

Regge approach to light meson photoproduction

Glòria Montaña

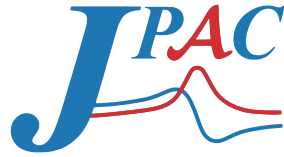
Theory Center, Thomas Jefferson National Accelerator Facility

Multi-Particle Reactions Workshop 2025, UC Berkeley, July 28, 2025



Joint Physics Analysis Center

2



Spectroscopy analyses

→ Understand production mechanisms

- **Reggeization of pion exchange**
- **Double Regge exchange**



Adam Szczepaniak
Indiana University / Jefferson Lab



Alessandro Pilloni
Università di Messina



Alex Akridge
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Arkaitz Rodas Bilbao
Old Dominion University /
Jefferson Lab



César Fernández
Ramírez
Universidad Nacional de
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Emilie Passemar
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Giorgio Foti
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Mikhail Mikhasenko
Ruhr-Universität Bochum



Nadine Hammoud
University of Barcelona



Robert Perry
Massachusetts Institute of
Technology



Sergi González-Solís
University of Barcelona



Vanamali Shastry
Indiana University



Viktor Mokeev
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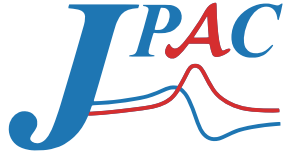
Vincent Mathieu
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This talk



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Regge theory 101

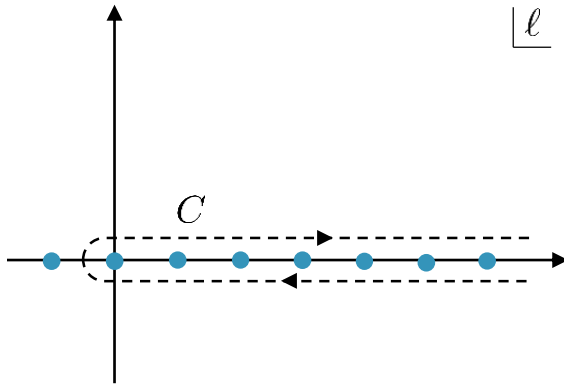
[T. Regge, *Nuovo Cim.* 18 (1960) 947-956]

Partial wave expansion

$$A(s, t) = \sum_{\ell=0}^{\infty} (2\ell + 1) P_{\ell}(z_t) f_{\ell}(t)$$

Sommerfeld-Watson transform

$$A(s, t) = -\frac{1}{2i} \int_C d\ell \frac{(2\ell + 1) P_{\ell}(-z_t) f(\ell, t)}{\sin \pi \ell}$$



Regge theory 101

3

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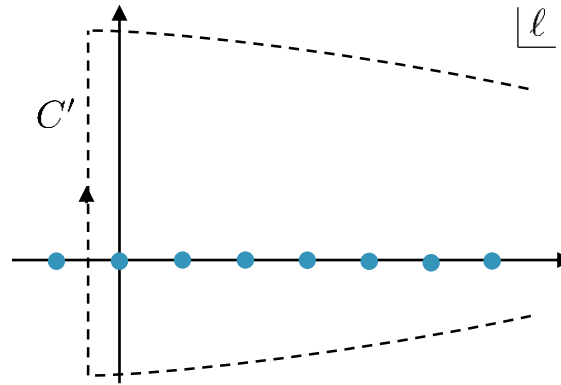
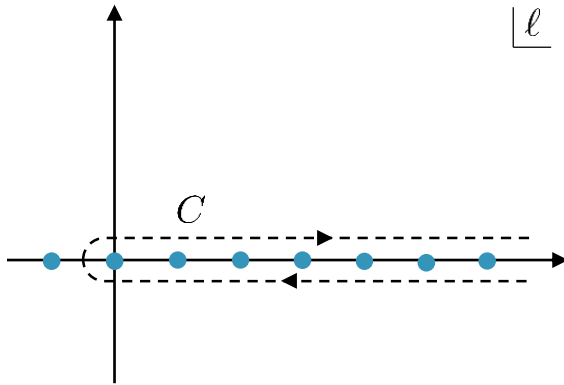
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Deform the contour

(assuming the partial waves are analytic, only singularities are poles)



Regge theory 101

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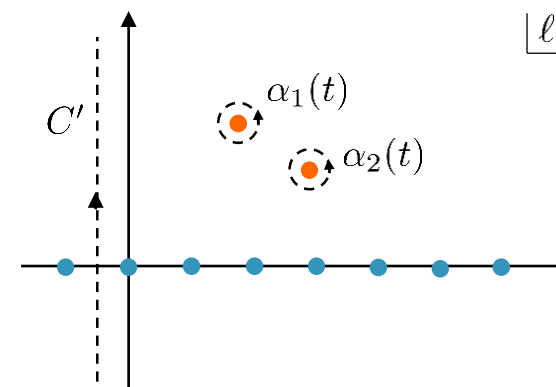
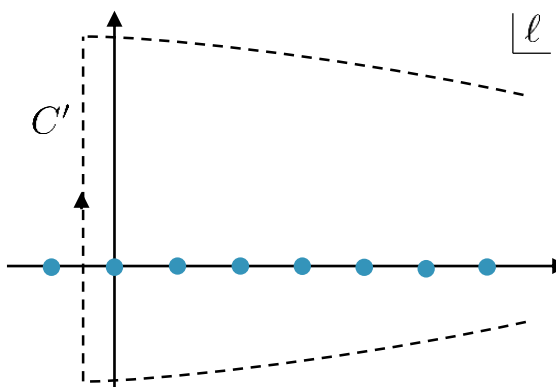
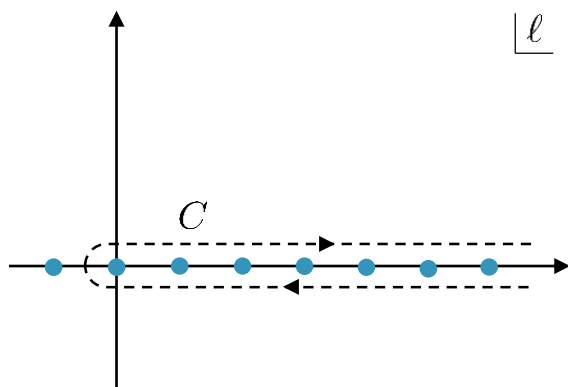
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$$A(s, t) = \underbrace{-\frac{1}{2i} \int_{-\frac{1}{2}-i\infty}^{-\frac{1}{2}+i\infty} d\ell \dots}_{\text{background} \sim s^{-1/2}} - \underbrace{\sum_i \frac{\pi(2\alpha_i(t) + 1)\beta_i(t)}{\sin(\pi\alpha_i^{\pm}(t))} \frac{1}{2} P_{\alpha_i}(-z_t)}_{\text{pole contributions} \sim s^{\alpha(t)}}$$

Regge poles



$$\beta_i(t) = \text{Res}_{\ell \rightarrow \alpha_i(t)} f_{\ell}(t)$$

Regge theory 101

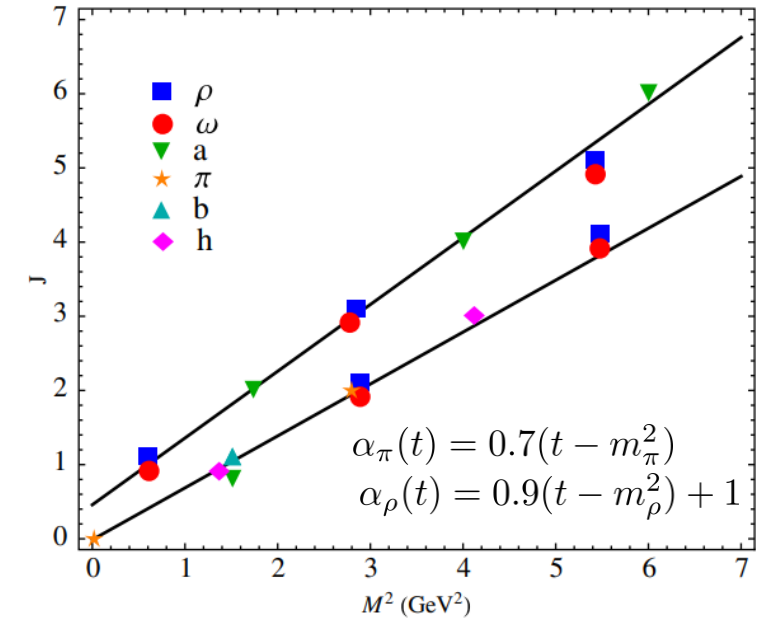
[T. Regge, *Nuovo Cim.* 18 (1960) 947-956]

Scattering amplitudes are analytic in s and t $\longrightarrow$ Partial waves must be analytic in angular momentum

$$f_\ell(t) = \int \frac{dz}{2} P_\ell(z) A(z, t) \sim \frac{\beta(t)}{\ell - \alpha(t)}$$

Resonances appear simultaneously as poles in energy and spin!

- **Regge trajectories:** families with same quantum numbers but different spin
- Almost straight lines (Chew-Frautschi plot)
- In standard Regge theory parameterized by: $\alpha(t) = \alpha' t + \alpha_0$



[V.Mathieu et al., *Phys.Rev.D* 98 (2018) 1, 014041]

Regge theory 101

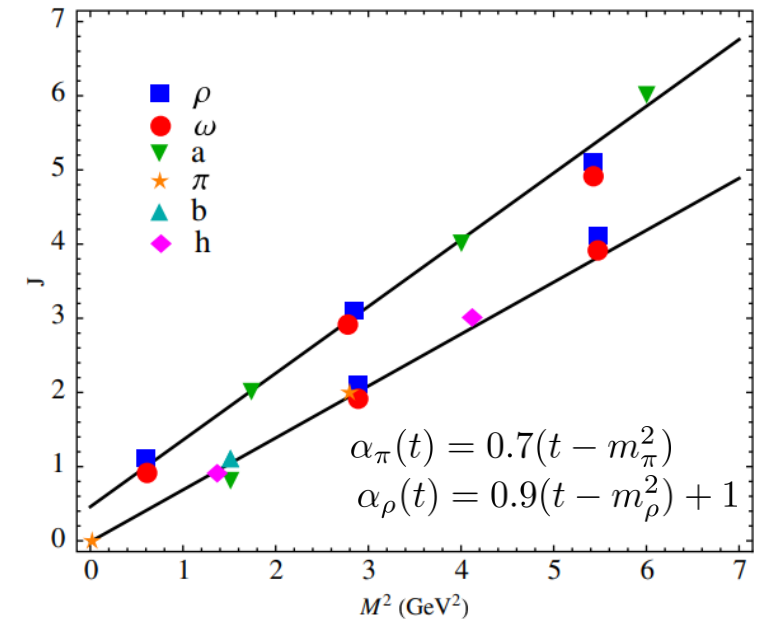
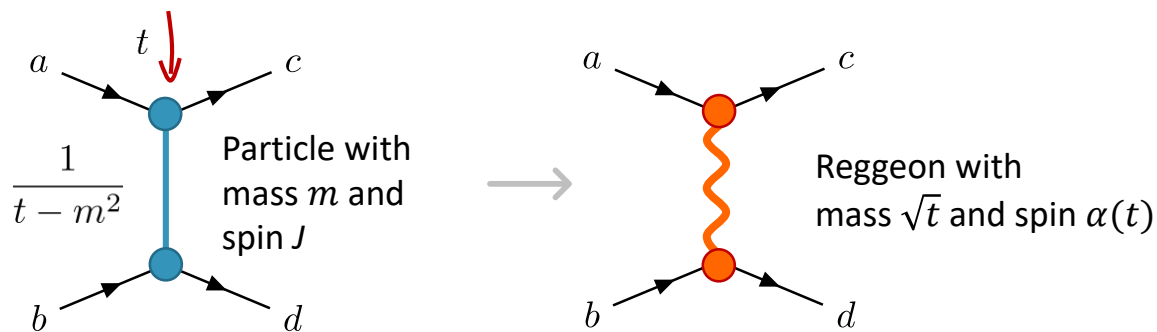
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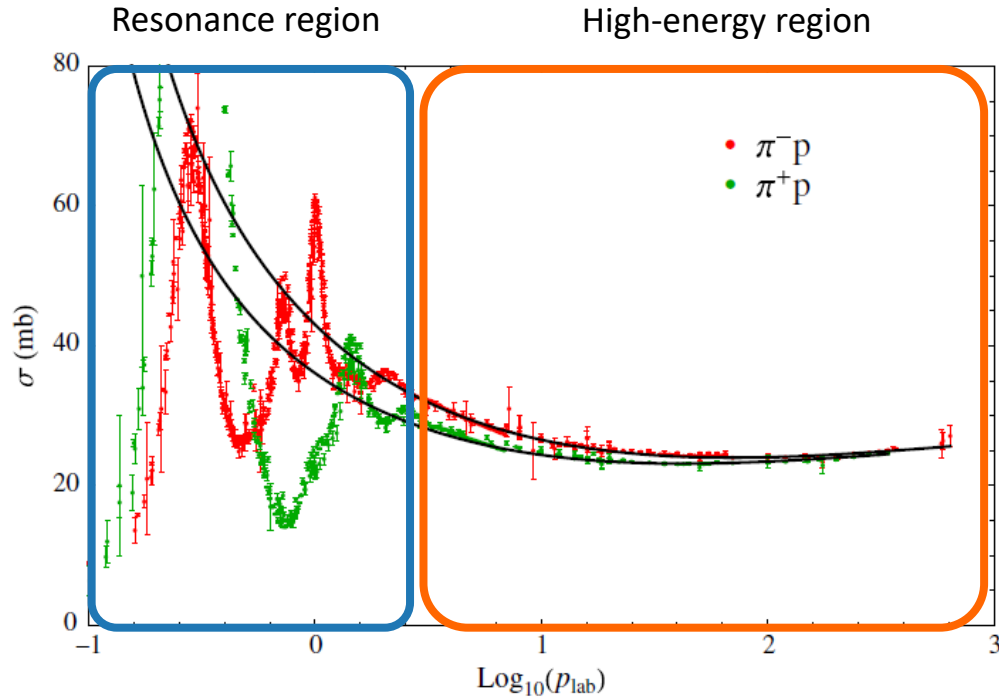
[V.Mathieu et al., *Phys.Rev.D* 98 (2018) 1, 014041]

$$\mathcal{P}_{\text{Regge}} = \frac{\pi\alpha'}{\sin(\pi\alpha(t))} \frac{1 + \eta e^{-i\pi\alpha(t)}}{2} \left(\frac{s}{s_0}\right)^{\alpha(t)} \frac{1}{\Gamma(1 + \alpha(t))}$$

Annotations for the equation components:

- $\frac{\pi\alpha'}{\sin(\pi\alpha(t))}$: poles for integer $\alpha(t)$
- $\frac{1 + \eta e^{-i\pi\alpha(t)}}{2}$: signature factor
- $\left(\frac{s}{s_0}\right)^{\alpha(t)}$: asymptotic behavior
- $\frac{1}{\Gamma(1 + \alpha(t))}$: cancel non-physical poles

When do we use Regge theory?



[V.Mathieu et al., *Phys.Rev.D* 92 (2015) 7]

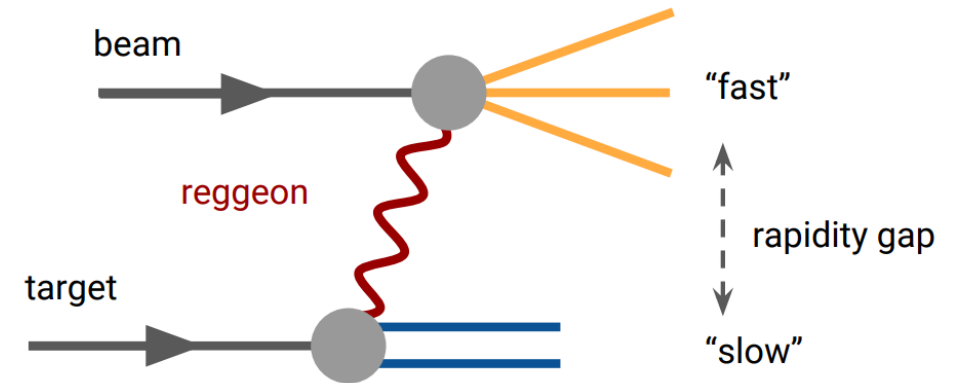
Low-energy: resonances

$$A(s, t) \sim \sum_r \frac{g_r}{s - s_r}$$

High energy: Regge exchanges

$$A(s, t) \sim \sum_i \frac{g_i}{t - t_i}$$

Analytically connected



Model for diffractive production

- We only need to know the (dominant) crossed-channel exchanges

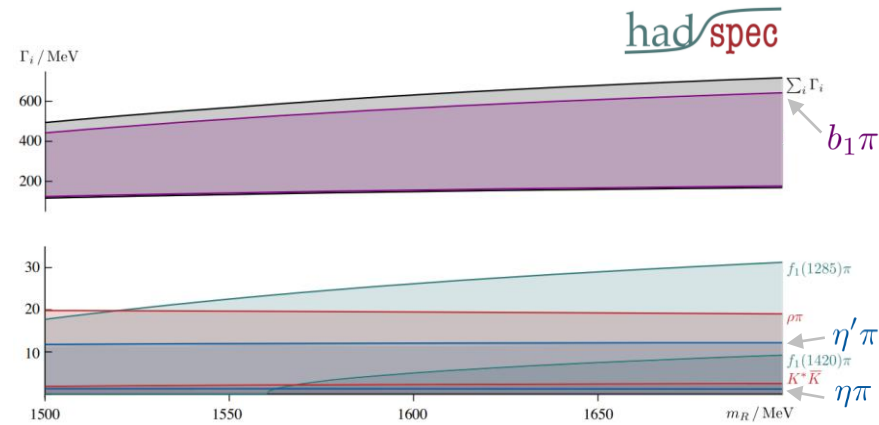
Photoproduction of the hybrid $\pi_1(1600)$

6

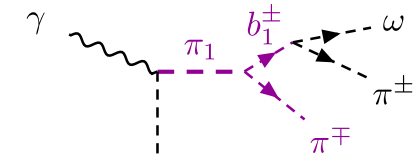
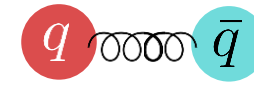
Lightest isoscalar with $J^{PC} = 1^{-+}$

Predicted by lattice QCD

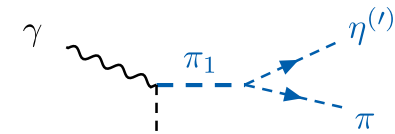
Experimental evidence from COMPASS)



[Woss et al., *Phys.Rev.D* 103 5, 054502 (2021)]

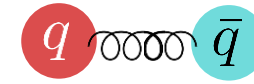


- ✓ More statistics
- ✗ Complicated final state



- ✗ Less statistics
- ✓ Easier (less particles)

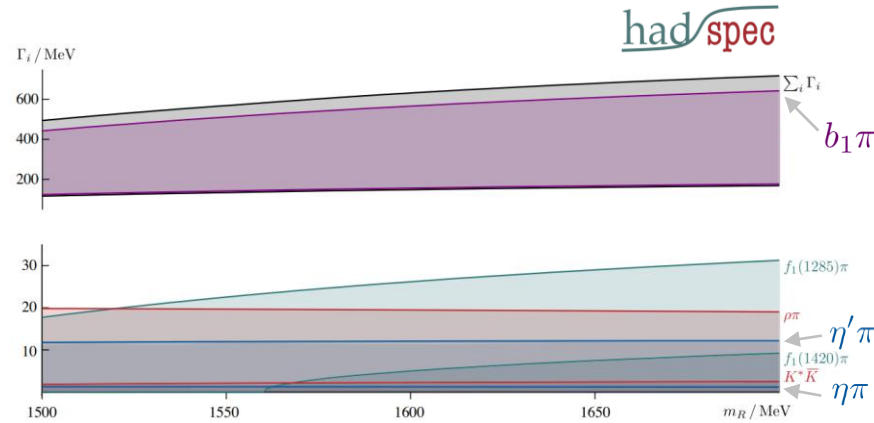
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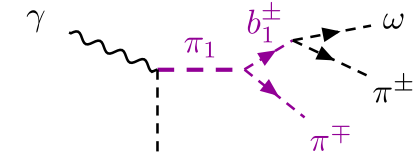
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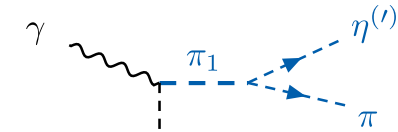


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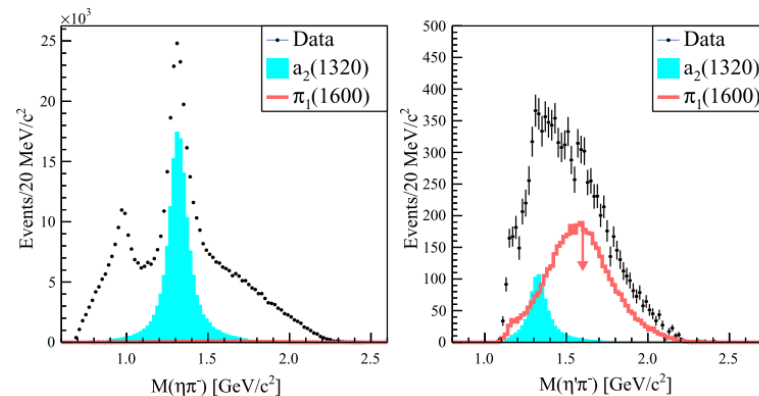
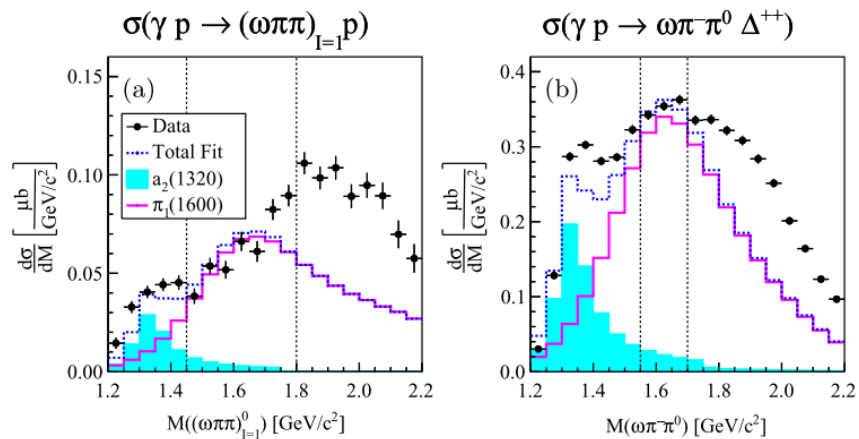
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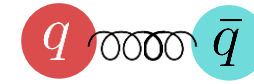
[Afzal et al., *Phys.Rev. Lett.* 133 261903 (2024)]



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Need a better understanding of production mechanisms!

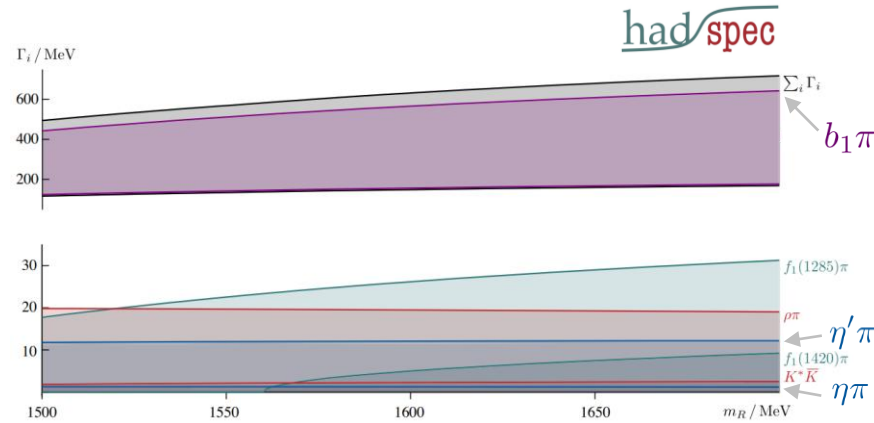
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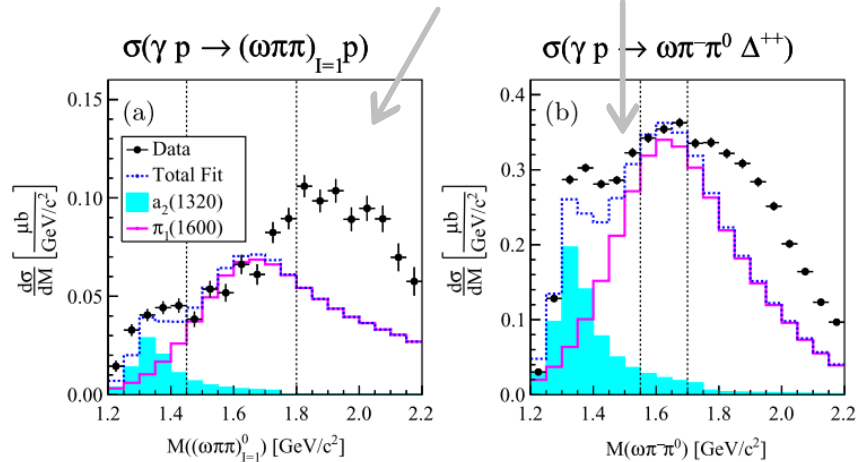
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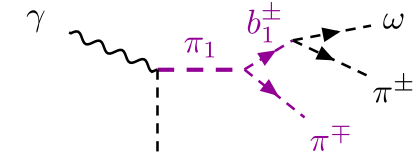
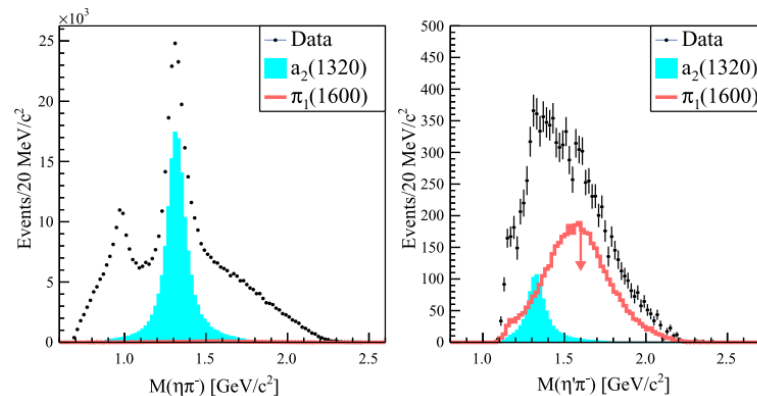
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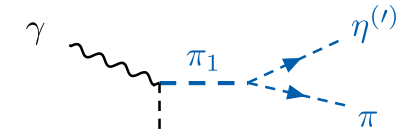
(neutral)
pomeron < pion
exchange exchange



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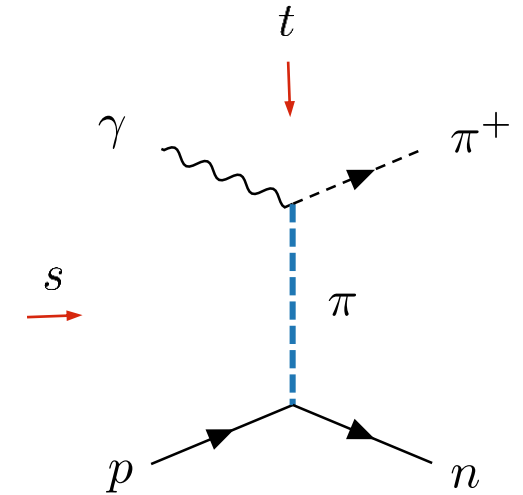
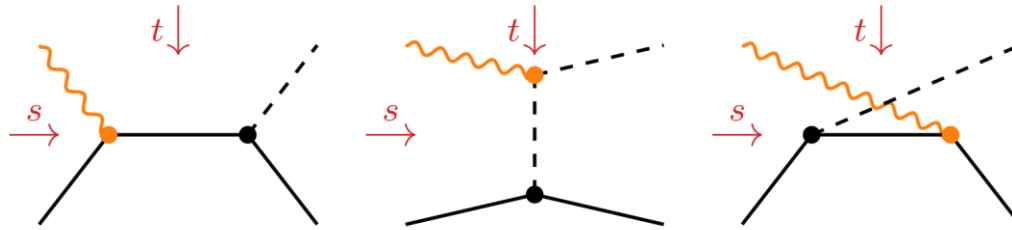
Understanding pion exchange

[GM, Daniel Winney, et al. (JPAC), *Phys.Rev.D* 110 (2024) 11, 114012]

Simplest charge-exchange process is pion photoproduction

Pion exchange diagram vanishes in the t -channel CM frame

In terms of Born diagrams, **gauge invariance** relates t -channel diagram (pion exchange) and s -, u -channel diagrams (nucleon exchanges)



Understanding pion exchange

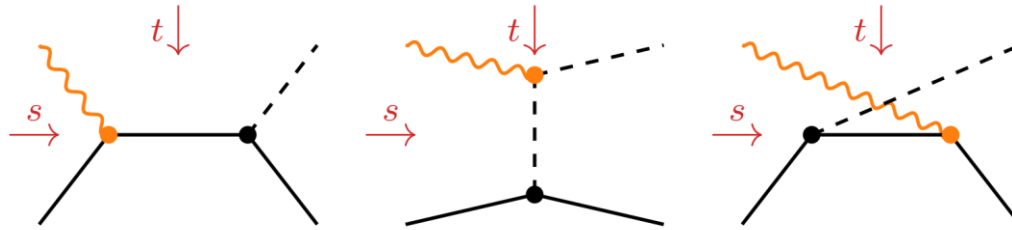
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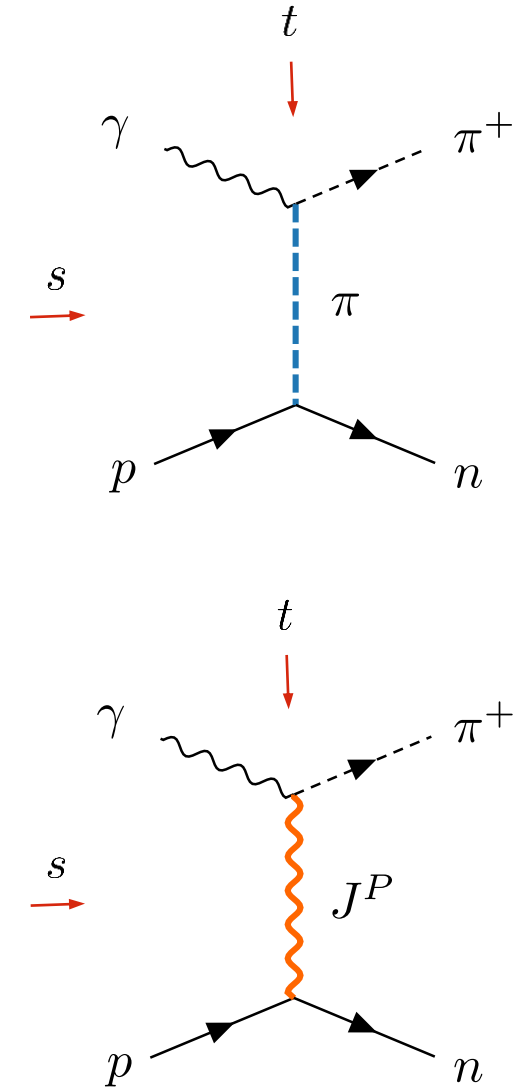
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What about the Reggeized pion?

$$A_{\lambda_\gamma \lambda_i \lambda_f}(s, t) = \sum_J (2J + 1) a_{\lambda_\gamma \lambda_i \lambda_f}^J(t) \underbrace{d_{\lambda_\gamma \lambda_i - \lambda_f}^J(\theta_t)}_{J \geq |\lambda_\gamma| = 1}$$

No pion pole?

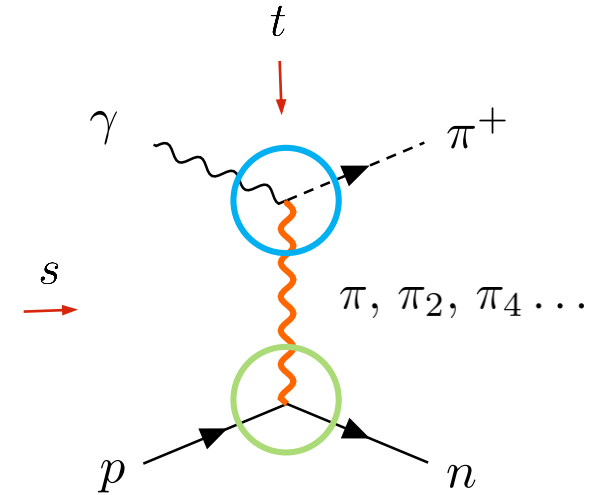


Reggeization of pion exchange

[GM, Daniel Winney, et al. (JPAC), *Phys.Rev.D* 110 (2024) 11, 114012]

1. Build an amplitude for the exchange of a particle of arbitrary spin $J > 0$ (gauge invariant by construction)

$$A_{\lambda_\gamma \lambda_i \lambda_f}^J(s, t) = \sum_{\sigma_J} \frac{V_{\lambda_\gamma}^J(\sigma_J) V_{\lambda_i \lambda_f}^J(\sigma_J)}{J - \alpha(t)} = \underbrace{a_{\lambda_\gamma \lambda_i \lambda_f}^J(t)}_{\propto \frac{1}{\sqrt{J}}} d_{\lambda_\gamma \lambda_i - \lambda_f}^J(\theta_t)$$



Reggeization of pion exchange

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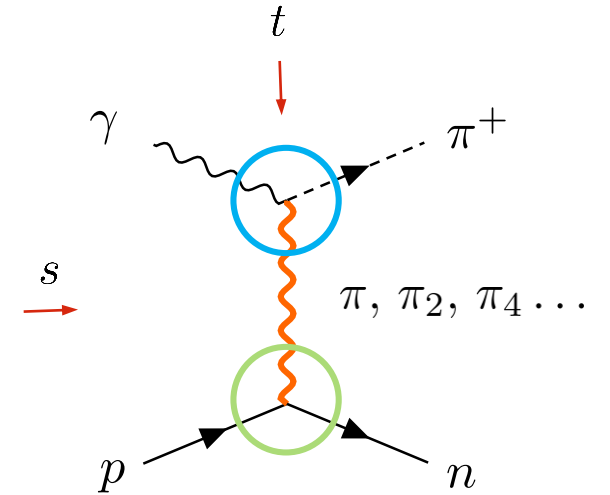
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2. Analytically continue to $J = 0$
 - Use the hypergeometric function

$$d_{\lambda_\gamma, 0}^J(\theta_t) \propto \sqrt{J}$$

$$A_{\lambda_\gamma \lambda_i \lambda_f}^{J \rightarrow 0}(s, t) \propto \frac{t}{-\alpha(t)} (2\lambda_i \lambda_\gamma \delta_{\lambda_i \lambda_f}) \frac{z_t}{\sqrt{1 - z_t}} \approx i \frac{t}{m_\pi^2 - t} (2\lambda_i \lambda_\gamma \delta_{\lambda_i \lambda_f})$$

The spin $J \rightarrow 0$ amplitude is finite!



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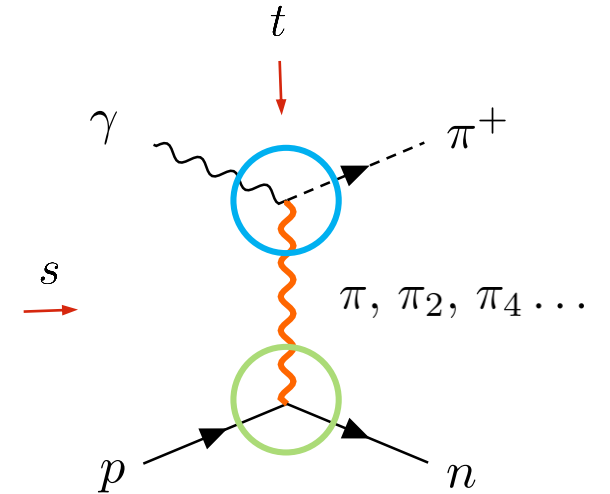
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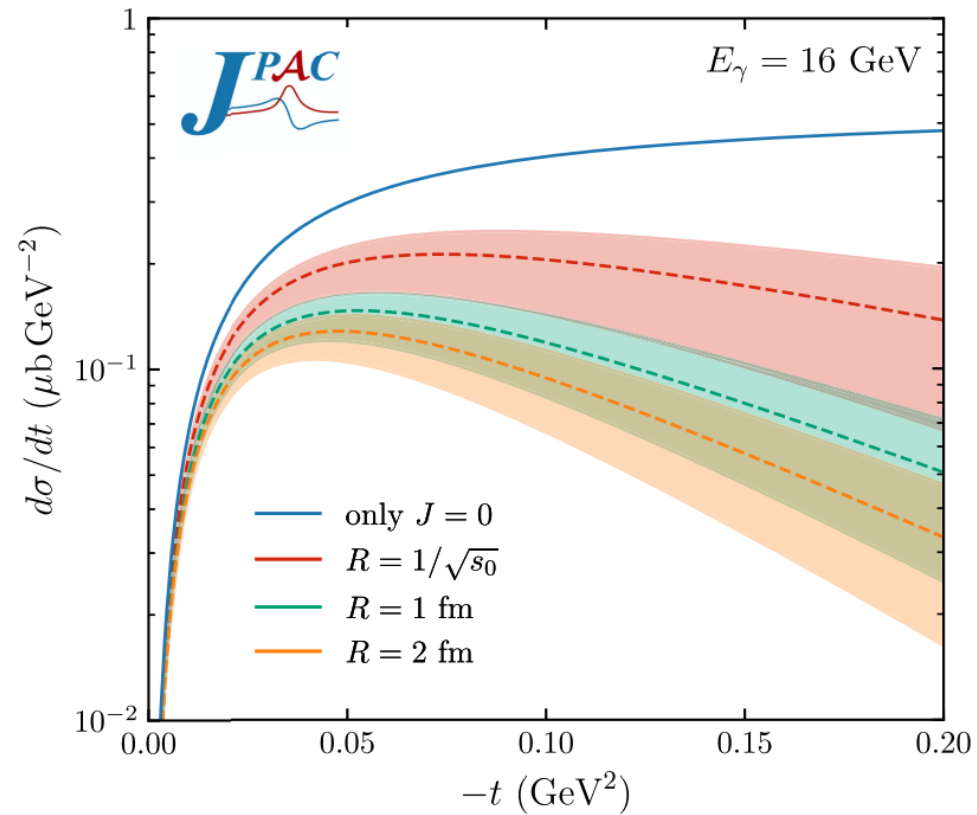
3. Sum the tower of exchanges (e.g. Sommerfeld-Watson transform, or using the generating function for the Jacobi polynomials)

$$A_{\lambda_\gamma \lambda_i \lambda_f}(s, t) = A_{\lambda_\gamma \lambda_i \lambda_f}^{J \rightarrow 0}(s, t) + \sum_{J=1} (2J + 1) a_{\lambda_\gamma \lambda_i \lambda_f}^J(t) d_{\lambda_\gamma \lambda_i - \lambda_f}^J(\theta_t) \propto t(sR^2)^{\alpha(t)}$$



Reggeized pion exchange

$$A_{\lambda_\gamma \lambda_i \lambda_f}^{\text{Regge}}(s, t) \propto t(sR^2)^{\alpha(t)}$$



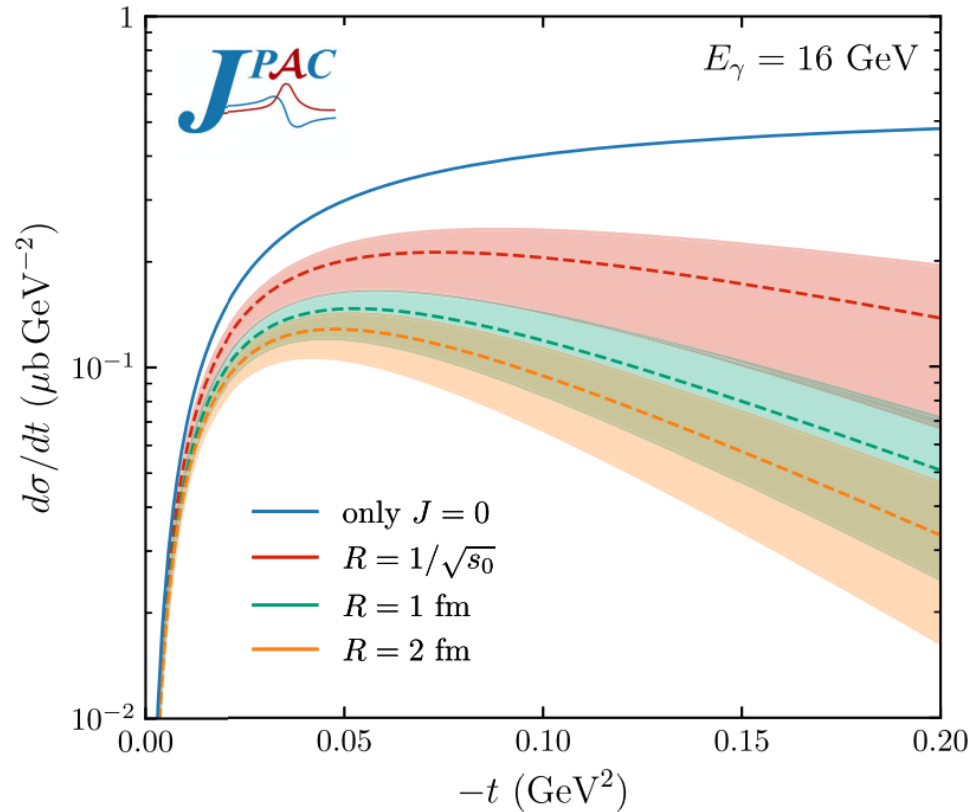
Results for pion photoproduction

[GM, Daniel Winney, et al. (JPAC), *Phys.Rev.D* 110 (2024) 11, 114012]

9

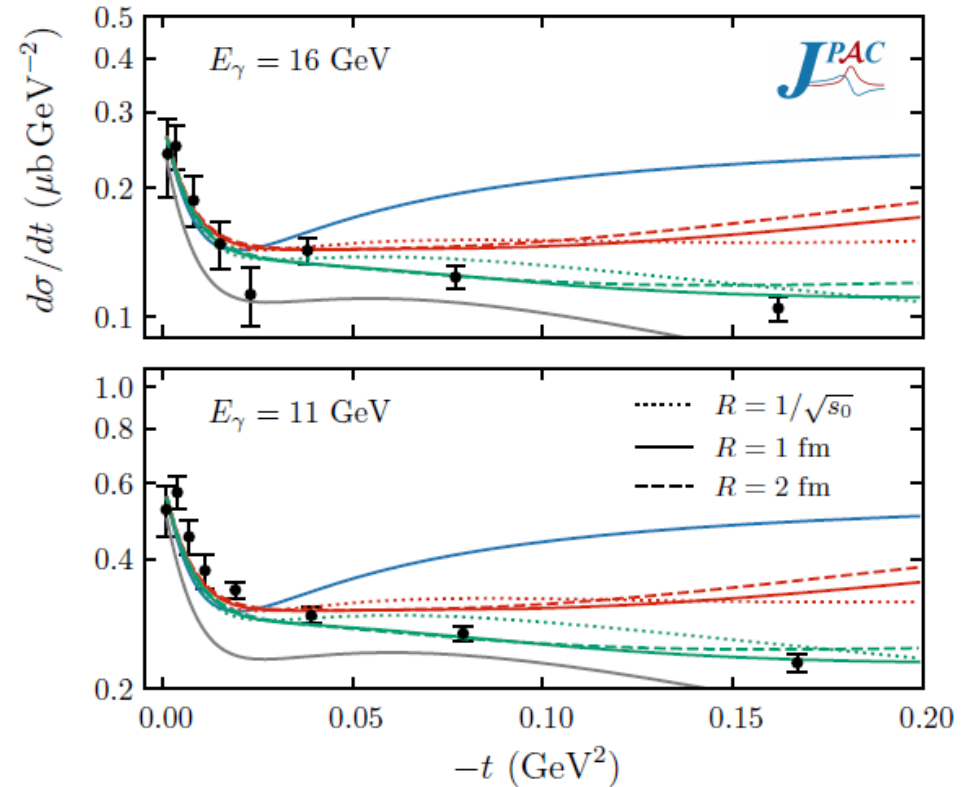
Reggeized pion exchange

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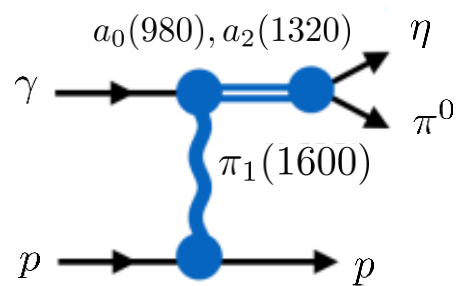
Adding the nucleon magnetic term

- Gauge invariant by itself

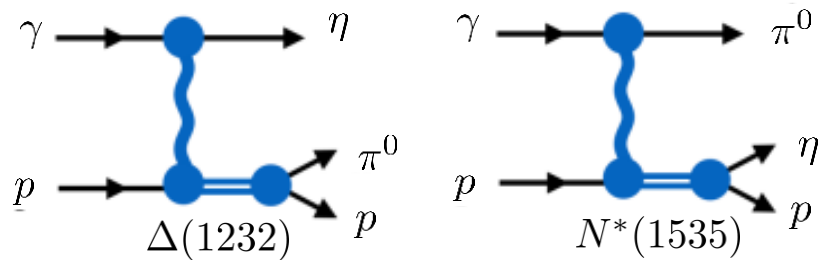


Photoproduction of $\eta\pi$ and $\eta'\pi$

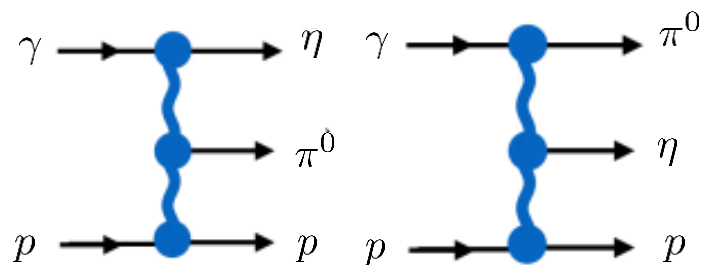
Meson resonances



Baryon resonances

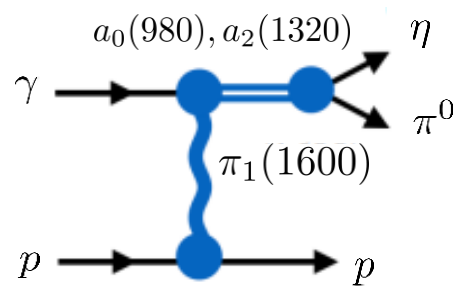


Double Regge production

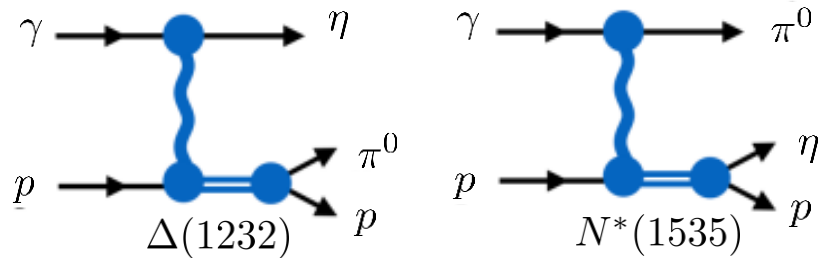


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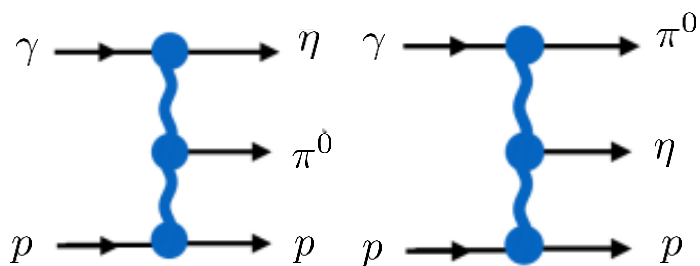
Meson resonances



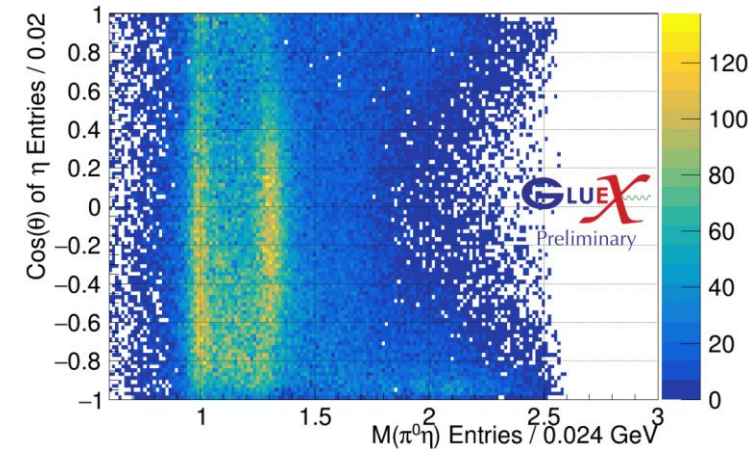
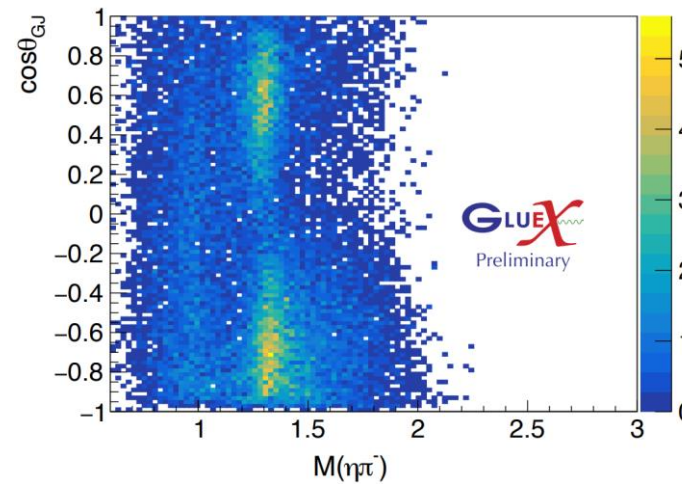
Baryon resonances



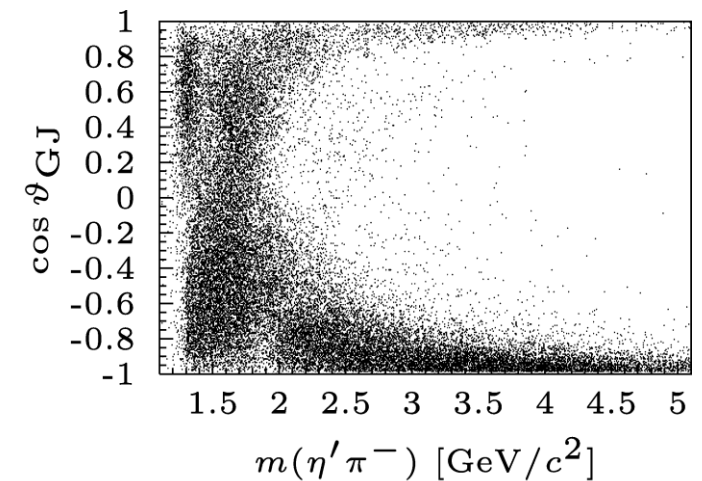
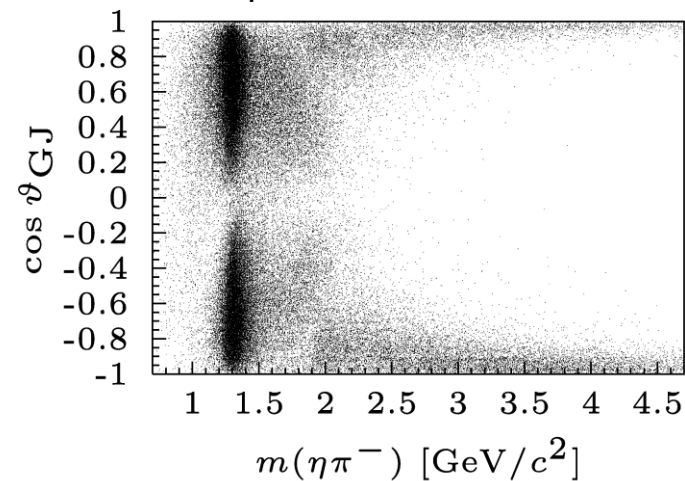
Double Regge production



With photon beam



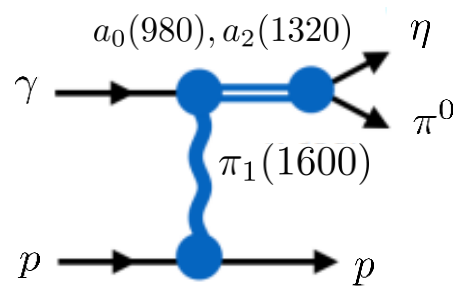
With pion beam



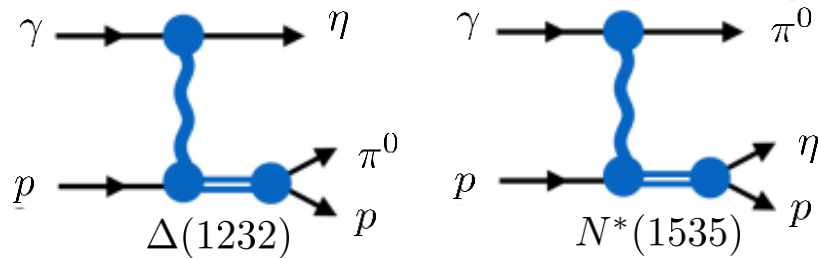
Photoproduction of $\eta\pi$ and $\eta'\pi$

- Clear signals of the $a_0(980)$ and $a_2(1320)$

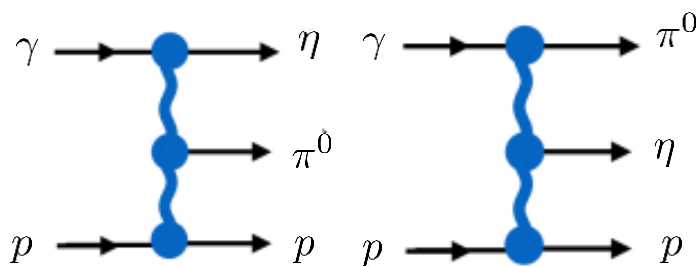
Meson resonances



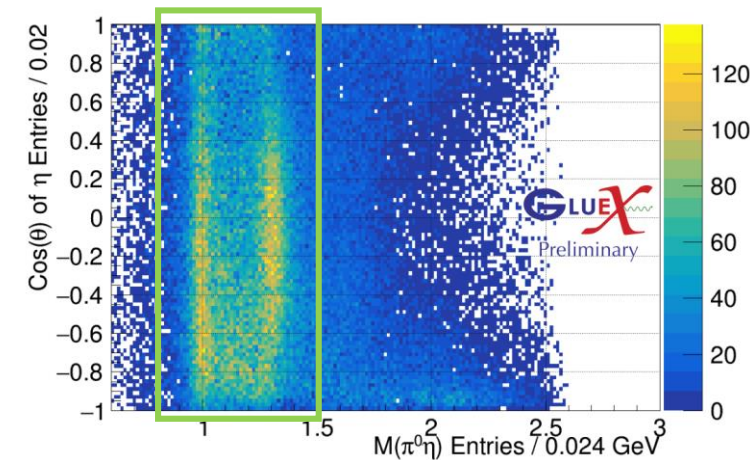
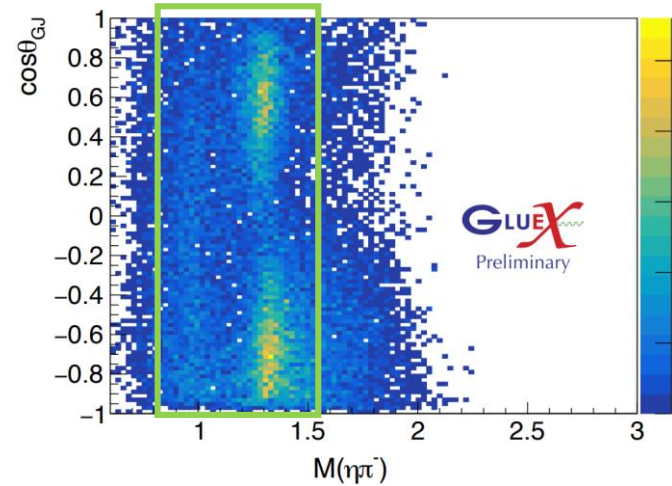
Baryon resonances



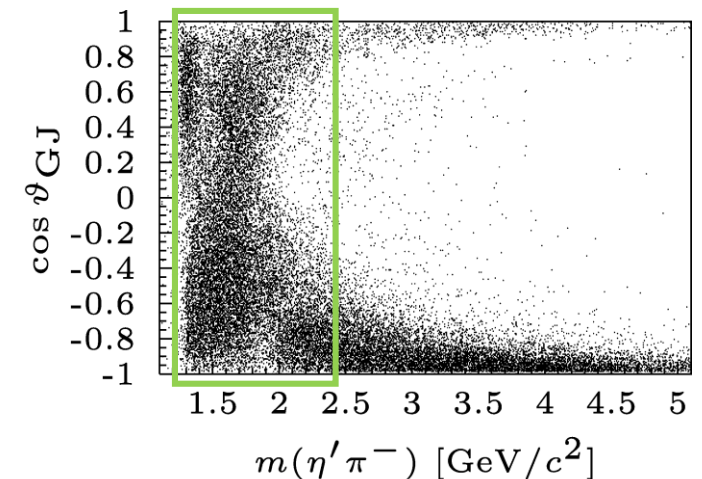
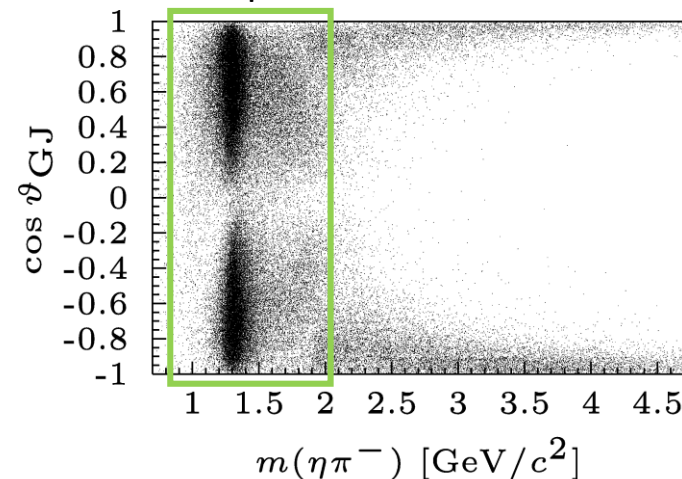
Double Regge production



With photon beam



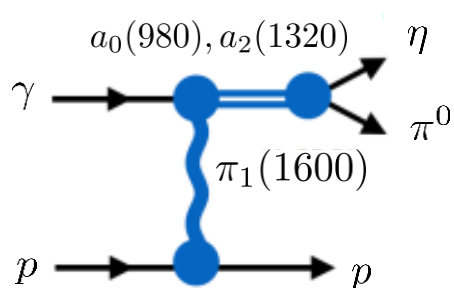
With pion beam



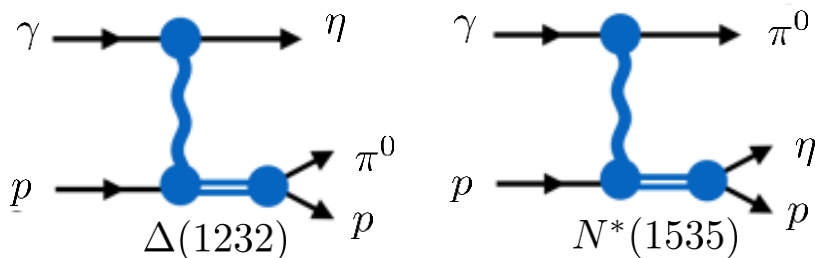
Photoproduction of $\eta\pi$ and $\eta'\pi$

- Clear signals of the $a_0(980)$ and $a_2(1320)$
- Forward-backward asymmetry at high energies

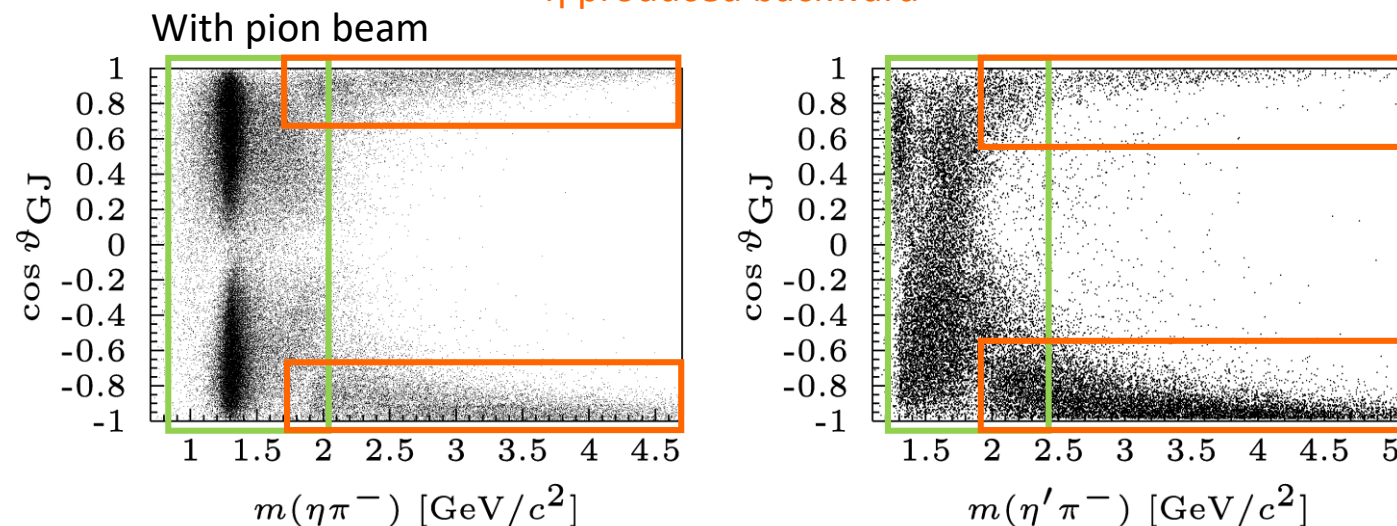
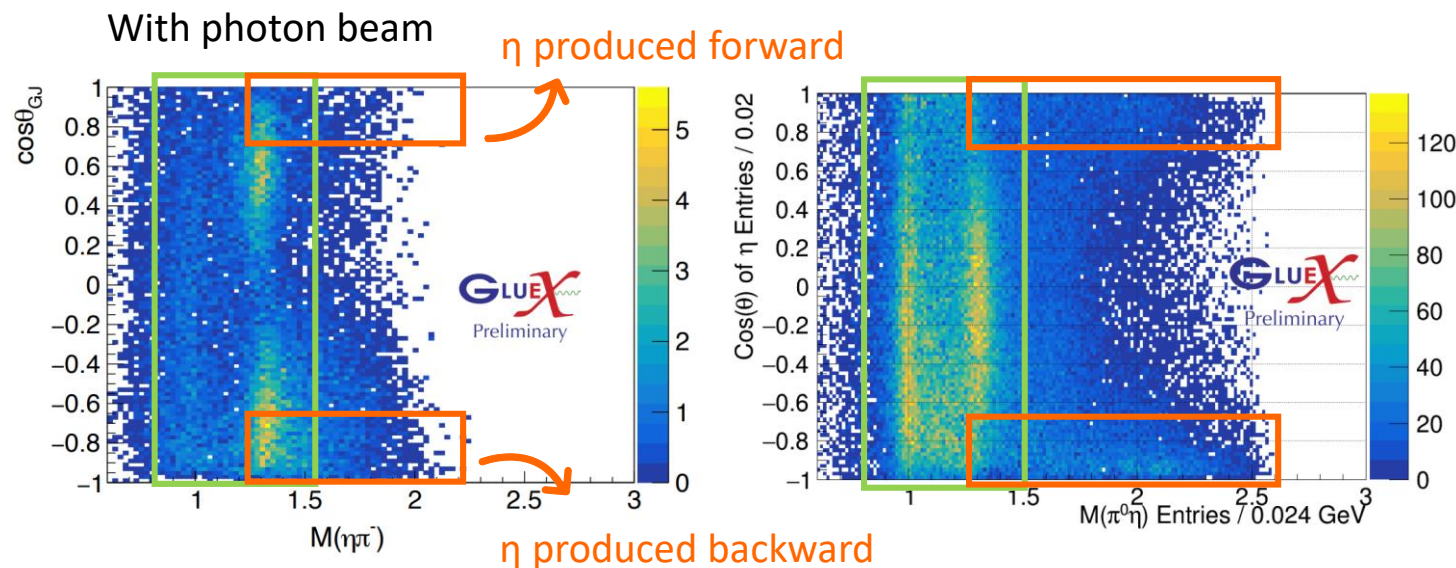
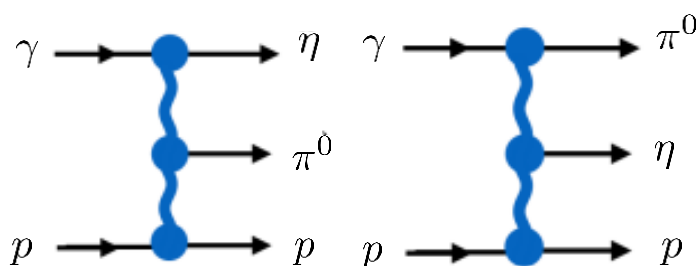
Meson resonances



Baryon resonances



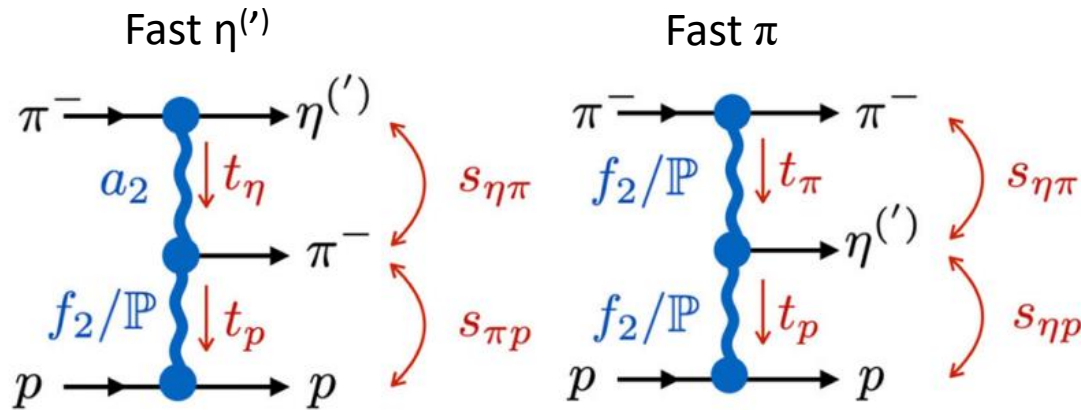
Double Regge production



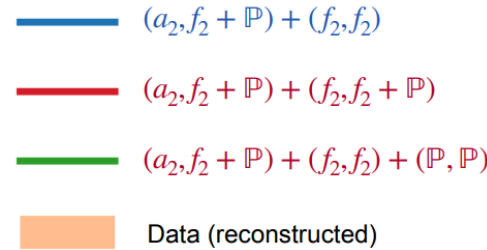
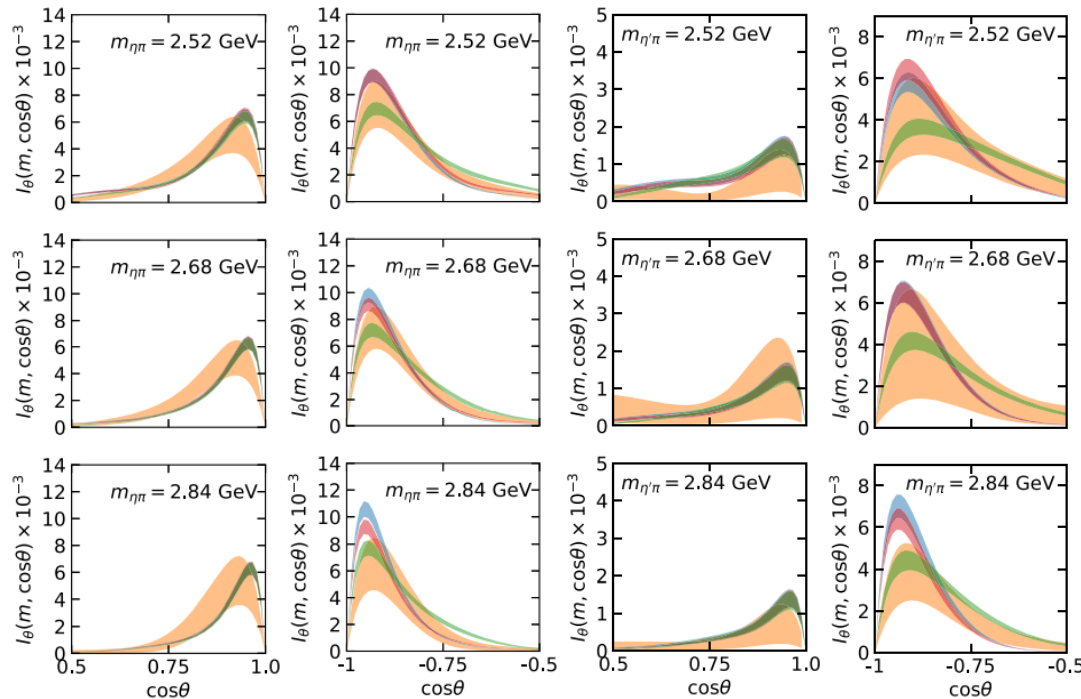
Double Regge $\eta^{(\prime)}\pi$ at COMPASS

11

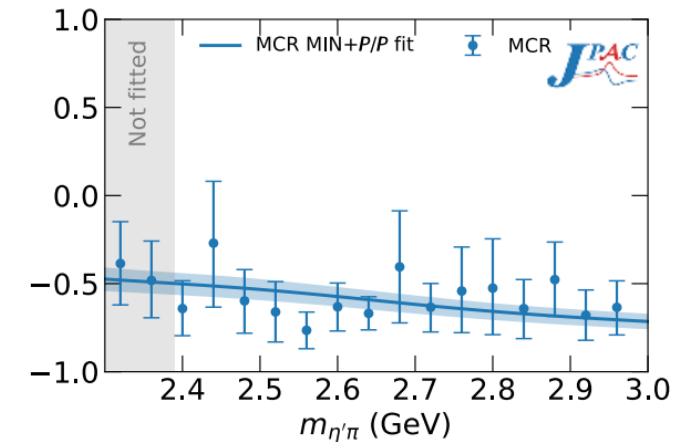
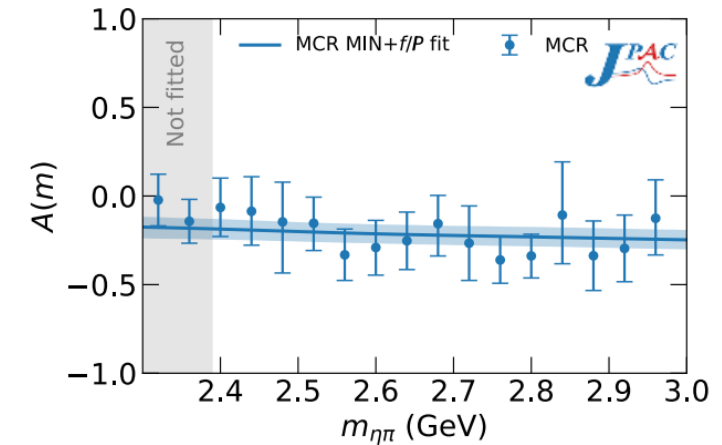
[L. Bibrzycki, C. Fernandez-Ramirez et al. (JPAC), *Eur.Phys.J.C* 81 (2021) 647]



$$\text{Asymmetry} = \frac{\text{Forward} - \text{Backward}}{\text{Forward} + \text{Backward}}$$



Asymmetry originating mainly
from $(a_2, f_2/\mathbb{P}) \neq (f_2, f_2/\mathbb{P})$
and from $(\mathbb{P}, \mathbb{P})$ in $\eta'\pi$



Double Regge model $\eta^{(\prime)}\pi$ photoproduction

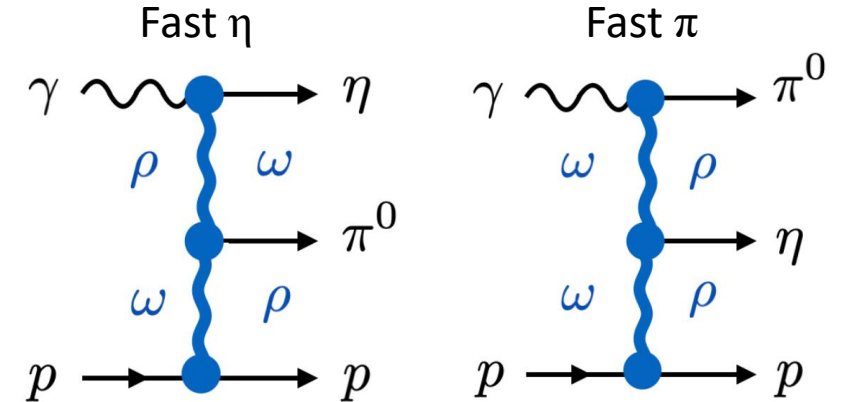
[GM, Vincent Mathieu, et al. (JPAC), arXiv:2008.XXXX]

1. Start from a vector-vector exchange model

$$A_{\lambda_\gamma \lambda \lambda'} = \sum_{\text{diagrams}} \left(\frac{g_{\gamma V_1 P_1}}{m_0} \frac{g_{V_1 V_2 P_2}}{m_0} \right) \frac{g_{\lambda_\gamma, \lambda' - \lambda} K_{\lambda_\gamma \lambda \lambda'}}{(m_{V_1}^2 - im_{V_1} \Gamma_{V_1} - t_1)(m_{V_2}^2 - im_{V_2} \Gamma_{V_2} - t_p)}$$

2. Reggeize

$$K_{\lambda_\gamma \lambda \lambda'} \rightarrow K_{\lambda_\gamma \lambda \lambda'} \times R \left(\alpha_1(t_\pi), \alpha_2(t_p), s_{\eta\pi}, s_{\eta p}, \frac{s}{\alpha' s_{\eta\pi} s_{\eta p}} \right)$$

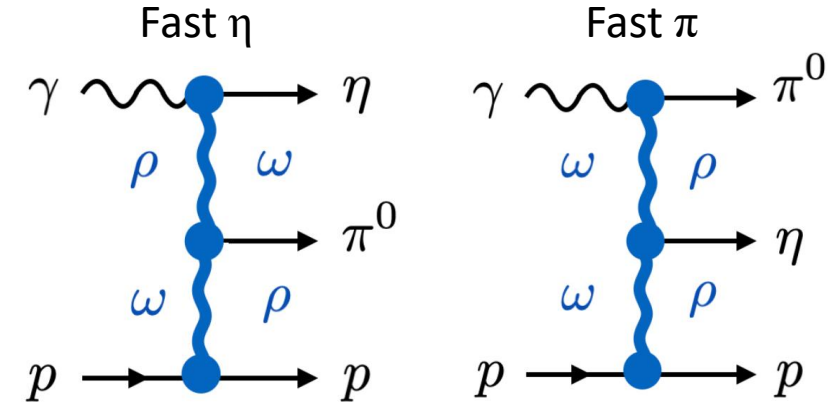


Double Regge model $\eta^{(\prime)}\pi$ photoproduction

[GM, Vincent Mathieu, et al. (JPAC), arXiv:2008.XXXX]

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[Shimada, Martin and Irving, *Nucl. Phys. B* 142 (1978)]

$$R(\alpha_1, \alpha_2, s_1, s_2, \eta) = \boxed{\xi_1 \xi_{21}} \eta^{\alpha_1 - 1} \boxed{V(\alpha_1, \alpha_2, \eta)} + \boxed{\xi_2 \xi_{12}} \eta^{\alpha_2 - 1} \boxed{V(\alpha_2, \alpha_1, \eta)} (\alpha' s_1)^{\alpha_1 - 1} (\alpha' s_2)^{\alpha_2 - 1} [(\alpha_1 - 1)\Gamma(1 - \alpha_1)] [(\alpha_2 - 1)\Gamma(1 - \alpha_2)]$$

vertex function

signature factors

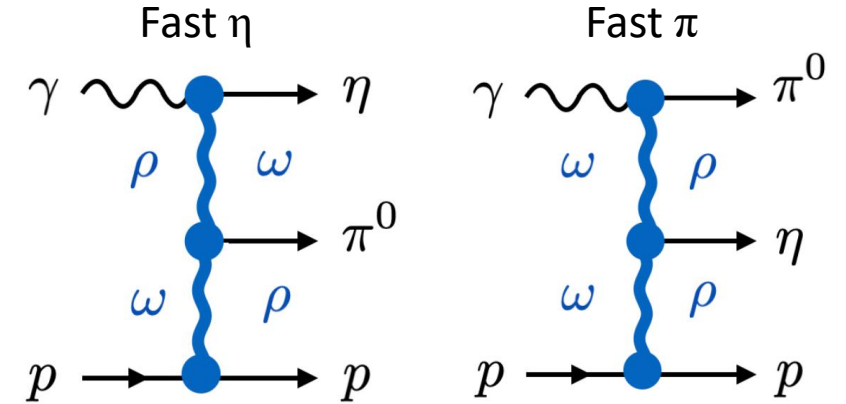
$$\eta = \frac{s}{\alpha' s_1 s_2}$$

Double Regge model $\eta^{(\prime)}\pi$ photoproduction

[GM, Vincent Mathieu, et al. (JPAC), arXiv:2008.XXXX]

1. Start from a vector-vector exchange model

$$A_{\lambda_\gamma \lambda \lambda'} = \sum_{\text{diagrams}} \left(\frac{g_{\gamma V_1 P_1}}{m_0} \frac{g_{V_1 V_2 P_2}}{m_0} \right) \frac{g_{\lambda_\gamma, \lambda' - \lambda} K_{\lambda_\gamma \lambda \lambda'}}{(m_{V_1}^2 - im_{V_1} \Gamma_{V_1} - t_1)(m_{V_2}^2 - im_{V_2} \Gamma_{V_2} - t_p)}$$



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vertex function

signature factors

$$\eta = \frac{s}{\alpha' s_1 s_2}$$

In the double Regge limit ($s, s_{\eta\pi}, s_p \rightarrow \infty, \eta \sim \text{ct.}, s_{\eta\pi}/s_{\eta p} \sim \text{ct.}$)

$$K_{\lambda_\gamma \lambda \lambda'} \sim s_{\eta\pi} s_{\eta p} \quad \text{and} \quad K_{\lambda_\gamma \lambda \lambda'} R \sim C_1 s^{\alpha_1(t_\pi)} s_{\eta p}^{\alpha_2(t_p) - \alpha_1(t_\pi)} + C_2 s_{\eta\pi}^{\alpha_1(t_\pi) - \alpha_2(t_p)} s^{\alpha_2(t_p)}$$

Forward-Backward asymmetry

(Preliminary results)

The asymmetry is defined as ($M = \sqrt{s_{\eta\pi}}$)

$$A(M) = \frac{F(M) - B(M)}{F(M) + B(M)}$$

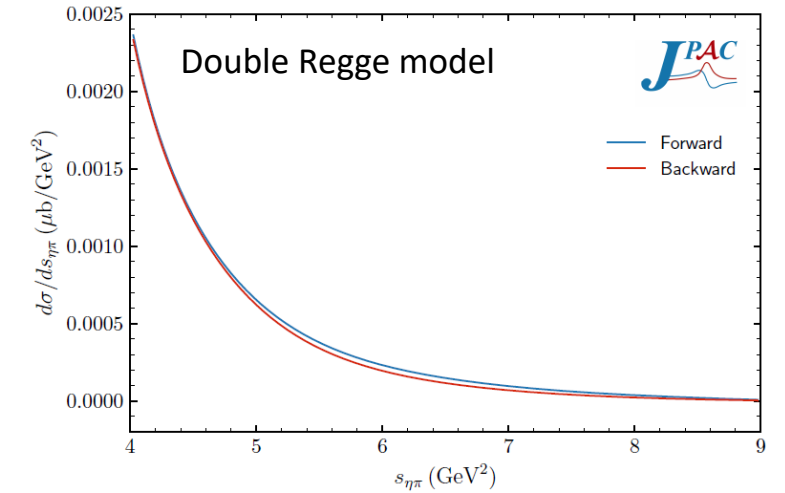
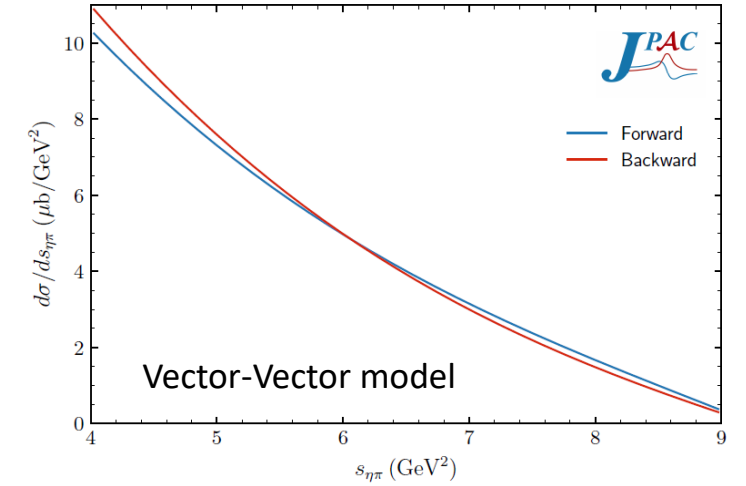
The forward and backward intensities contain cuts in kinematic variables

$$F(M) = \int d^2\Omega_{GJ} dz_{cm} \frac{d\sigma}{d\Omega_{GJ} dz_{cm} dM_{\eta\pi}^2} \theta(t_\eta > -2) \theta(t_p > -2) \theta(s_{\eta\pi} > 4) \theta(s_{\pi p} > 4) \theta(s_{\eta p} > 4.75)$$

$$B(M) = \int d^2\Omega_{GJ} dz_{cm} \frac{d\sigma}{d\Omega_{GJ} dz_{cm} dM_{\eta\pi}^2} \theta(t_\pi > -2) \theta(t_p > -2) \theta(s_{\eta\pi} > 4) \theta(s_{\pi p} > 4) \theta(s_{\eta p} > 4.75)$$

and the physics model is captured in $\frac{d\sigma}{d\Omega_{GJ} dz_{cm} ds_{\eta\pi}} \propto |A_{\lambda_\gamma \lambda \lambda'}|^2$

[GM, Vincent Mathieu, et al. (JPAC), arXiv:2008.XXXX]



Forward-Backward asymmetry

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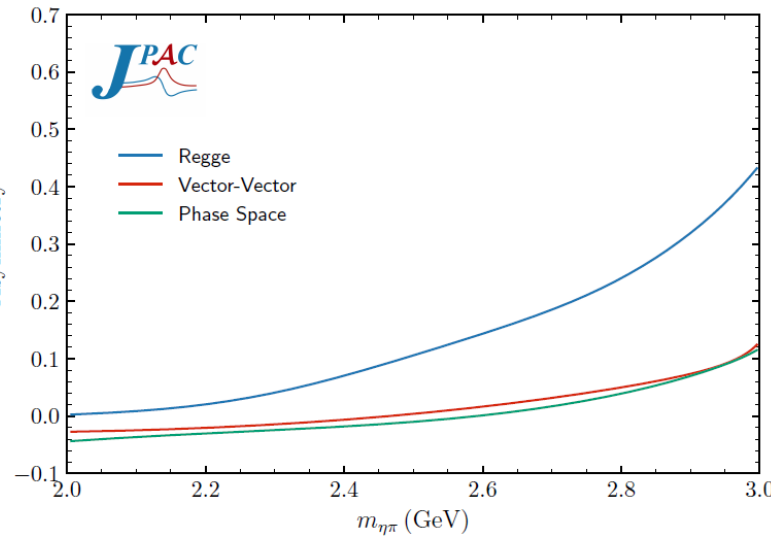
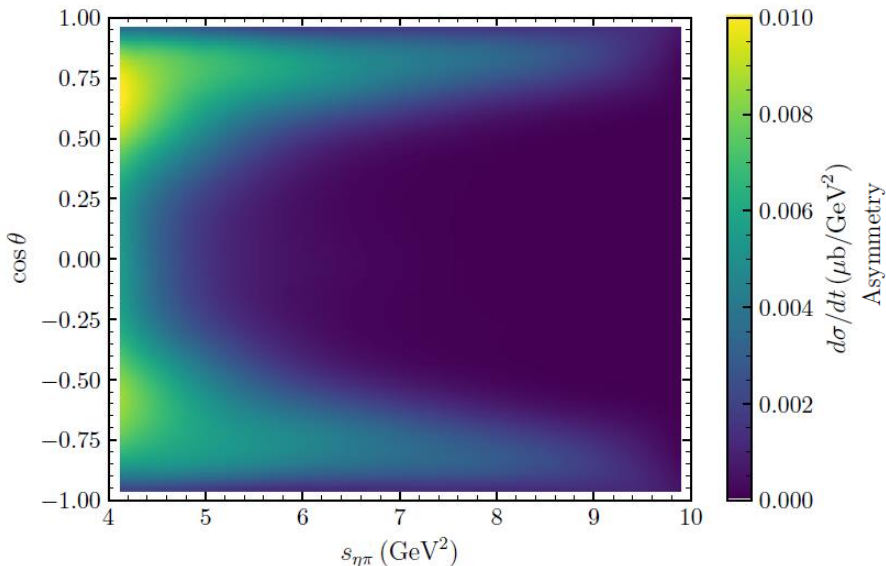
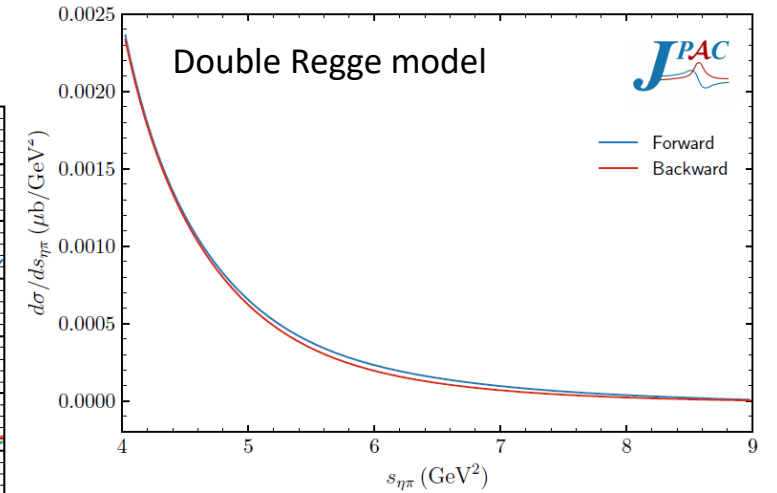
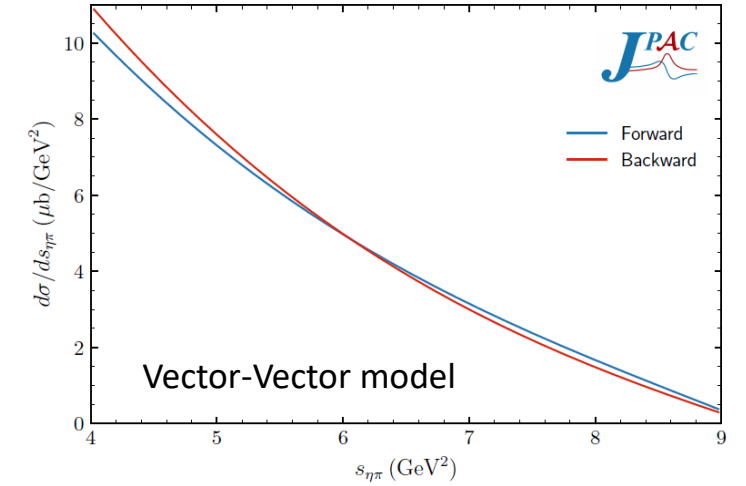
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[GM, Vincent Mathieu, et al. (JPAC), arXiv:2008.XXXX]



Summary

Regge theory is a practical tool to model and understand meson production at **high energies**

- Step forward in understanding pion exchange in pion photoproduction (gauge invariance)
- Interesting physics in $\eta^{(\prime)}\pi$, in the resonance region (hybrid meson) but also in the double Regge region

Strong effort from JPAC to **provide theory support to experimental analyses**
(GlueX, but also CLAS, COMPASS...)

- Understand the data
- Provide predictions
- Close collaboration with experimentalists