

Model simulations for BES-II

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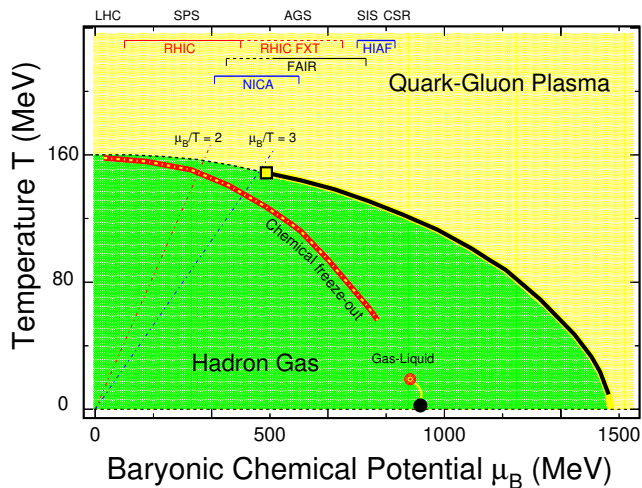
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Simulations = dynamical simulations
= hydro or transport (hadronic cascades)

The goal

Heavy-ion collisions \Leftrightarrow hydro simulations



\Leftrightarrow

Possible signals of the **1st order phase transition** in:

- Shape of net proton distribution
- Flow
- HBT
- ...

Model construction

Hydro-based models for RHIC BES region include: UrQMD+vHLLLE (2015), SMASH+vHLLLE (2022), EPOS4 2024), 3D Glauber+MUSIC (2018?), SMASH+MUSIC (a.k.a. JETSCAPE, 2025), CCAKE+SMASH (a.k.a. MUSES, 2025)

The recipe is mostly imported from higher energies:

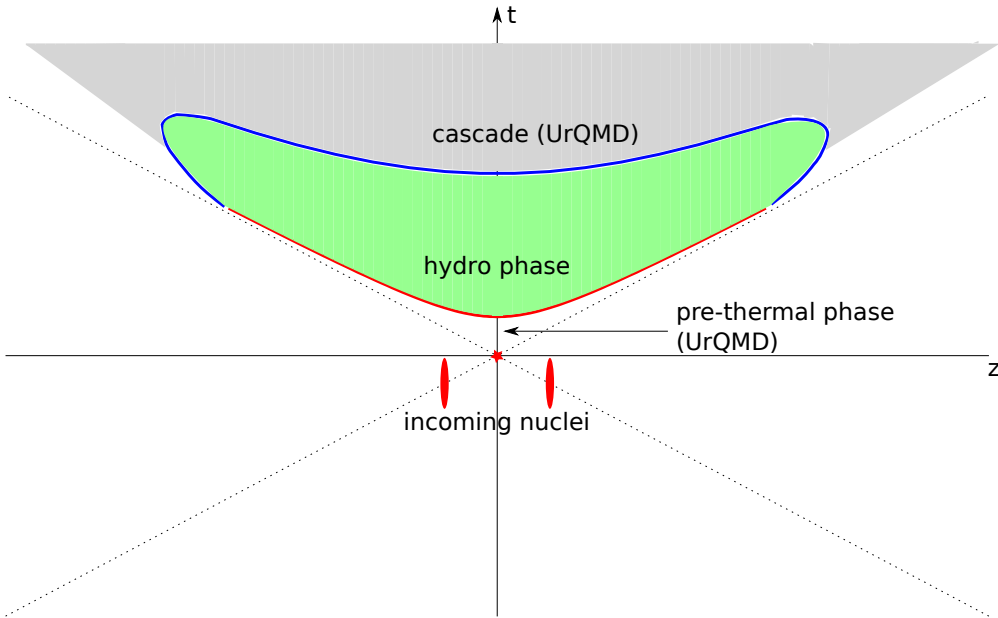
- 3D initial state: from transport or parametrized
Start hydro later (nuclei are not thin pancakes)
- 3D hydro with finite viscosity (not a challenge)
and conserved charges (not a challenge either)
- EoS at finite μ_B – an input to the modelling
- Same or slightly adapted Cooper-Frye
- final-state hadronic cascade

Most of basic observables are reproduced (see next slides)

Modelling is complex

- multi-step simulation, e.g.
(SMASH) + (vHLLLE) +
(smash-hadron-sampler) +
(SMASH)
- 3 to 4 different codes
stacked one after another
- Need to take care of every
interface

The first interface: initial state to hydro



The easy choice: start fluid-dynamic picture at $\tau = \tau_0$

Motivated at high collision energies:

- Longitudinal boost invariance
- Physics changes as a function of τ

Parametrized 3D initial state at $\tau=\tau_0$

superMC code (?)

$$e(x, y, \eta_s; y_{\text{CM}}) = \mathcal{N}_e(x, y) \times \exp \left[-\frac{(|\eta_s - y_{\text{CM}}| - \eta_0)^2}{2\sigma_\eta^2} \theta(|\eta_s - y_{\text{CM}}| - \eta_0) \right]$$

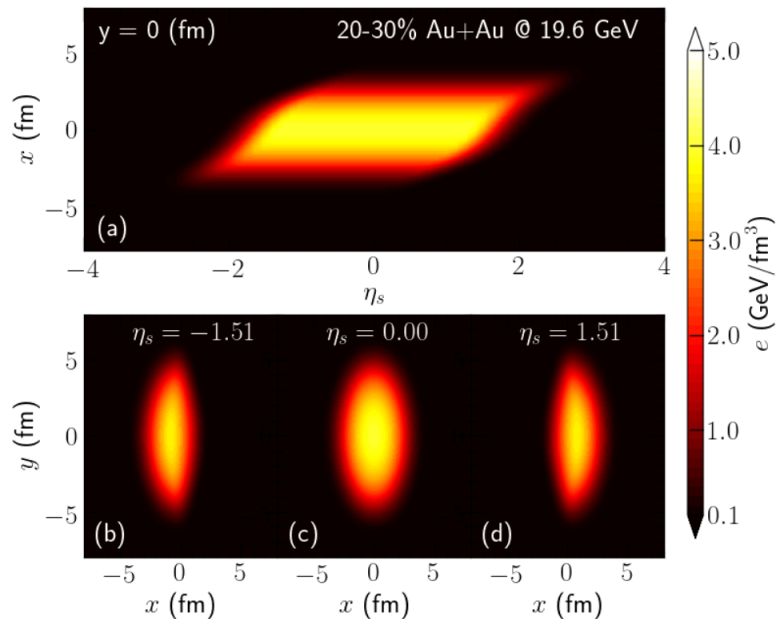
$\mathcal{N}_e(x, y)$ determined from the full energy count

Parameter tuning:

$\sqrt{s_{\text{NN}}} \text{ (GeV)}$	$\tau_0 \text{ (fm/c)}$	η_0	σ_η	$\eta_{B,0}$	$\sigma_{B,\text{in}}$	$\sigma_{B,\text{out}}$
AuAu & dAu @ 200	1.0	2.5	0.6	3.5	2.0	0.1
AuAu & dAu @ 62.4	1.0	2.25	0.3	2.7	1.9	0.2
AuAu & dAu @ 39	1.3	1.9	0.3	2.2	1.6	0.2
AuAu@27	1.4	1.6	0.3	1.8	1.5	0.2
AuAu & dAu @ 19.6	1.8	1.3	0.3	1.5	1.2	0.2
AuAu@14.5	2.2	1.15	0.3	1.4	1.15	0.2
AuAu@7.7	3.6	0.9	0.2	1.05	1.0	0.1
PbPb@17.3	1.8	1.25	0.3	1.6	1.2	0.2
PbPb@8.77	3.5	0.95	0.2	1.2	1.0	0.1

Shen, Alzhrani, [Phys.Rev.C 102 \(2020\) 1, 014909](#)

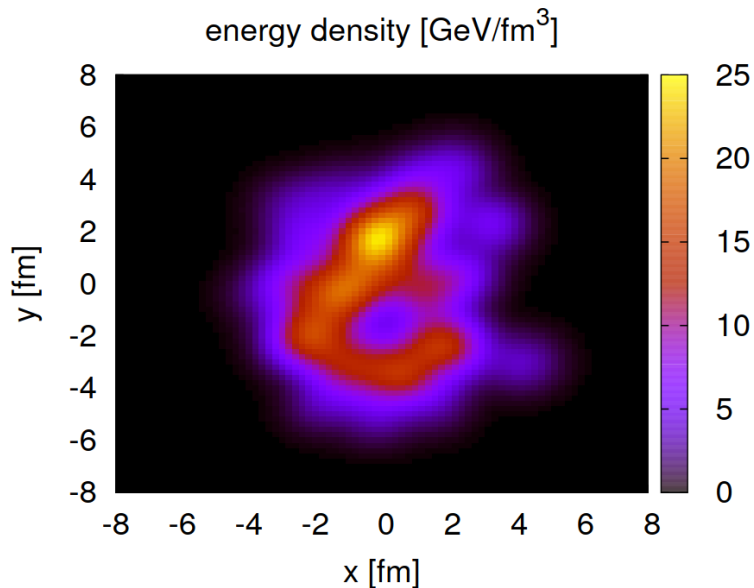
+ follow-up by Lipei du et al



3D initial state at $\tau=\tau_0$ from transport

...which hopefully has fewer parameters as initial state dynamics dictates the shape

(UrQMD, SMASH, EPOS, JAM)



- Patch a transport code to stop evolution or disable interactions after after $\tau=\tau_0$
- Collect hadrons at $\tau=\tau_0$
- Assume that hadrons melt into fluid:

$$\Delta P_{ijk}^\alpha = P^\alpha \cdot C \cdot \exp\left(-(\Delta x_i^2 + \Delta y_j^2)/R_\perp^2 - \Delta \eta_k^2 \gamma_\eta^2 \tau_0^2 / R_\eta^2\right)$$

$$\Delta N_{ijk}^0 = N^0 \cdot C \cdot \exp\left(-(\Delta x_i^2 + \Delta y_j^2)/R_\perp^2 - \Delta \eta_k^2 \gamma_\eta^2 \tau_0^2 / R_\eta^2\right)$$

$$E/\Delta V = (\epsilon + p)(u^0)^2 - p, \quad P^i/\Delta V = (\epsilon + p)u^0 u^i$$

- 3D initial state with finite n_B, n_Q, n_S

Does the medium look like a fluid when it is fluidized?

In case of event averaged initial state, in some cases the answer is yes.

Gabriele Inghirami, Hannah Elfner, *Eur. Phys. J. C* 82, 796 (2022)

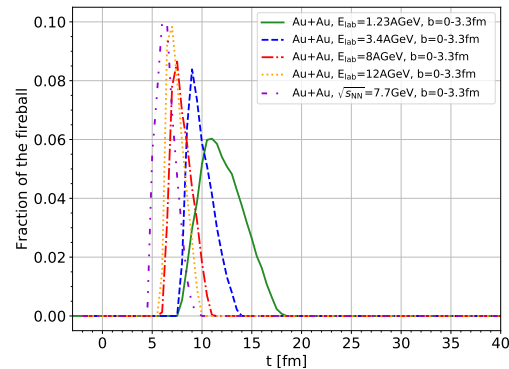
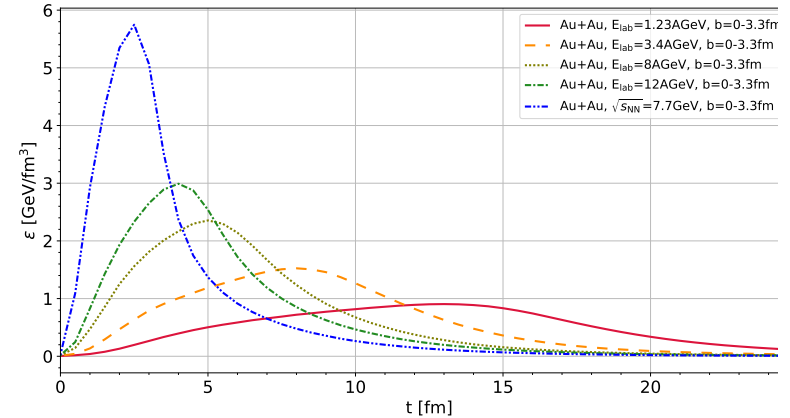
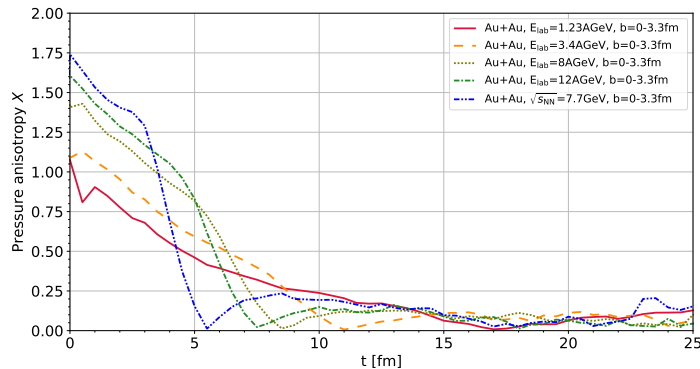
From $E_{\text{lab}} = 1.23 \text{ GeV}$ up to $\sqrt{s} = 7.7 \text{ GeV}$

Examining the pressure anisotropy:

$$T^{\mu\nu} = \sum_i \frac{p_i^\mu p_i^\nu}{p_i^0} K(\mathbf{r} - \mathbf{r}_i, \mathbf{p}_i)$$

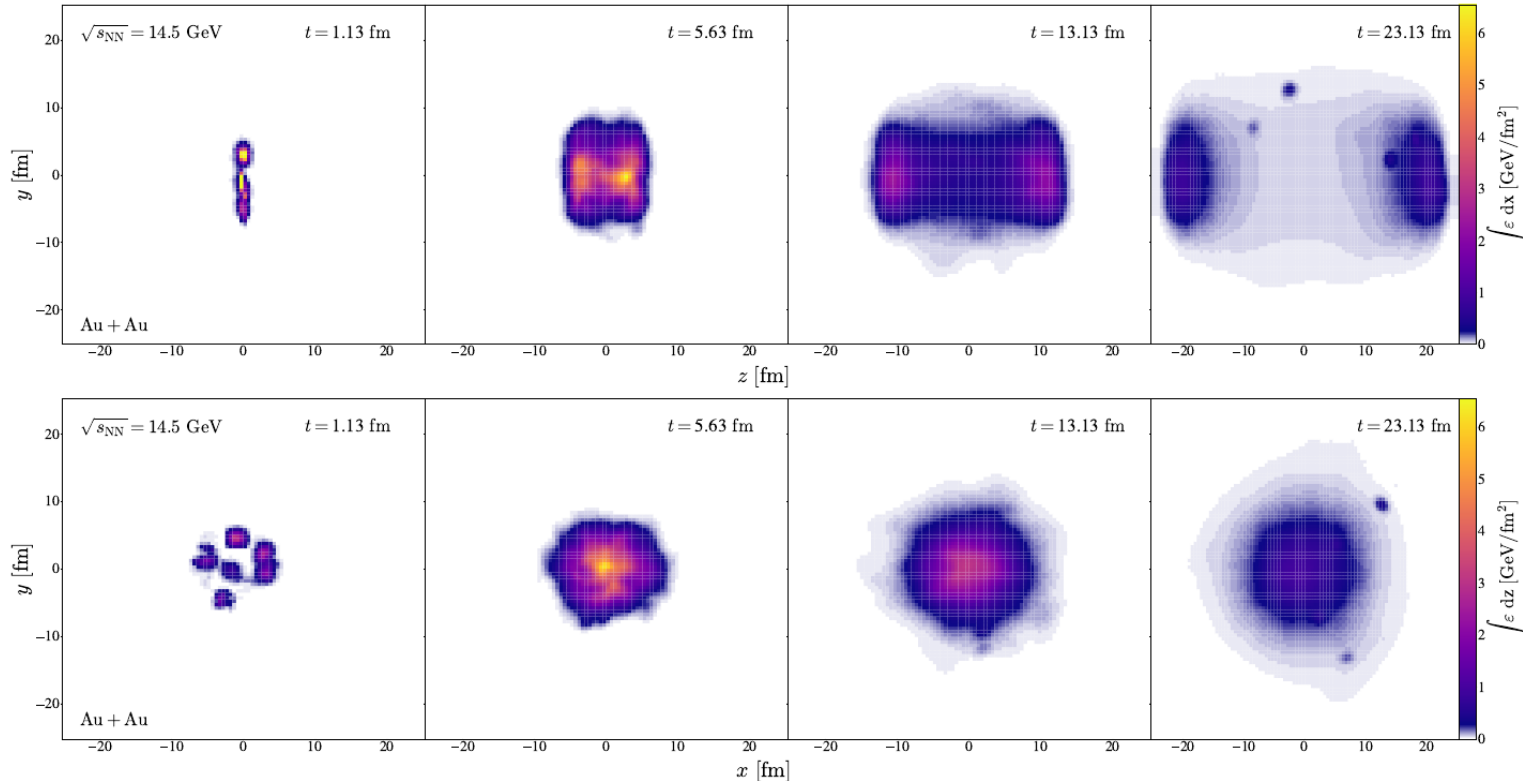
$$X \equiv \frac{|\langle T_L^{11} \rangle - \langle T_L^{22} \rangle| + |\langle T_L^{22} \rangle - \langle T_L^{33} \rangle| + |\langle T_L^{33} \rangle - \langle T_L^{11} \rangle|}{\langle T_L^{11} \rangle + \langle T_L^{22} \rangle + \langle T_L^{33} \rangle},$$

$$Y \equiv \frac{3(|\langle T_L^{12} \rangle| + |\langle T_L^{23} \rangle| + |\langle T_L^{13} \rangle|)}{\langle T_L^{11} \rangle + \langle T_L^{22} \rangle + \langle T_L^{33} \rangle}$$



A more sophisticated switching idea: dynamical fluidization

The idea: once the energy density in LRF is high enough *locally*, fluidize this part.



JAM + dynamical fluidisation + hydro + JAM

Y. Akamatsu, M. Asakawa, T. Hirano, M. Kitazawa, K. Morita, K. Murase, Y. Nara, C. Nonaka, A. Ohnishi, [Phys. Rev. C 98, 024909 \(2018\)](#)

JAM IS: HIJING string excitation + PYTHIA6 fragmentation + rescatterings of produced hadrons.

Hadrons are converted to fluid if the local energy density $e > e_f = 0.5 \text{ GeV/fm}^3$.

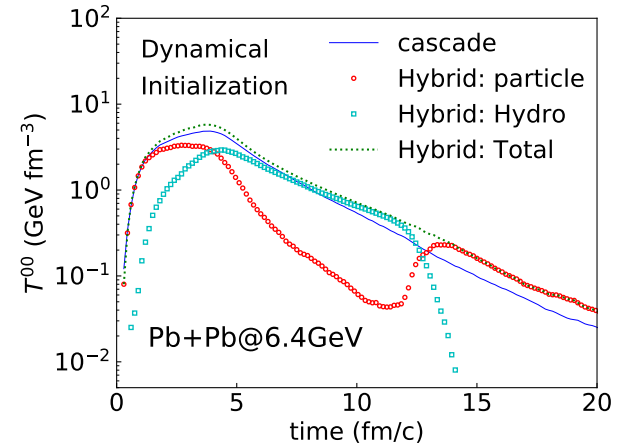
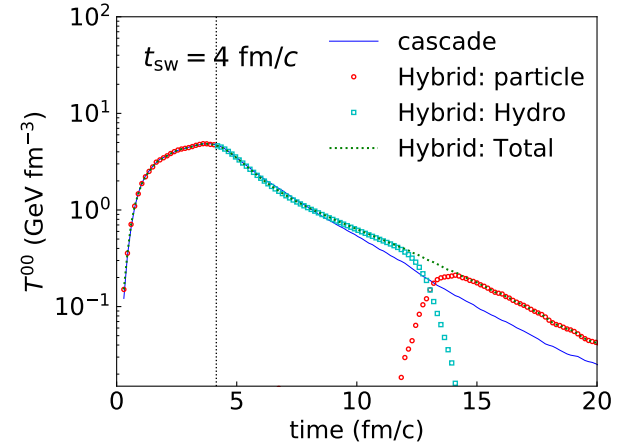
$$\partial_\mu T_f^{\mu\nu} = J^\nu, \quad \partial_\mu N_f^\mu = \rho$$

$$J^\mu(r) = \frac{1}{\Delta t} \sum_i p_i^\mu G(r - r_i(t))$$

$$\rho(r) = \frac{1}{\Delta t} \sum_i B_i G(r - r_i(t))$$

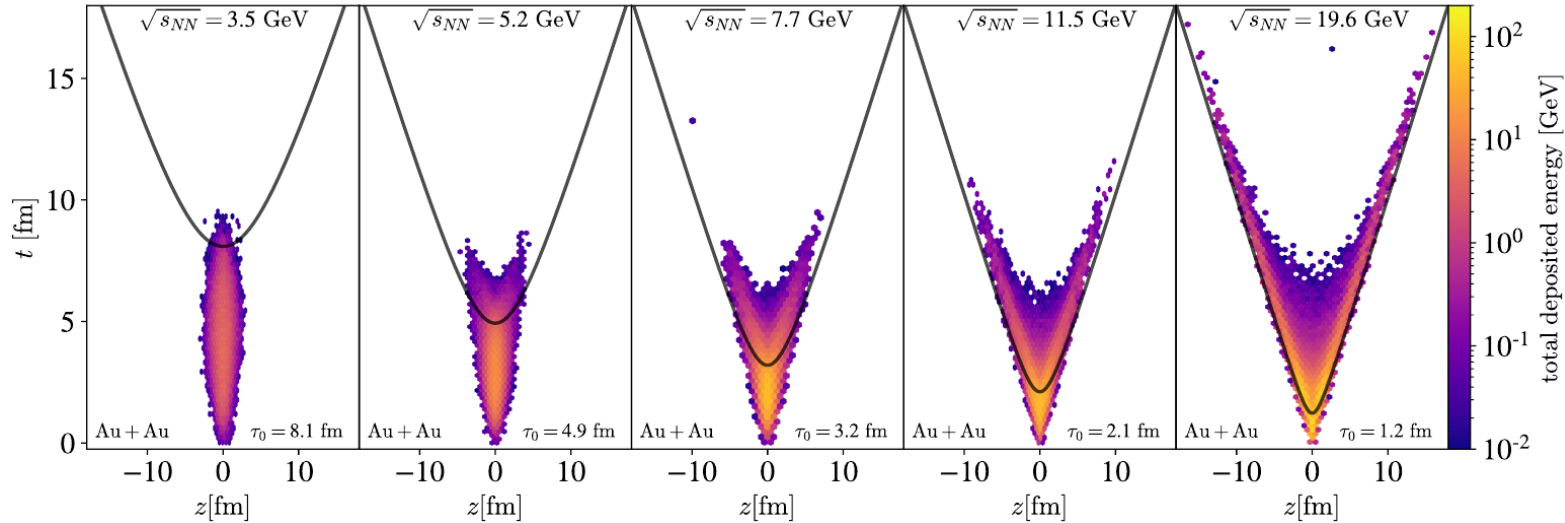
$G(r)$ is a Gaussian smearing profile.

EoS: EoS Q, a.k.a. BM+HRG EoS



As a result, the simulation does not wait for the two nuclei to completely pass through each other, which is a long time at low \sqrt{s}

Fluid approximation still applies for the most dense part of the evolution



Skipping the first switching/interface altogether: *multi-fluid approach*

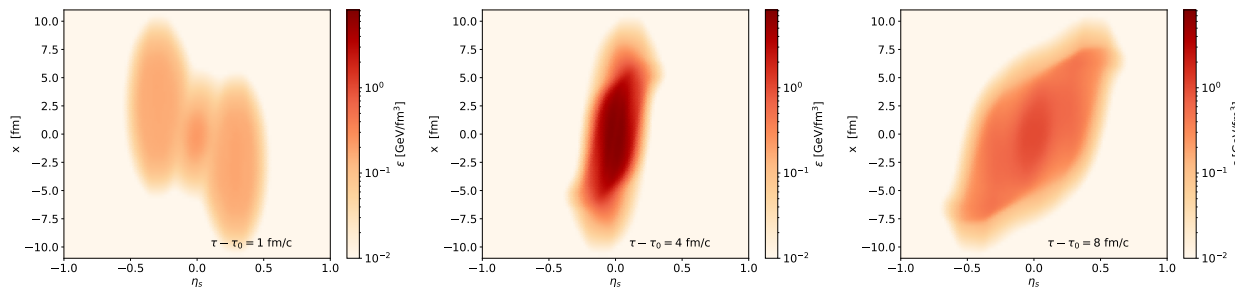
- Incoming nuclei = two fluids labelled as projectile and target
- Interaction of the fluids (slowing down) via “friction” terms
- Friction transports energy and momentum into the third fluid labelled as fireball
- It is a minimal setup to reproduce baryon stopping at low \sqrt{s} and baryon transparency at high \sqrt{s} .

$$\partial_\mu T_p^{\mu\nu}(x) = -F_p^\nu(x) + F_{fp}^\nu(x),$$

$$\partial_\mu T_t^{\mu\nu}(x) = -F_t^\nu(x) + F_{ft}^\nu(x),$$

$$\partial_\mu T_f^{\mu\nu}(x) = F_p^\nu(x) + F_t^\nu(x) - F_{fp}^\nu(x) - F_{ft}^\nu(x),$$

Cimerman, IK, Tomasik,
Huovinen,
[Phys.Rev.C 107 \(2023\) 4, 4](#)



Werthmann, IK,
Huovinen,
[Phys.Rev.C 113 \(2026\), 034908](#)

**Reichert, Spieles
, Bleicher**

Snapshots of multi-fluid evolution in $x-\eta_s$ plane, Au-Au collision at $\sqrt{s_{NN}} = 7.7$ GeV

(Hadronic) Transport

- ✓ natively 3D
- ✓ EoS can be emulated with potentials
- ✓ Exact energy-momentum, charge conservation
- ✓ no interfaces (to/from hydro)

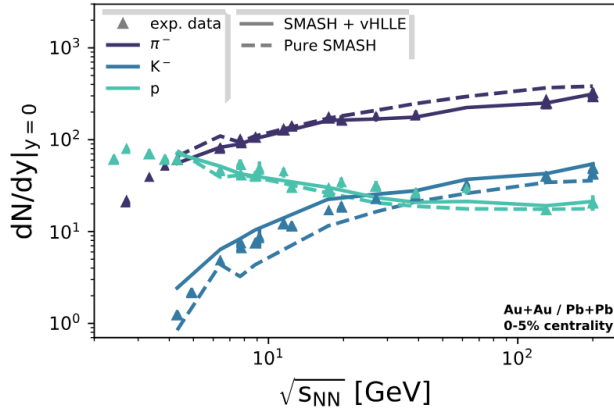
Caveats:

- ⚠ No change in the degrees of freedom (UrQMD, SMASH, HSD, JAM)
Unless: partonic phase + coalescence to hadrons (AMPT, PHSD, PHQMD(?))
- ⚠ No good agreement with the data above $\sqrt{s} > 10$ GeV (except for JAM!)

What does the modelling tell us

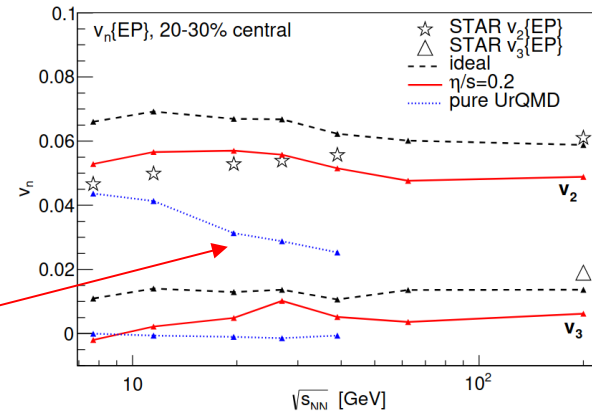
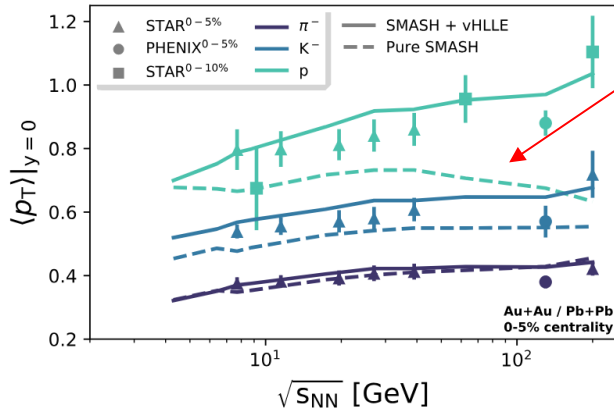
Aggregate conclusions from the hydro modelling

! Fluid phase improves agreement with the data; pure cascade fails to describe flow and mean p_T above $\sqrt{s} \sim 7$ GeV



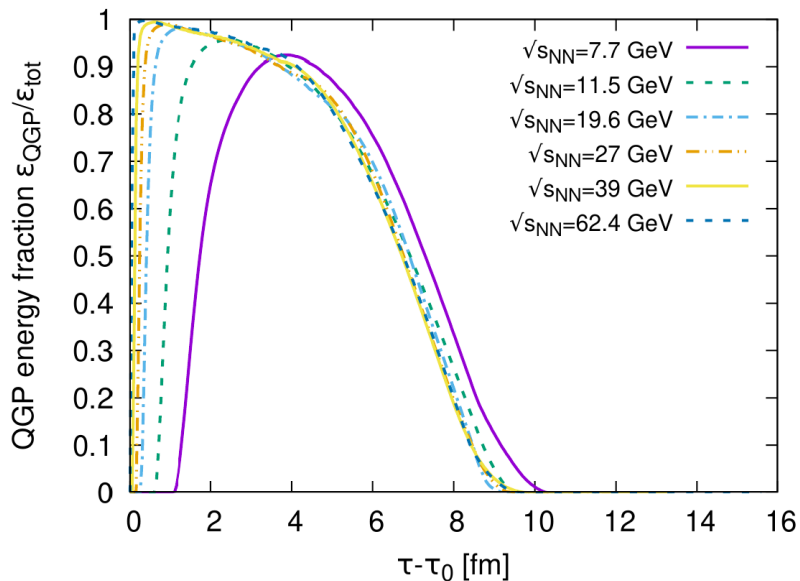
Schäfer, Karpenko, Wu, Hammelmann, Elfner,
[Eur.Phys.J.A 58 \(2022\) 11, 230](#)

Karpenko, Huovinen, Petersen, Bleicher,
[Phys.Rev.C 91 \(2015\) 6, 064901](#)

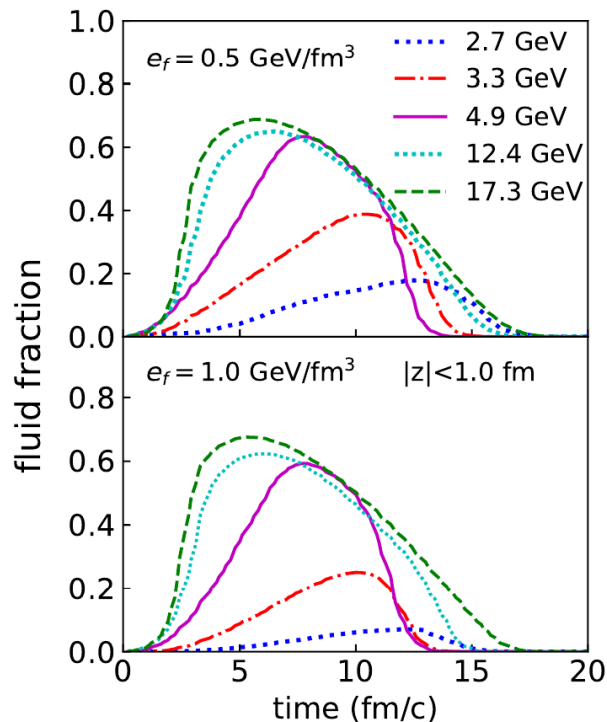


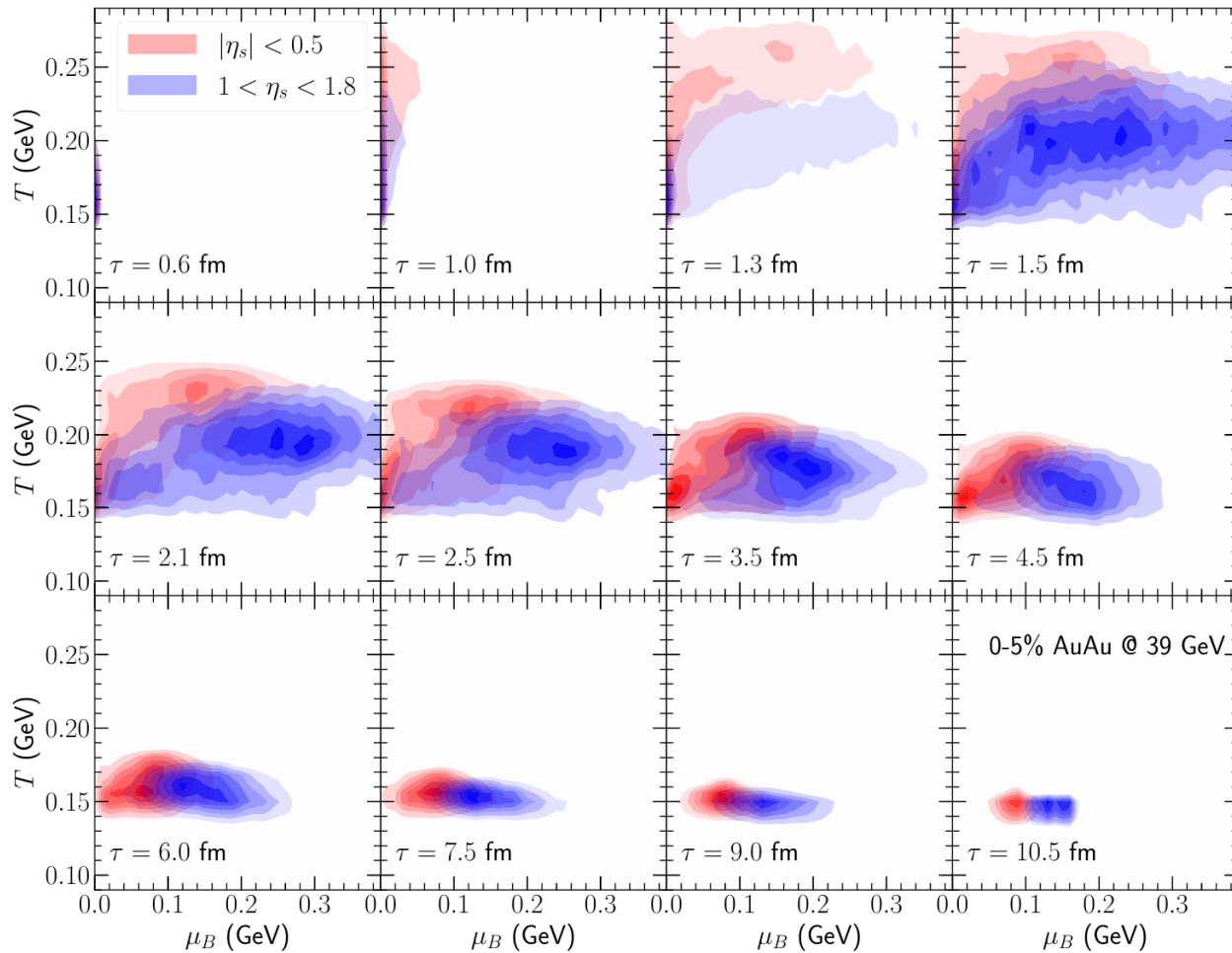
! High densities are reached at low collision energies

MUFFIN 1.0: Cimerman, Karpenko, Tomasik,
Huovinen, [Phys.Rev.C 107 \(2023\) 4](#)



JAM+hydro: Yasushi Nara et al





! The fireball is quite inhomogeneous

A \sqrt{s} point does not map to a (T, μ_B) point.

This plot is taken from Chun Shen, [2108.04987](https://arxiv.org/abs/2108.04987)

Other hybrids show a similar picture

For particle number ratios, the whole system is integrated out, which results in one point at the phase diagram

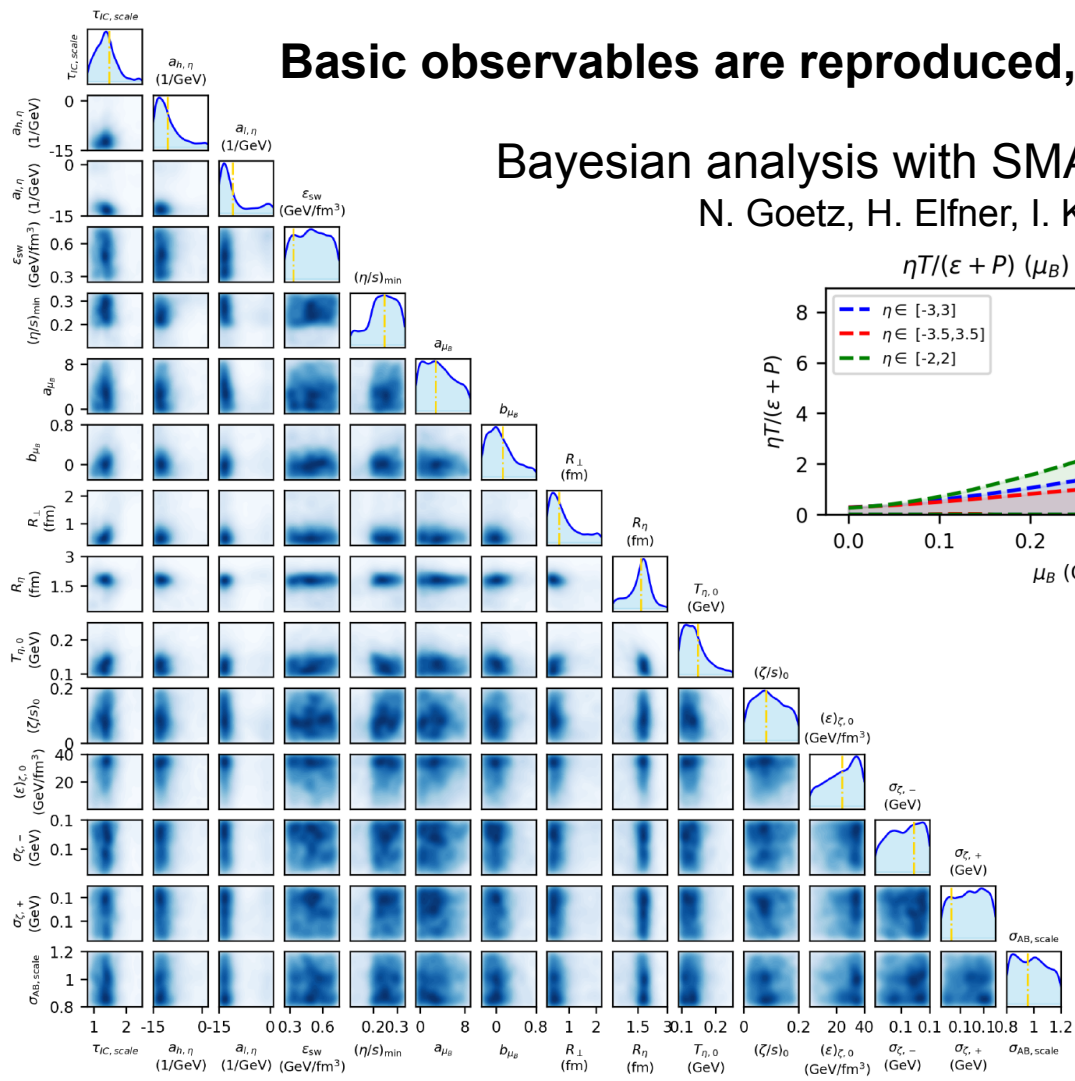
Evolution movies: <https://smash-transport.github.io/movies-hybrid.html>

Bayesian analyses

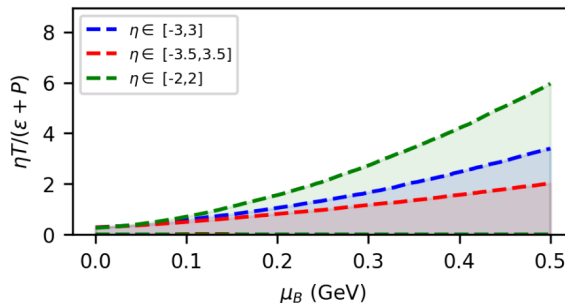
Basic observables are reproduced, models constrained via BA

Bayesian analysis with SMASH+vHLLC, $\sqrt{s}=7.7\dots 200$ GeV

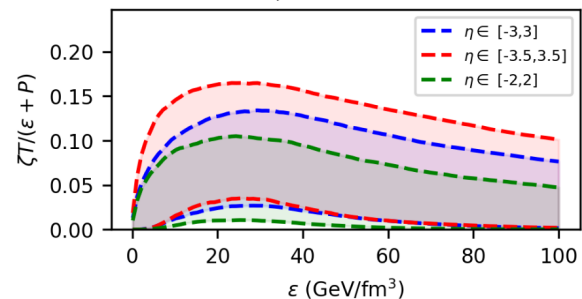
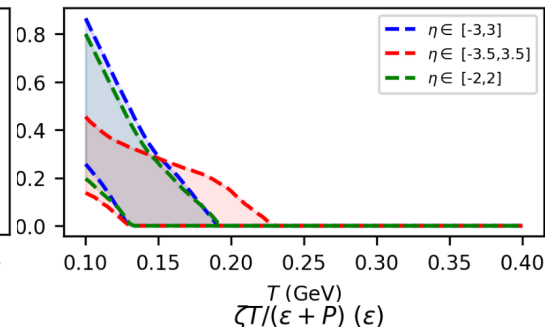
N. Goetz, H. Elfner, I. Karpenko, [Phys.Rev.C 112 \(2025\) 1, 014910](https://arxiv.org/abs/2408.14910)



$\eta T/(\epsilon + P)$ (μ_B) at $T=150$ MeV

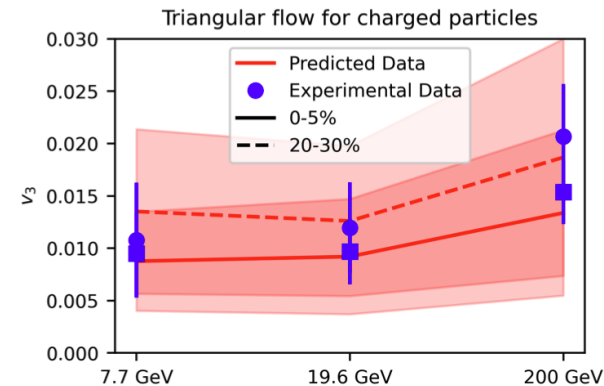
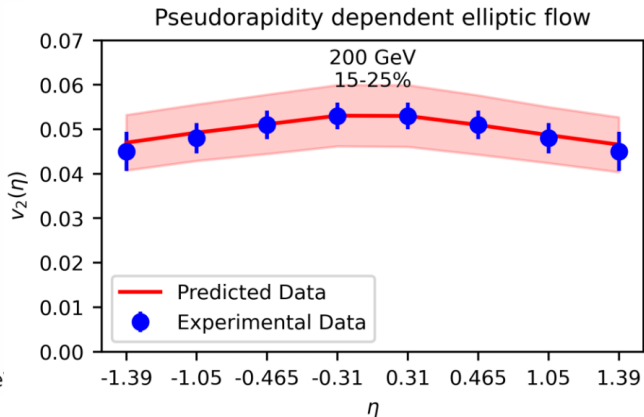
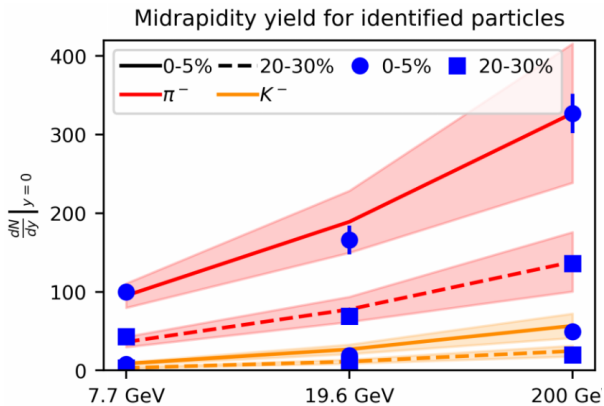
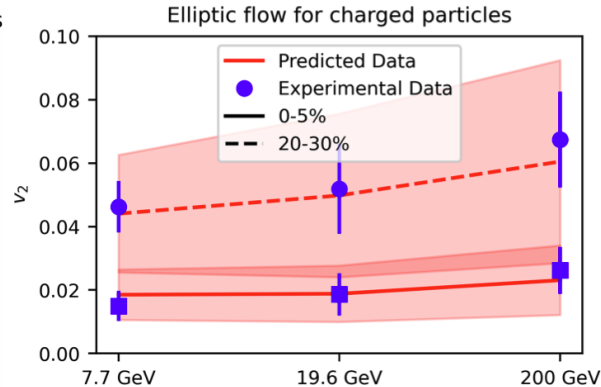
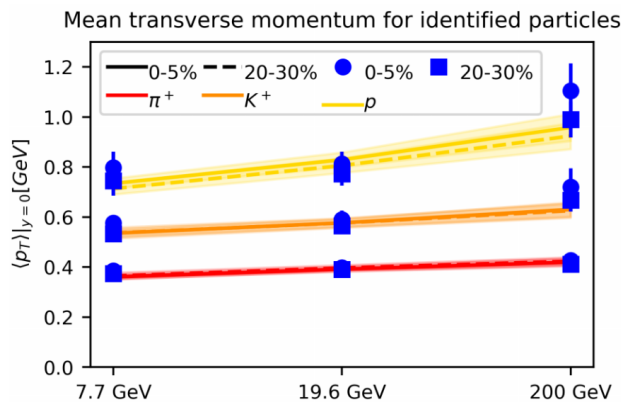
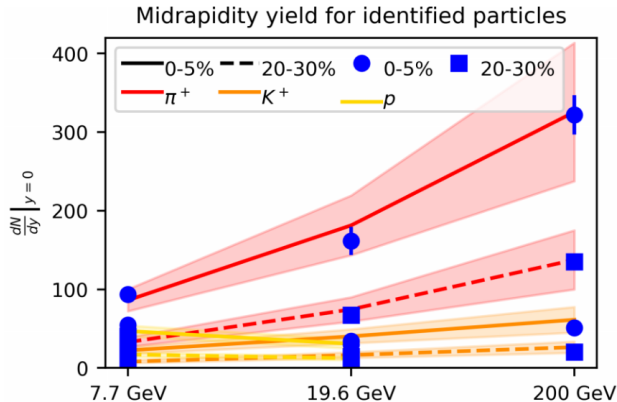


$\eta T/(\epsilon + P)$ (T) at $\mu_B=0$ MeV



Parameter space is large,
Bayesian analysis is challenging,
exp. data is not always high quality¹⁹

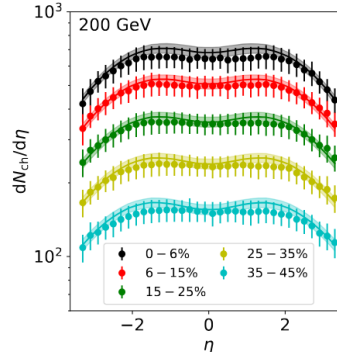
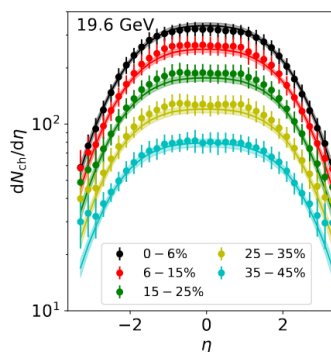
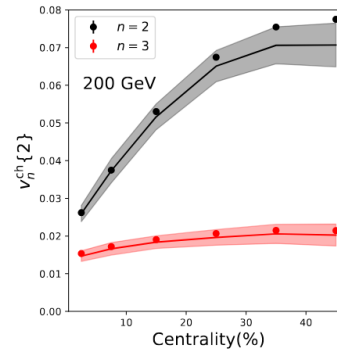
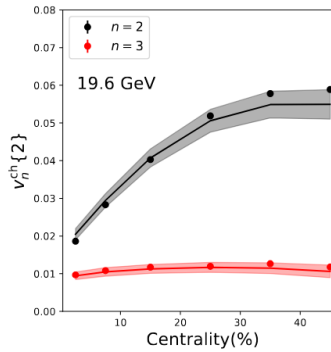
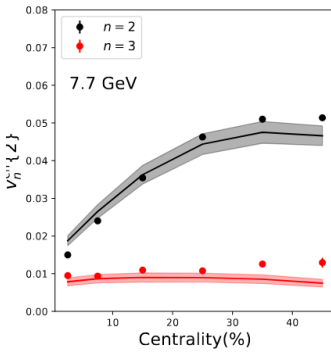
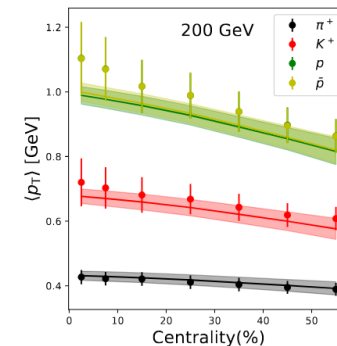
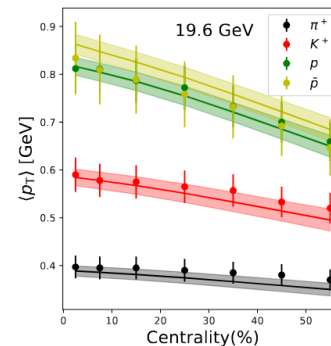
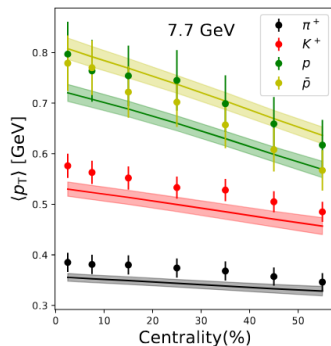
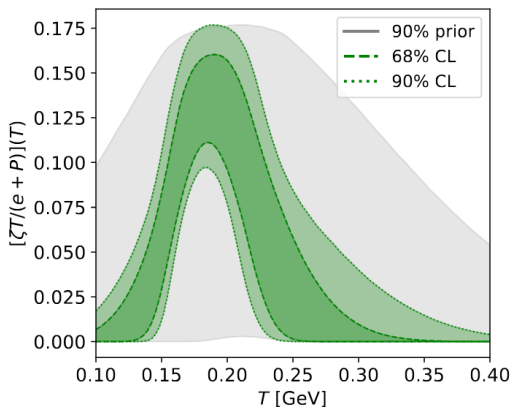
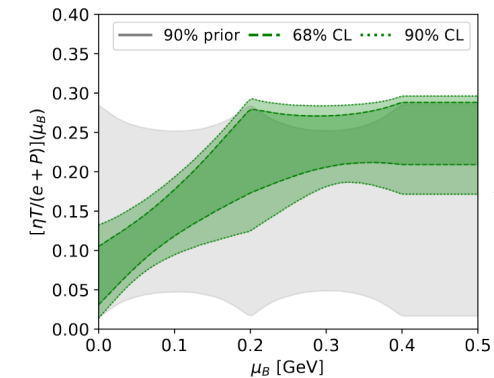
What experimental data is reproduced in (being used for) the Bayesian analysis:



Another Bayesian analysis: string deceleration IS

Jahan, Roch, Shen, [Phys.Rev.C 110 \(2024\) 5, 054905](#)

Constraints



Caveats for the BAs above:

- no agreement so far in the extracted viscosities
- tied to particular initial state(s)

EoS sensitivity / constraints?

There is no simple answer for that.

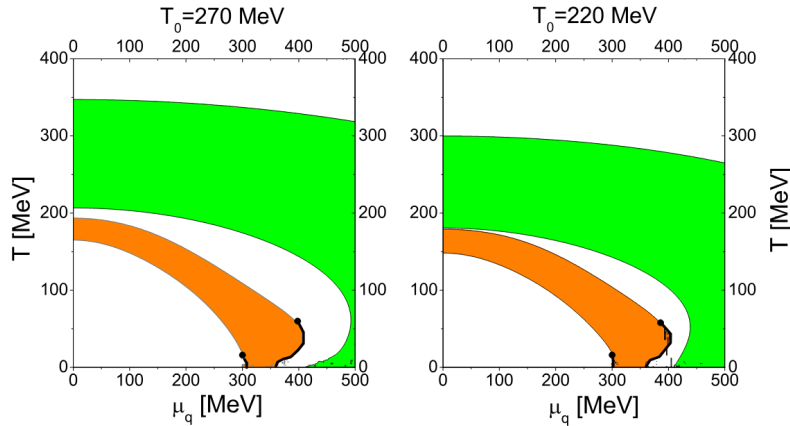
- EoS sensitivity in the models can be obscured by other sensitivities (just like at high energies).
- E.g. EoS constraining via directed flow does not seem to work.
- State-of-the-art EoS with CP location as a free parameter are applied in *some* fluid-dynamic models (MUSIC code, NEOS).
- But there is **no** BA constraining the EoS at non-zero baryon density yet (?)

New EoS at finite n_B (optionally n_Q, n_S) are coming

EoS can include first-order PT but only Maxwell construction

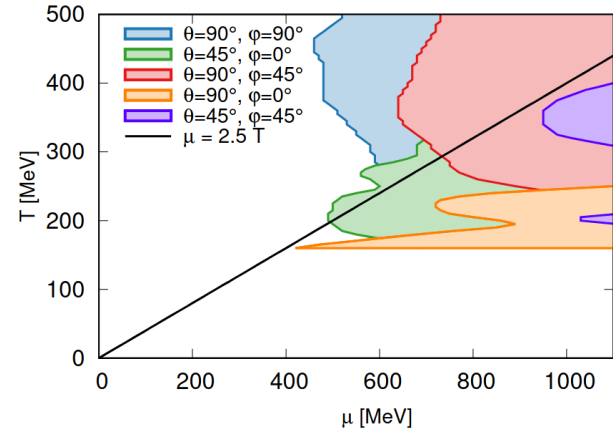
EoS has to be known/tabulated down to very low T and high μ_B → limits the candidate EoS to use.

Older choices: Chiral model EoS



Phase diagram of the EoS

State-of-the-art: **MUSES**



White region is where the used expansion works

Steinheimer, Schramm, [J.Phys.G 38 \(2011\) 035001](#)

- Energy and baryon density (2D)
- Compatible with latticeQCD at $\mu_B=0$
- Reproduces ground state of nucleus (!?)
- Crossover transition at all μ_B

Ahmed Abuali et al (MUSES), [Phys.Rev.D 112 \(2025\) 5, 054502](#)

- 4D (T, μ_B, μ_Q, μ_S)
- latticeQCD at $\mu_B=0$
- CP location parameterized

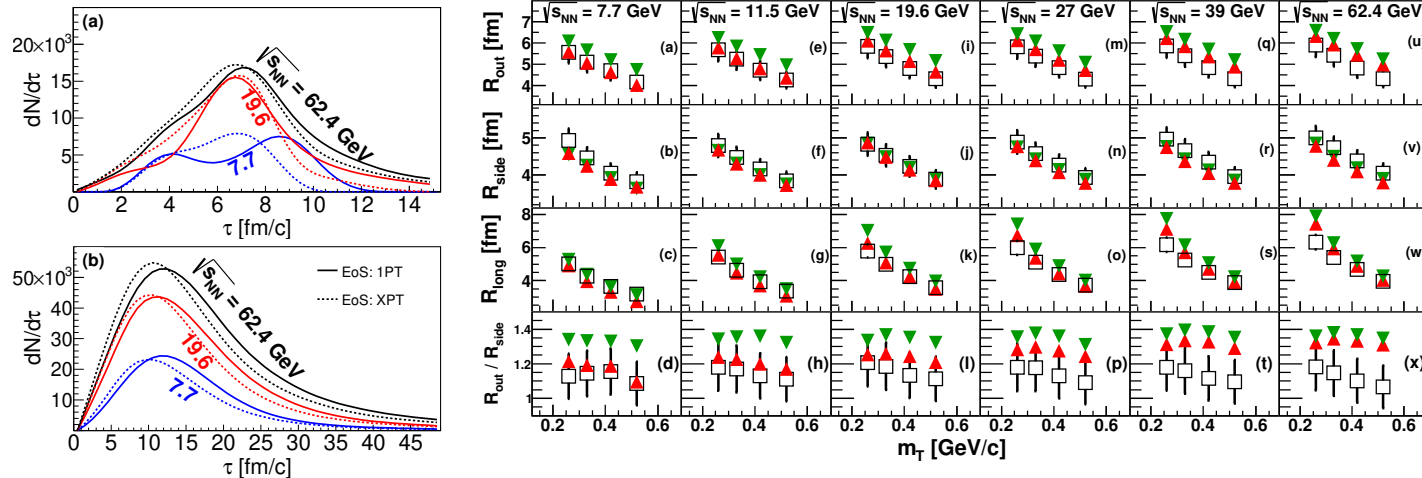
EoS/PT sensitivity via HBT/femtoscopy

Initial idea: enhancement of $R_{\text{out}}^2 - R_{\text{side}}^2$ would signal a passage through the mixed phase.

“Reality check”:

UrQMD (initial state) + fluidisation at fixed τ_0 + vHLLC (hydro) + UrQMD

Batyuk, IK, Lednicky, Malinina, Mikhaylov, Rogachevsky, Wielanek, [Phys. Rev. C 96, 024911 \(2017\)](#)



1PT = 1st order PT, XPT = crossover; crossover EoS is red, 1PT EoS is green

There is weak EoS sensitivity, crossover EoS is preferred.

Directed flow indicates a cross-over deconfinement transition in relativistic nuclear collisions

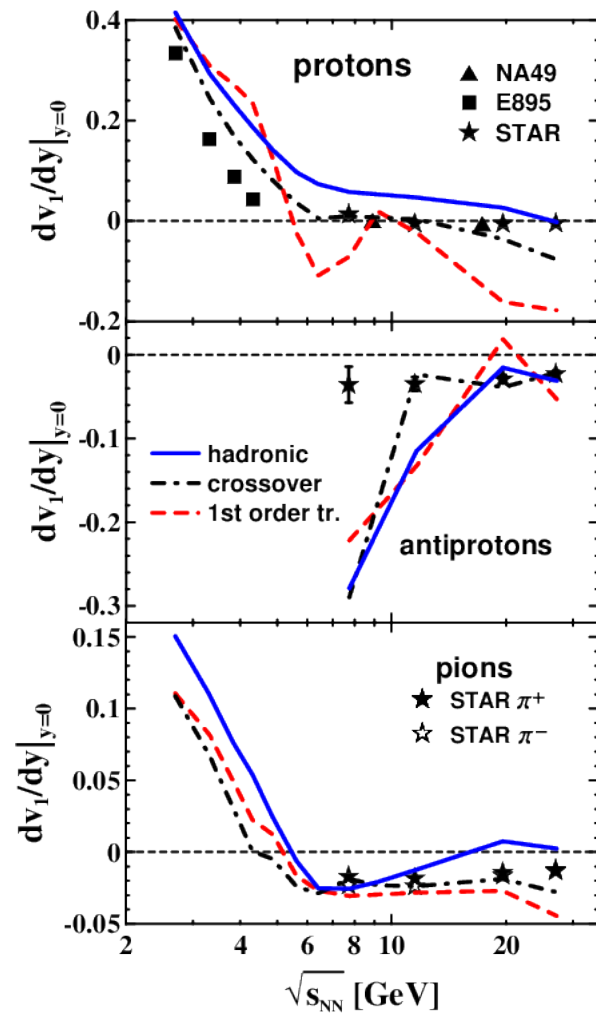
[Yu. B. Ivanov](#)^{1,2,*} and [A. A. Soldatov](#)^{2,†}

Show more

Phys. Rev. C **91**, 024915 – Published 26 February, 2015

DOI: <https://doi.org/10.1103/PhysRevC.91.024915>

Crossover-type EoS preferred



Summary

- There is a zoo of models for RHIC BES energies, covering a region of $\sqrt{s}=5-200$ GeV.
- There is WIP to extend the picture further down to a few GeV, with either dynamical fluidisation or multi-fluid dynamics.
- Most of models cover most of basic observables: rapidity distributions, net protons, radial flow, elliptic flow, (triangular flow). HBT seems to be not much off from the data.
- There has been a few Bayesian analyses that constrain
- The studies above are done mostly with crossover-type EoS. There has been no constraints on the EoS yet.
- Hadronic cascades allow to emulate EoS softening ($\sim 1PT$ EoS) but fundamentally their applicability stops (gradually) at around $\sqrt{s}=10$ GeV.