} $c^2 < 1/3$	$\frac{10^{15}}{c_s^2/c^2} \simeq 1$	10 ⁴ c_s^2/	6

2000





Inside Neutron Stars







General Constraints on the Equation of State















Beane, Bedaque, Epelbaum, Kaplan, Machliedt, Meisner, Phillips, Savage, van Klock, Weinberg, Wise ...

Equation of State of Dense Nuclear Matter

Neutron Matter NLO 40 $N^{3}LC$ 30 per Particle (MeV) 20 10 Hebeler and Schwenk (2009), Gandolfi, Carlson, Reddy (2010), Gezerlis et al. **Nuclear Matter** (2013), Tews, Kruger, Hebeler, Schwenk (2013), Holt Kaiser, Weise (2013), 0 Hagen et al. (2013), Roggero, Mukherjee, Pederiva (2014), Wlazlowski, Holt, Energy Moroz, Bulgac, Roche (2014), Tews et al. (2018), Drischler et al., (2020). -10 -20 Nuclear **Saturation** 0.10.20.3Density $n \, [\mathrm{fm}^{-3}]$

Quantum many-body calculations of neutron matter and nuclear matter using EFT potentials show convergence up to about twice nuclear saturation density. Many-body perturbation theory and Quantum Monte Carlo methods have both been employed to calculate the energy on dense neutron matter. Drischler et al. used Bayesian methods to systematically estimate the EFT truncation errors in neutron and nuclear matter. Drischler, Furnstahl, Melendez, Phillips, (2020).

Equation of State of Neutron Star Matter

In neutron stars, matter is in equilibrium with respect to weak interactions and contains a small fraction (about 5-10%) of protons, electrons and muons:

Many-body perturbation theory and Bayesian estimates of the EFT truncation errors predict:

 $P_{\rm NSM}(n_B = 0.16 \text{ fm}^{-3}) = 3.0 \pm 0.2 \text{ MeV/fm}^3$ $P_{\rm NSM}(n_B = 0.34 \text{ fm}^{-3}) = 20.0 \pm 5 \text{ MeV/fm}^3$



Christian Drischler



Sophia Han



Tianqi Zhao





Bounds on Neutron Star Radii

EFT predictions for the EOS can be combined with extremal high-density EOS (with $c_s^2 = 1$) to derive robust bounds on the radius of a NS of any mass.

The lower limit on the NS maximum mass obtained from observations strengthen these bounds:

- $M_{\rm max} > 2.0 M_{\odot}$, 9.2 km < R_{1.4} < 13.2 km
- $M_{\rm max} > 2.6 M_{\odot}$, 11.2 km < R_{1.4} < 13.2 km

If $R_{1.4 is}$ small (<11.5 km) or large (>12.5 km), it would imply a very large speed of sound in the cores of massive neutron stars.





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Radius (km)









Speed of Sound in Dense Matter

<u>d</u> $\partial \epsilon$

Large maximum mass and observed radii, combined with neutron matter calculations suggests a rapid increase in pressure in the neutron star core.

This implies a large and nonmonotonic sound speed in dense QCD matter.

Suggests the existence of a strongly interacting phase of relativistic matter.



Tews, Carlson, Gandolfi and Reddy (2018) Steiner & Bedaque (2016)



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Is the NS core a fermi liquid or superconductor or superfluid ?

- The equation of state is insensitive to pairing at the Fermi surface.
- Need access to low-temperature properties such as response and specific heat to distinguish between phases.
- Neutrino cooling rates and heat capacity are exponentially supressed due to pairing at low temperature $\propto \exp(-\Delta/T)$
- Observations of accreting neutron stars can constrain the fraction of the baryons or quarks that can form Cooper pairs.



Transiently Accreting Neutron Stars



small variability

Deep Crustal Heating

During accretion nuclear reactions release: ~ 2-4 MeV / nucleon

Sato (1974), Haensel & Zdunik (1990), Brown, Bildsten Rutledge (1998) Gupta et al (2007,2011).

H/He burning r-p process ¹²C burning e- capture β⁻ decay n emission & capture fusion

density (g/cm³) depth (m) envelope 0.1 **10**⁵ 108 10 **10**² **10**¹¹ neutron drip superfluid neutrons **10**¹³ nuclear pasta **10**³ **10**¹⁴ core



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n emission & capture fusion

The energy is conducted to the core. Warms up old neutron stars. A steady state is reached as neutrino losses from the core balance crustal heating.

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KS 1731-260: High accretion 1988-2000

Image credit: NASA/CXC/Wijnands et al.

Measuring the Heat Capacity of the Core

Heat the star, allow it to relax, and observe the change in temperature:



Cumming et al. (2016)

 $C_{NS} dT = dQ$

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$$C_{NS} dT = dQ$$

$$_{\rm VS} = \alpha \ T: \quad \frac{\alpha}{2} \ (T_f^2 - T_i^2) = \Delta Q$$

limit:
$$C_{NS}(T_f) > 2 \frac{\Delta Q}{T_f}$$

$$t_H - L_{\nu} \times (t_H + t_{obs})$$

9 neutrino cooling rate duration of heating

time of observation (after heating ceases)

Constraining the Neutrino Emissivity and Heat Capacity

In systems with repeated accretion outbursts, we can infer the neutrino luminosity by estimating the net crustal heat deposition.

If cooling is observed between accretion events, it would constrain the specific heat. A 10% change in the core temperature on a 10 year time scale, would require that most fermions are gapped- either superfluid or superconducting. 10³⁹

> 10³⁶ 1



Brown at al. (2018)



- High sound speed requires strong repulsive interactions.
- In neutron matter, these interactions are modeled by the exchange of vector mesons Questions:
 - 1. Can fermions pair if their interactions are mediated by massive vector bosons? 2. Can the pairing gap be related to the sound speed?

Answers:

- 1. Yes. ${}^{3}P_{2}$ pairing is likely.
- 2. In some scenarios.



Pairing in the core (typical density $n_R \simeq 4 n_{sat}$)

Pairing through the Kohn-Luttinger Mechanism:





An old subject, first studies by Fay and Lazer in 1968!

see review by Gezerlis, Pethick and Schwenk arXiv:1406.6109 (2014)

Large short-range repulsion due heavy vector meson exchange naturally leads to pairing in the ${}^{3}P_{2}$ channel.

Interesting interplay between central and spin-orbit forces.

The strength of the vector repulsion is related to the sound speed.

$$c_s^2 = \frac{k_F^2}{3mm^*} \ (1 + F_0)$$

Induced Pairing in Neutron Matter





Kumamoto & Reddy (2024)

Summary & Conclusions

- speed in the neutron star core.
- stringent constraints.
- We will need to observe and interpret thermal phenomena in accreting (and magnetized) neutron stars to constrain low temperature properties.
- neutron stars appear to be cooling rapidly.
- \bullet
- the ${}^{3}P_{2}$ channel.

Mass and radius measurements can and will provide strong constraints on the sound

• X-ray observations can provide valuable guidance but GW's are needed to obtains

• There appears to some diversity in the cooling of neutron stars. A few old accreting

Long term monitoring of these systems can constrain the specific heat of the core.

• Screening of repulsive vector interactions provides a mechanism to pair fermions in