

Spin content of the proton at NNLO and beyond

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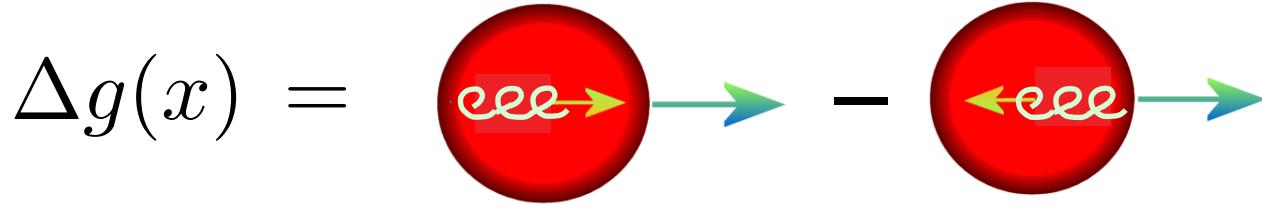
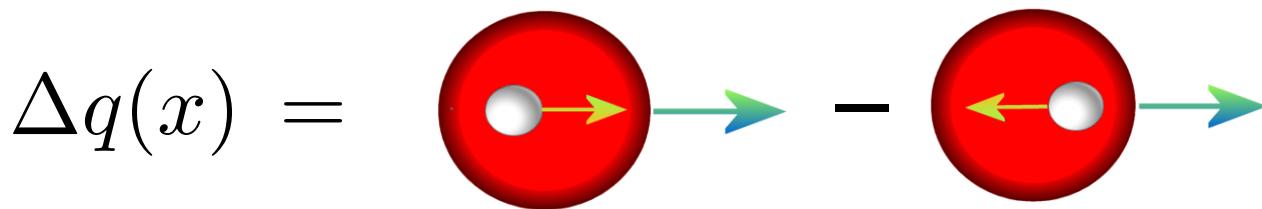
Berkeley, 09/18/2019

Outline:

- Introduction
- Evolution equations
- Singlet case
- Non-singlet case

Collaboration with D. de Florian (PRD 2019)

Introduction



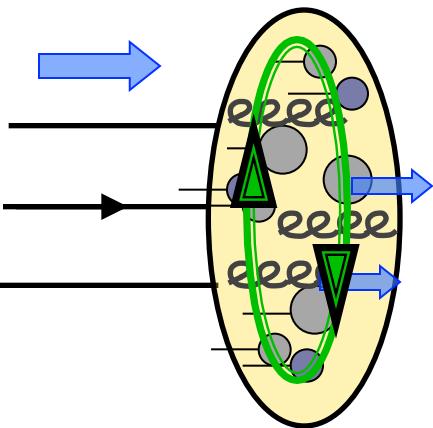
$$\Delta q(x) = \frac{1}{4\pi} \int dy^- e^{-iy^- xP^+} \langle P, S | \bar{\psi}(0, y^-, \mathbf{0}_\perp) \gamma^+ \gamma_5 \psi(0) | P, S \rangle$$

$$\Delta g(x) = \frac{1}{4\pi x P^+} \int dy^- e^{-iy^- xP^+} \langle P, S | F^{+\alpha}(0, y^-, \mathbf{0}_\perp) \tilde{F}_\alpha^+(0) | P, S \rangle$$

- in QCD: $\Delta q(x, \mu^2), \Delta g(x, \mu^2)$

- DGLAP evolution:

$$\mu^2 \frac{d}{d\mu^2} \begin{pmatrix} \Delta q(x, \mu^2) \\ \Delta g(x, \mu^2) \end{pmatrix} = \int_x^1 \frac{dz}{z} \begin{pmatrix} \Delta \mathcal{P}_{qq} & \Delta \mathcal{P}_{qg} \\ \Delta \mathcal{P}_{gq} & \Delta \mathcal{P}_{gg} \end{pmatrix} \begin{pmatrix} \Delta q \\ \Delta g \end{pmatrix} \left(\frac{x}{z}, \mu^2 \right)$$



$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma(Q^2) + \Delta G(Q^2) + L_q(Q^2) + L_g(Q^2)$$

Jaffe, Manohar; Ji, Hoodbhoy, Lu; Ji, Yuan;
Wakamatsu; Chen et al.; Lorce et al.;
Hatta

in terms of helicity PDFs $\Delta q(x, Q^2)$, $\Delta \bar{q}(x, Q^2)$, $\Delta g(x, Q^2)$:

$$\Delta\Sigma(Q^2) = \sum_q^{N_f} \int_0^1 dx \left(\Delta q(x, Q^2) + \Delta \bar{q}(x, Q^2) \right)$$

$$\Delta G(Q^2) = \int_0^1 dx \Delta g(x, Q^2)$$

from now on, consider only 1st moments
(obvious connection to small-x)

total orbital angular momentum: $\mathcal{L} \equiv L_q + L_g$

$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma(Q^2) + \Delta G(Q^2) + L_q(Q^2) + L_g(Q^2)$$

$$\Rightarrow \frac{d \mathcal{L}(Q^2)}{d \ln Q^2} = -\frac{1}{2} \frac{d \Delta\Sigma(Q^2)}{d \ln Q^2} - \frac{d \Delta G(Q^2)}{d \ln Q^2}$$

general structure of properties of orbital part clarified in

Hatta, Yoshida JHEP 2012

Hatta, Yao arXiv:1906.07744

Boussarie, Hatta, Yuan arXiv:1904.02693

Courtoy,...,Rajan PLB 2014

evolution of individual L_q and L_g only known to LO

Evolution equations

$$a, b \equiv u, \bar{u}, d, \bar{d}, s, \bar{s}, \dots, G$$

$$\frac{d\Delta a(Q^2)}{d \ln Q^2} = \sum_b \Delta P_{ab}(a_s(Q^2)) \Delta b(Q^2) \quad a_s \equiv \frac{\alpha_s}{4\pi}$$

$$\Delta P_{ab} = a_s \Delta P_{ab}^{(0)} + a_s^2 \Delta P_{ab}^{(1)} + a_s^3 \Delta P_{ab}^{(2)} + \mathcal{O}(a_s^4)$$

Ahmed, Ross
Altarelli, Parisi

Mertig, van Neerven
WV

**Moch, Rogal,
Vermaseren, Vogt**

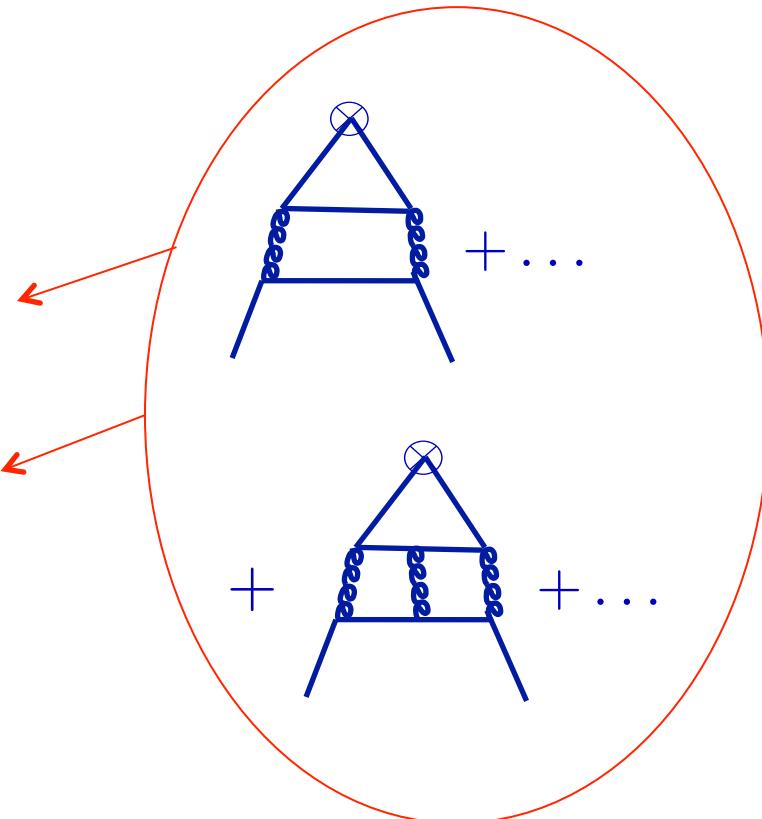
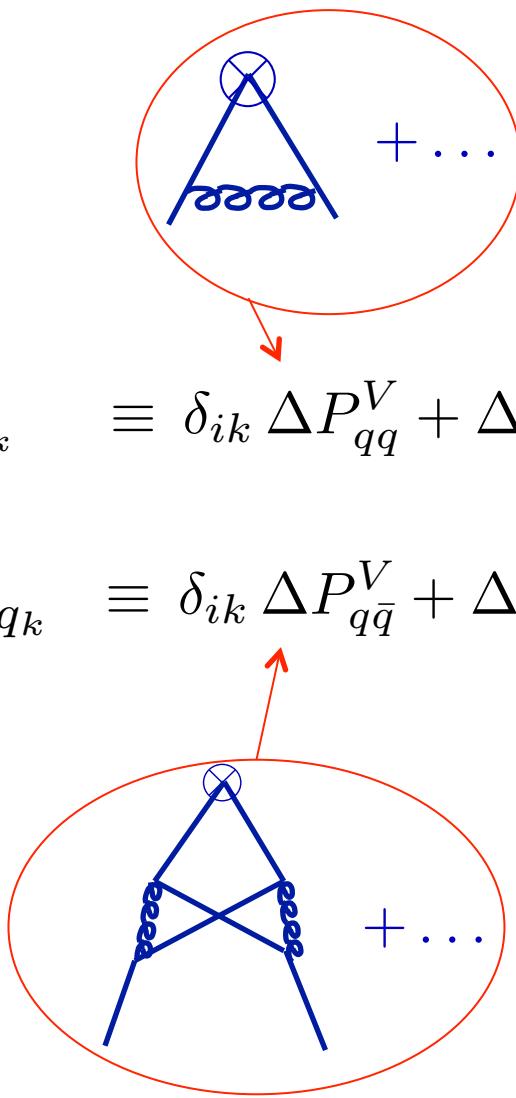
$$\frac{d \ln a_s(Q^2)}{d \ln Q^2} \equiv \frac{\beta(a_s)}{a_s} = -\beta_0 a_s - \beta_1 a_s^2 - \beta_2 a_s^3 + \mathcal{O}(a_s^4)$$

$$\beta_0 = \frac{11}{3}C_A - \frac{2}{3}N_f , \quad \beta_1 = \frac{34}{3}C_A^2 - \frac{10}{3}C_AN_f - 2C_FN_f , \quad \dots$$

to disentangle evolution equations, introduce:

$$\Delta P_{q_i q_k} = \Delta P_{\bar{q}_i \bar{q}_k} \equiv \delta_{ik} \Delta P_{qq}^V + \Delta P_{qq}^S$$

$$\Delta P_{q_i \bar{q}_k} = \Delta P_{\bar{q}_i q_k} \equiv \delta_{ik} \Delta P_{q\bar{q}}^V + \Delta P_{q\bar{q}}^S$$



$$\Delta P_{q_i q_k} = \Delta P_{\bar{q}_i \bar{q}_k} \equiv \delta_{ik} \Delta P_{qq}^V + \Delta P_{qq}^S$$

$$\Delta P_{q_i \bar{q}_k} = \Delta P_{\bar{q}_i q_k} \equiv \delta_{ik} \Delta P_{q\bar{q}}^V + \Delta P_{q\bar{q}}^S$$

and

$$\left. \begin{aligned} \Delta q^{(V)} &\equiv \sum_q (\Delta q - \Delta \bar{q}) \\ \Delta q^{(\pm)} &\equiv \Delta q \pm \Delta \bar{q} - \frac{1}{N_f} \sum_{q'} (\Delta q' \pm \Delta \bar{q}') \end{aligned} \right\} \text{non-singlet}$$

$$\Delta \Sigma = \sum_q (\Delta q + \Delta \bar{q}) \quad \text{singlet}$$

then:

$$\frac{d \Delta q^{(A)}(Q^2)}{d \ln Q^2} = \Delta P^{(A)}(a_s(Q^2)) \Delta q^{(A)}(Q^2) , \quad (A = V, \pm)$$

$$\frac{d}{d \ln Q^2} \begin{pmatrix} \Delta \Sigma \\ \Delta G \end{pmatrix} = \begin{pmatrix} \Delta P_{\Sigma\Sigma} & 2N_f \Delta P_{qG} \\ \Delta P_{Gq} & \Delta P_{GG} \end{pmatrix} \begin{pmatrix} \Delta \Sigma \\ \Delta G \end{pmatrix}$$

where:

$$\Delta P^{(V)} = \Delta P_{qq}^V - \Delta P_{q\bar{q}}^V + N_f (\Delta P_{qq}^S - \Delta P_{q\bar{q}}^S)$$

$$\Delta P^{(\pm)} = \Delta P_{qq}^V \pm \Delta P_{q\bar{q}}^V$$

$$\Delta P_{\Sigma\Sigma} \equiv \Delta P_{qq}^V + \Delta P_{q\bar{q}}^V + N_f (\Delta P_{qq}^S + \Delta P_{q\bar{q}}^S)$$



vanishes to all orders

e.g., $\Delta a_3 \equiv \Delta u^+ - \Delta d^+$ evolves with ΔP^+

$$S^\mu \Delta a_3 = \langle P, S | (\bar{u} \gamma^\mu \gamma_5 u - \bar{d} \gamma^\mu \gamma_5 d) | P, S \rangle$$

Bjorken; Kodaira et al.

Singlet case

from now on, focus on singlet:

$$\frac{d}{d \ln Q^2} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix} = \begin{pmatrix} \Delta P_{\Sigma\Sigma} & 2N_f \Delta P_{qG} \\ \Delta P_{Gq} & \Delta P_{GG} \end{pmatrix} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix}$$

$$\Delta P_{\Sigma\Sigma} = N_f (\Delta P_{qq}^S + \Delta P_{q\bar{q}}^S)$$

what's known (in $\overline{\text{MS}}$)?

Altarelli, Parisi;
Mertig, van Neerven; WV;
Moch, (Rogal),Vermaseren,Vogt

$$\frac{d}{d \ln Q^2} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix} = \left\{ a_s \begin{pmatrix} 0 & 0 \\ 3C_F & \beta_0 \end{pmatrix} + a_s^2 \begin{pmatrix} -6N_f C_F & 0 \\ C_F \left(\frac{71}{3}C_A - 9C_F - \frac{2}{3}N_f \right) & \beta_1 \end{pmatrix} + a_s^3 \begin{pmatrix} -2N_f C_F \left(\frac{71}{3}C_A - 9C_F - \frac{2}{3}N_f \right) & 0 \\ \# & \beta_2 \end{pmatrix} \right\} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix}$$

pattern:

$$\Delta P_{qG} = 0 \quad \Delta P_{GG} = -\frac{\beta(a_s)}{a_s}$$

$$\Delta P_{\Sigma\Sigma} = -2N_f a_s \Delta P_{Gq}$$

not really surprising:

Altarelli, Lampe '90

$$S^\mu \Delta\Sigma = \langle P, S | \bar{\psi} \gamma^\mu \gamma_5 \psi | P, S \rangle \equiv \langle P, S | j_5^\mu | P, S \rangle$$

$$\begin{aligned} \partial_\mu j_5^\mu &= 2N_f a_s \text{Tr} [F_{\mu\nu} \tilde{F}^{\mu\nu}] \\ &\equiv 2N_f a_s \partial_\mu \left\{ \varepsilon^{\mu\nu\rho\sigma} \text{Tr} \left[A^\nu \left(F^{\rho\sigma} - \frac{2}{3} A^\rho A^\sigma \right) \right] \right\} \\ &\quad \underbrace{\qquad\qquad\qquad}_{\equiv K^\mu} \end{aligned}$$

$$\Rightarrow j_5^\mu - 2N_f a_s K^\mu \quad \text{conserved}$$

$$\rightarrow \frac{d}{d \ln Q^2} \left(\Delta\Sigma + 2N_f a_s \Delta G \right) = 0$$

recall

$$\frac{d}{d \ln Q^2} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix} = \begin{pmatrix} \Delta P_{\Sigma\Sigma} & 2N_f \Delta P_{qG} \\ \Delta P_{Gq} & \Delta P_{GG} \end{pmatrix} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix}$$

$$\begin{aligned} \Rightarrow \frac{d}{d \ln Q^2} & \left(\Delta\Sigma(Q^2) + 2N_f a_s(Q^2) \Delta G(Q^2) \right) \\ &= \left(\Delta P_{\Sigma\Sigma} + 2N_f a_s \Delta P_{Gq} \right) \Delta\Sigma \\ &+ \left(2N_f \Delta P_{qG} + \Delta P_{GG} + \frac{\beta(a_s)}{a_s} \right) 2N_f a_s \Delta G \end{aligned}$$

thus

$$\Delta P_{qG} = 0 \quad \Delta P_{GG} = -\frac{\beta(a_s)}{a_s}$$

$$\Delta P_{\Sigma\Sigma} = -2N_f a_s \Delta P_{Gq}$$

evolution equation becomes:

$$\frac{d}{d \ln Q^2} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix} = \begin{pmatrix} \Delta P_{\Sigma\Sigma}(a_s) & 0 \\ -\frac{1}{2N_f a_s} \Delta P_{\Sigma\Sigma}(a_s) & -\frac{\beta(a_s)}{a_s} \end{pmatrix} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix}$$

write in terms of

$$\Delta\Gamma(Q^2) \equiv a_s(Q^2)\Delta G(Q^2)$$

$$\frac{d}{d \ln Q^2} \begin{pmatrix} \Delta\Sigma \\ \Delta\Gamma \end{pmatrix} = \begin{pmatrix} \Delta P_{\Sigma\Sigma}(a_s) & 0 \\ -\frac{1}{2N_f} \Delta P_{\Sigma\Sigma}(a_s) & 0 \end{pmatrix} \begin{pmatrix} \Delta\Sigma \\ \Delta\Gamma \end{pmatrix}$$

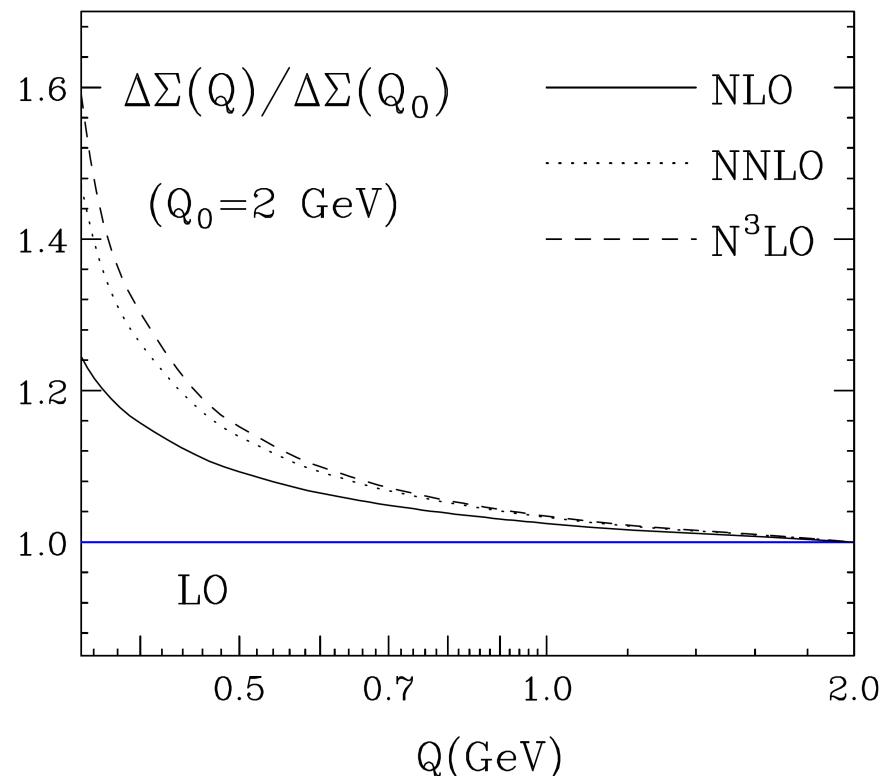
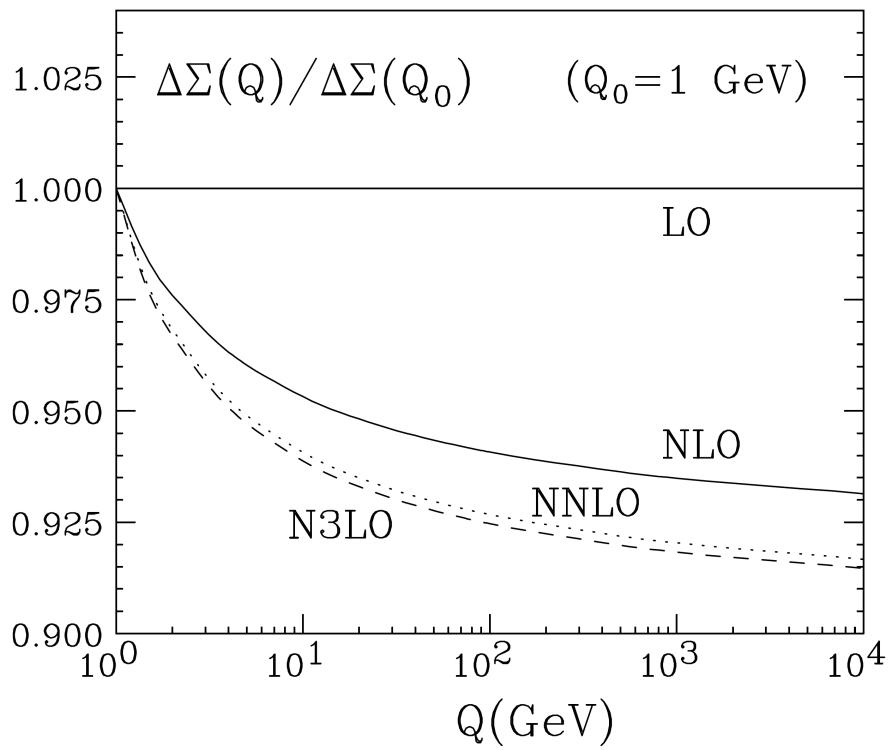
- single anomalous dimension controls all evolution
- straightforward to solve analytically

$$\Delta P_{\Sigma\Sigma} \equiv -2N_f a_s \Delta P_{gq}$$

from known $\Delta P_{gq}^{(2)}$ obtain

$$\begin{aligned}\Delta P_{\Sigma\Sigma}^{(3)} = -2N_f \Delta P_{Gq}^{(2)} &= -2N_f C_F \left[\frac{1607}{12} C_A^2 - \frac{461}{4} C_F C_A + \frac{63}{2} C_F^2 \right. \\ &\quad \left. + \left(\frac{41}{3} - 72\zeta_3 \right) C_A N_f - \left(\frac{107}{2} - 72\zeta_3 \right) C_F N_f - \frac{13}{3} N_f^2 \right]\end{aligned}$$

$$\begin{aligned}
\frac{\Delta\Sigma(Q^2)}{\Delta\Sigma(Q_0^2)} &= \exp[0] \times \exp\left[-\frac{a_Q - a_0}{\beta_0} \Delta P_{\Sigma\Sigma}^{(1)}\right] \times \exp\left[\frac{a_Q^2 - a_0^2}{2\beta_0^2} \left(\beta_1 \Delta P_{\Sigma\Sigma}^{(1)} - \beta_0 \Delta P_{\Sigma\Sigma}^{(2)}\right)\right] \\
&\times \exp\left[\frac{a_Q^3 - a_0^3}{3\beta_0^3} \left(-\beta_1^2 \Delta P_{\Sigma\Sigma}^{(1)} + \beta_0 \beta_2 \Delta P_{\Sigma\Sigma}^{(1)} + \beta_0 \beta_1 \Delta P_{\Sigma\Sigma}^{(2)} - \beta_0^2 \Delta P_{\Sigma\Sigma}^{(3)}\right)\right] \\
&\equiv K^{(\text{LO})} \times K^{(\text{NLO})} \times K^{(\text{NNLO})} \times K^{(\text{N}^3\text{LO})}
\end{aligned}$$



gluon spin:

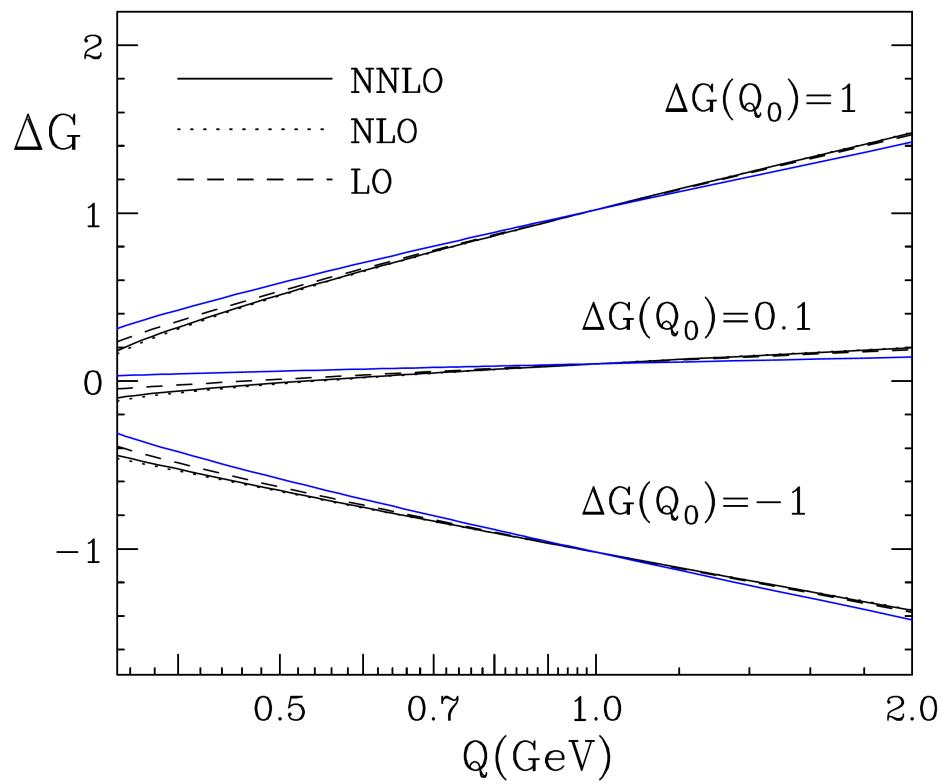
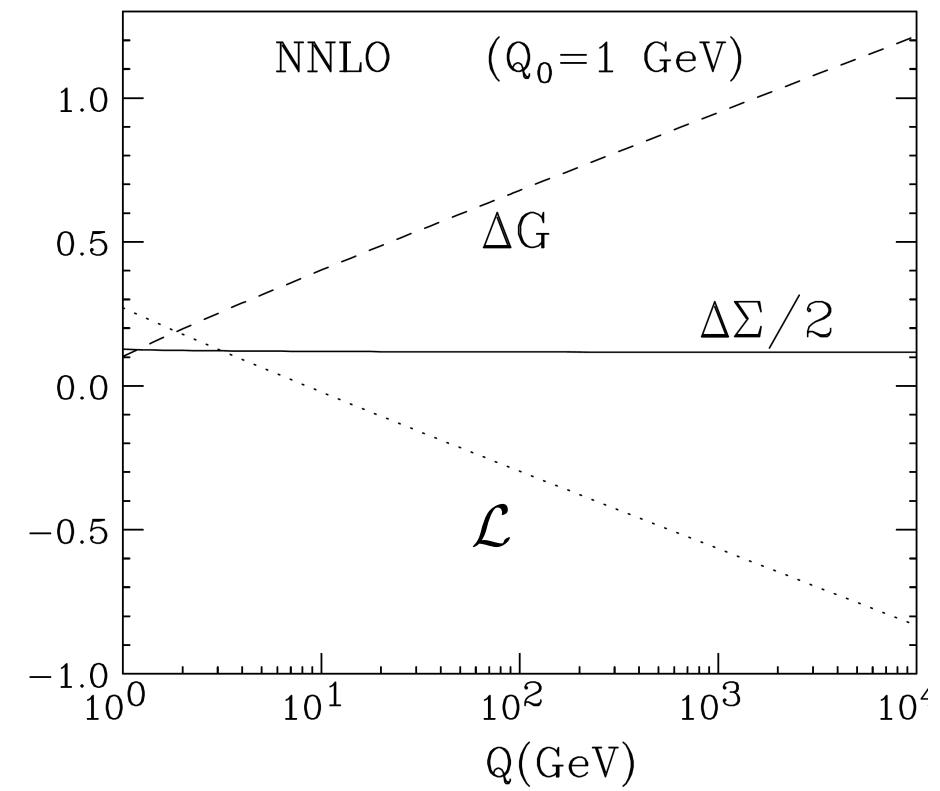
$$\Delta G(Q^2) = \frac{a_s(Q_0^2)}{a_s(Q^2)} \Delta G(Q_0^2) + \Delta \Sigma(Q_0^2) F\left(a_s(Q^2), a_s(Q_0^2)\right)$$

$$F(a_Q, a_0) = F^{\text{LO}}\left(\frac{a_0}{a_Q}\right) + a_Q F^{\text{NLO}}\left(\frac{a_0}{a_Q}\right) + a_Q^2 F^{\text{NNLO}}\left(\frac{a_0}{a_Q}\right)$$

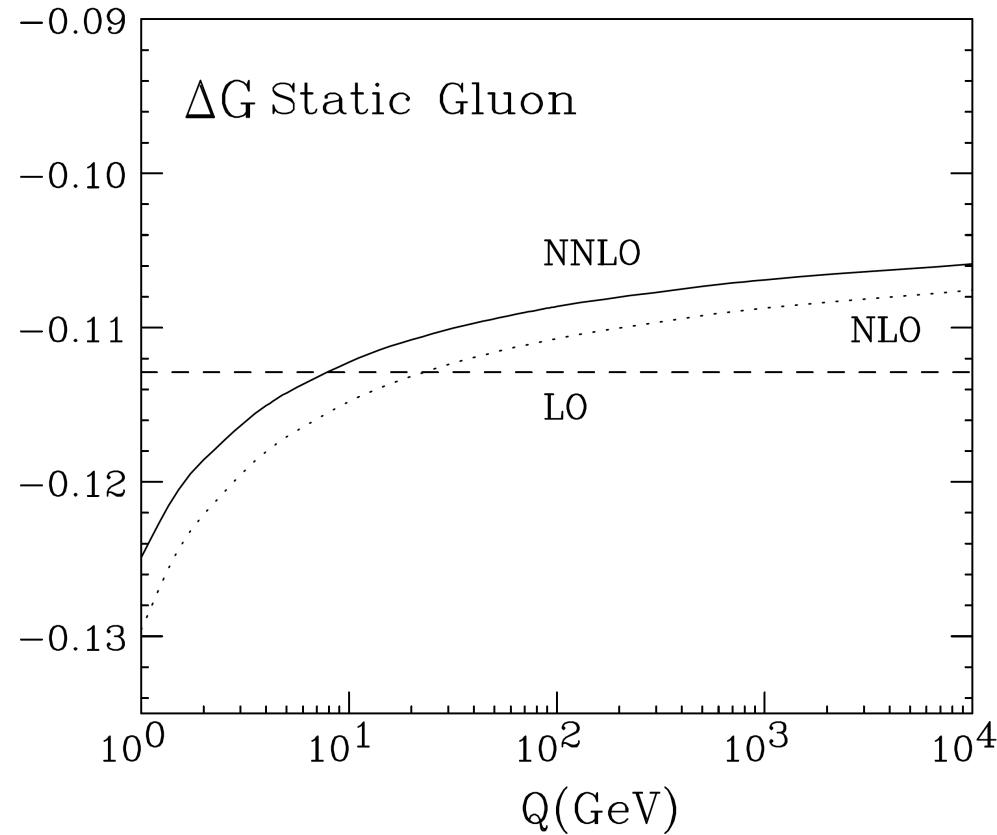
$$F^{\text{LO}}(r) = -(1-r) \frac{\Delta P_{Gq}^{(0)}}{\beta_0}$$

$$F^{\text{NLO}}(r) = \frac{1-r^2}{2\beta_0^2} \left(\beta_1 \Delta P_{Gq}^{(0)} - \beta_0 \Delta P_{Gq}^{(1)} \right) + \frac{(1-r)^2}{2\beta_0^2} \Delta P_{Gq}^{(0)} \Delta P_{\Sigma\Sigma}^{(1)}$$

$|\Delta G(Q^2)|$ will in general rise as $1/\alpha_s(Q^2)$



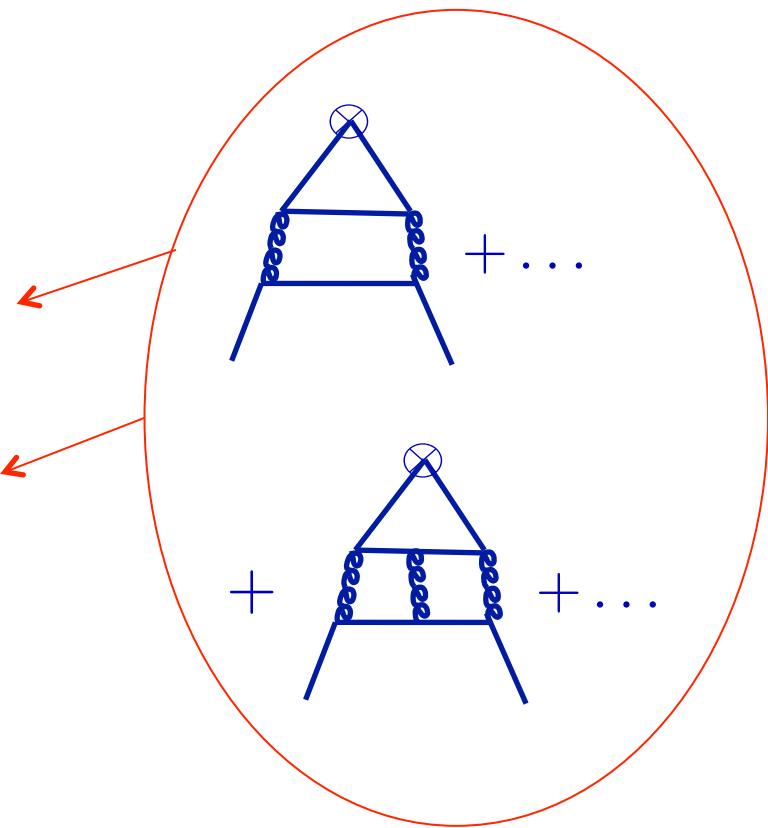
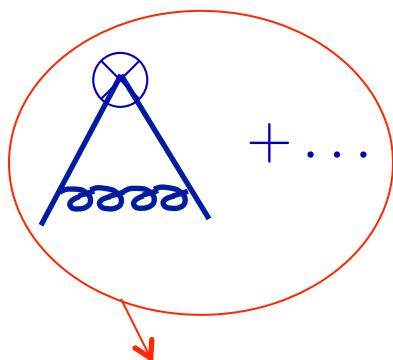
“fine-tuned” input $\Delta G(Q_0^2) \sim -0.1$



Non-singlet case

$$\Delta P_{q_i q_k} = \Delta P_{\bar{q}_i \bar{q}_k} = \delta_{ik} \Delta P_{qq}^V + \Delta P_{qq}^S$$

$$\Delta P_{q_i \bar{q}_k} = \Delta P_{\bar{q}_i q_k} = \delta_{ik} \Delta P_{q\bar{q}}^V + \Delta P_{q\bar{q}}^S$$



$P_{qq}^S \neq P_{q\bar{q}}^S$ generates strangeness asymmetry $s \neq \bar{s}$

Catani, de Florian, Rodgrigo, WV

$$\Delta q^{(V)} \equiv \sum_q (\Delta q - \Delta \bar{q})$$

$$\Delta q^{(\pm)} \equiv \Delta q \pm \Delta \bar{q} - \frac{1}{N_f} \sum_{q'} (\Delta q' \pm \Delta \bar{q}')$$

$$\frac{d \Delta q^{(A)}(Q^2)}{d \ln Q^2} = \Delta P^{(A)}(a_s(Q^2)) \Delta q^{(A)}(Q^2) , \quad (A = V, \pm)$$

$$U^{(A)}(Q, Q_0) = \exp \left\{ \int_{Q_0^2}^{Q^2} \frac{dq^2}{q^2} \Delta P^{(A)}(a_s(q^2)) \right\}$$

$$(\Delta q - \Delta \bar{q})(Q^2) = U^{(-)}(Q, Q_0) \left[(\Delta q - \Delta \bar{q})(Q_0^2) + \underbrace{\frac{1}{N_f} \left(\frac{U^{(V)}(Q, Q_0)}{U^{(-)}(Q, Q_0)} - 1 \right) \Delta q^{(V)}(Q_0^2)}_{total valence} \right]$$

total valence

$$\begin{aligned}
(\Delta s - \Delta \bar{s})_{\text{pert}}(Q^2) &= -\frac{\Delta P_{qq}^{(2)S} - \Delta P_{q\bar{q}}^{(2)S}}{2\beta_0} (a_Q^2 - a_0^2) (\Delta u - \Delta \bar{u} + \Delta d - \Delta \bar{d})(Q_0^2) \\
&= -\frac{5(23 - 12\zeta_2 - 16\zeta_3)}{72\beta_0\pi^2} (\alpha_s(Q^2) - \alpha_s(Q_0^2)) (\Delta u - \Delta \bar{u} + \Delta d - \Delta \bar{d})(Q_0^2)
\end{aligned}$$

