DETECTING NEUTRINOS FROM MATTER-RICH BINARY MERGERS

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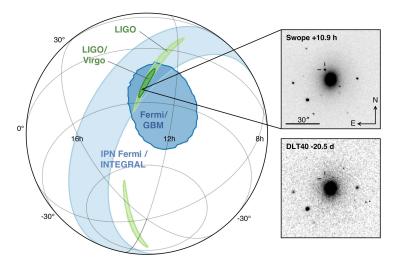
Arizona State University

NSF

Zidu Lin and CL, arXiv:1907.00034, accepted in PRD

INTRODUCTION

The multi-messenger era



• GW170817 : the first binary neutron star merger observation

- GW and electromagnetic counterpart
- Insights into origin of heavy elements, short gamma ray bursts, properties of gravity, etc.

Abbott et al. (LIGO-Virgo) ApJ Lett., 848:L12, 2017 Abbott et al., ApJ 848 (2017) no.2, L12 Abbott et al. (LIGO-Virgo, Fermi-GBM, Integral), ApJ. 848 (2017) no.2, L13

- Neutrinos still missing in binary mergers observations
- A future neutrino detection will be *important* to:
 - Learn about the post-merger phase
 - Test r-process nucleosynthesis
 - Understand the genesis of gamma ray bursts from mergers

Thermal neutrinos from mergers: detectable?

- Supernova-like emission: E \sim 10-20 MeV, E_{tot} \sim 10⁵¹ 10⁵² ergs
- BUT : mergers are at least 10² times less frequent than supernovae!
- The diffuse merger neutrino flux is detectable *in principle*
 - Up to O(1) events/year in Mt-scale detectors
 - *Overwhelmed by background* (diffuse supernova flux, atmospheric, etc.)

O. L. Caballero, G. C. McLaughlin, R. Surman, and R. Surman, PRD80, 123004 (2009). T. S. H. Schilbach, O. L. Caballero, and G. C. McLaughlin (2018), arXiv: 1808.03627.

Idea: use time-coincidence with GW for background reduction

K. Kyutoku and K. Kashiyama, PRD97, 103001 (2018) Zidu Lin and CL, arXiv:1907.00034

Next generation detectors: synergies

- HyperKamiokande neutrino detector
 - 0.3 1 Mt water Cherenkov, start construction April 2020
 - Main channel: $\bar{\nu}_e + p \rightarrow n + e^+$
 - Option of Gadolinium addition for background reduction
- Next GW projects: Voyager, Cosmic Explorer, Einstein Telescope
 - Improved distance reach, up to redshift z ~ 2
 - > 10 times enhanced observed merger rate (compared to Adv. LIGO)

K. Abe et al. (Hyper-Kamiokande) (2018), arXiv:1805.04163.

Gadolinium addition: J. F. Beacom and M. R. Vagins, PRL 93, 171101 (2004)

Advanced LIGO: J. Abadie et al. (LIGO Scientific, VIRGO), Class. Quant. Grav. 27, 173001 (2010).

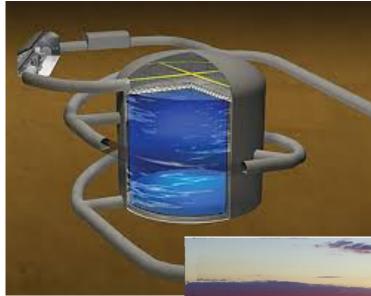
Voyager: B. P. Abbott et al., https://dcc.ligo.org/LIGO-T1600119/public (2016), the lsc-virgo white paper on instrument science (-2017 edition). Technical Report T1600119, LIGO and Virgo Scientific Collaborations.

Einstein Telescope: M. Punturo et al., Class. Quant. Grav. 27, 194002 (2010).

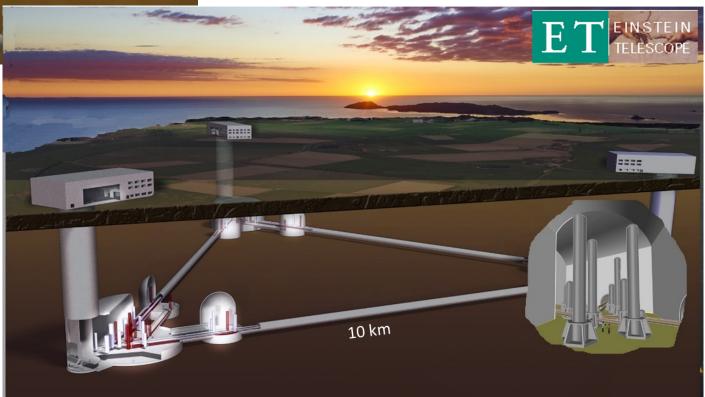
Cosmic explorer: B. P. Abbott et al. (LIGO Scientific), Class. Quant. Grav. 34, 044001 (2017).

redshift reach : C. Mills, V. Tiwari, and S. Fairhurst, Phys. Rev. D97, 104064 (2018).

http://www.hyperk.org

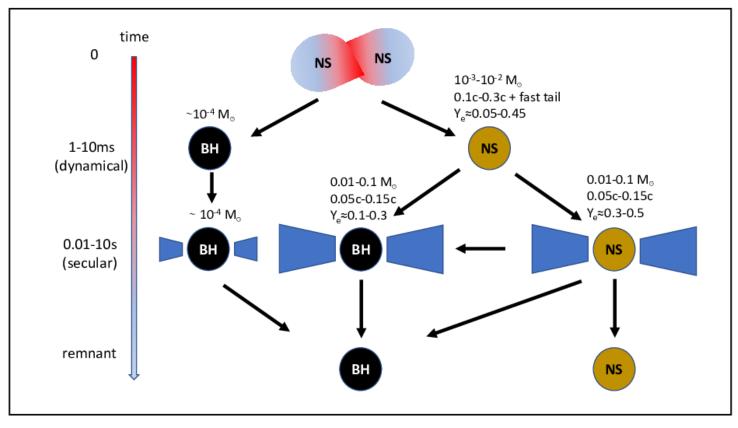


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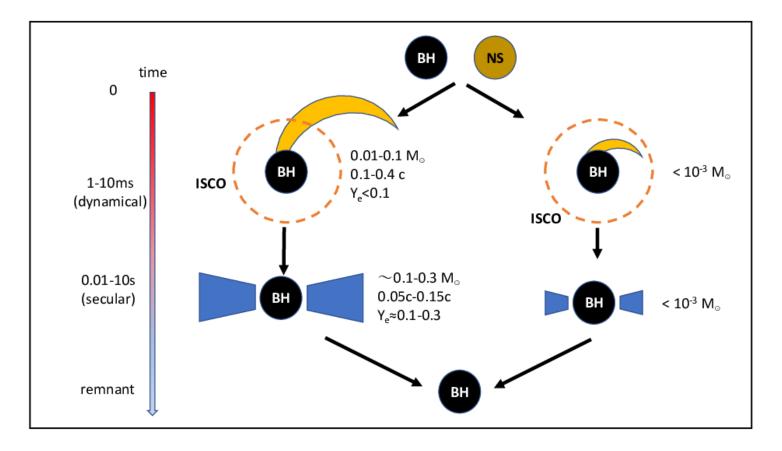
GENERALITIES: NEUTRINOS AND MERGERS

Binary neutron star (BNS) mergers



Nakar et al., arXiv:1912.05659

Neutron star-black hole (NSBH) mergers



Nakar et al., arXiv:1912.05659

Thermal neutrino emission

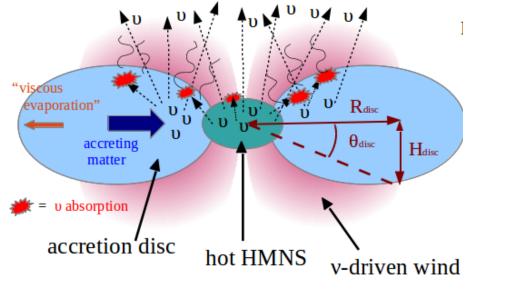


Image from Perego et al., MNRAS 443 (2014) no.4

- BNS and NSBH: neutrinos emitted from neutrinosphere in the accretion disc
- For BNS : emission from hot hypermassive neutron star

$T_{99}(s)$	$\mathcal{E}_{\bar{\nu}_e}(10^{51} erg)$	$\langle E_{\bar{\nu}_e} \rangle (MeV)$	type	\mathbf{CR}	${\rm CRM}~(M_\odot)$	DM (M_{\odot})	Model	Ref
0.58	4.4	18	BNS	HMNS	3	0.03	Hinf	J.Lippuner(2017)
0.40	2.0	16.5	BNS	BH	3	0.03	B090	
0.30	1.8	15.4	NSBH	BH	8.1	0.1	BF15	
0.10	1.0	17.8	BNS	BH	3	0.03	M3A8m03a5	O.Just(2015)
0.27	11.2	16	NSBH	BH	6	0.3	M6A8m3a5	
0.99	14	10	BNS	HMNS	2.7	0.2	DD2-1351350-On-H	S.Fujibayashi(2017)
0.58	40	20	BNS	HMNS	3		М	Y.Sekiguchi(2011)
0.08	19.8	24	NSBH	BH	6.6	0.49	A5	H.T.Janka (1999)
0.16	23.2	28	NSBH	BH	11.6	0.47	B10	

• From numerical simulations:

- T99 : duration of neutrino burst
- CR = central remnant
- D = disc

H. T. Janka, T. Eberl, M. Ruffert, and C. L. Fryer, ApJ 527, L39 (1999).
Y. Sekiguchi, K. Kiuchi, K. Kyutoku, and M. Shibata, PRL 107, 051102 (2011).
O. Just et al., MNRAS 448, 541 (2015).

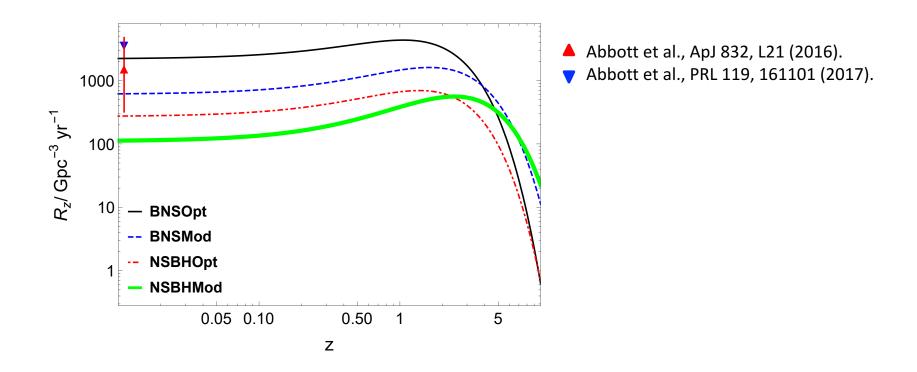
S. Fujibayashi, Y. Sekiguchi, K. Kiuchi, and M. Shibata, ApJ 846, 114 (2017).

J. Lippuner et al., MNRAS 472, 904 (2017).

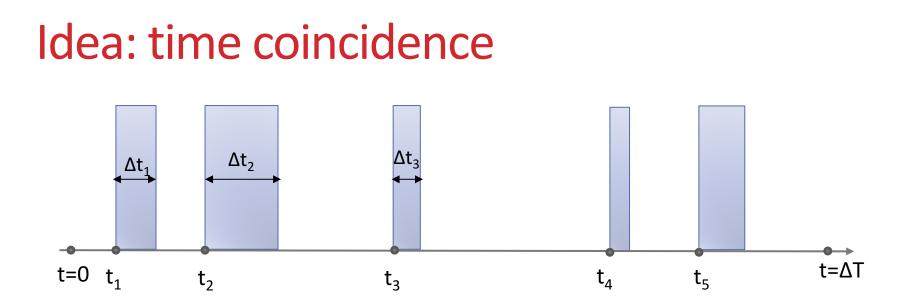
Cosmological rates

Large uncertainty from evolution models

M. Mapelli and N. Giacobbo, MNRAS 479, 4391 (2018). J. J. Eldridge, E. R. Stanway, and P. N. Tang, MNRAS 482, 870 (2019).



DETECTABILITY



- From GW: i=1,2,3,...,N mergers. Merger type, time t_i, redshift z_i
- From theory: estimated neutrino burst duration, Δt_i
 - Conservative: $\Delta t_i = \Delta t = 1 \text{ s}$. Captures entire neutrino emission.
- <u>to reduce background</u>: keep only neutrino data in [t_i, t_i+Δt]; impose z < z_{max};

K. Kyutoku and K. Kashiyama, PRD97, 103001 (2018) Zidu Lin and CL, arXiv:1907.00034

Event rates

• (Observed) merger rate and number of background events:

$$\frac{dV}{dz} = \frac{4\pi D_c^2 c}{H(z)} \qquad D_c = \int_0^z \frac{c}{H_0 \sqrt{\Omega_m (1+\tilde{z})^3 + \Omega_\Lambda}} d\tilde{z}$$

$$N_{z_{max}} = \int_0^{z_{max}} \frac{R(z)dV}{(1+z)dz}dz$$

$$N_b \simeq (N_{z_{max}}\Delta T)(n_b\Delta t)$$

- Number of signal events (inverse beta decay in water):
 - Diffuse flux approximation

$$F_{\bar{\nu}_e}(E) = \frac{\mathcal{E}_{\bar{\nu}_e}}{\langle E_{\bar{\nu}_e} \rangle} \frac{2}{3T^3 \zeta(3)} \frac{E^2}{e^{E/T} + 1},$$

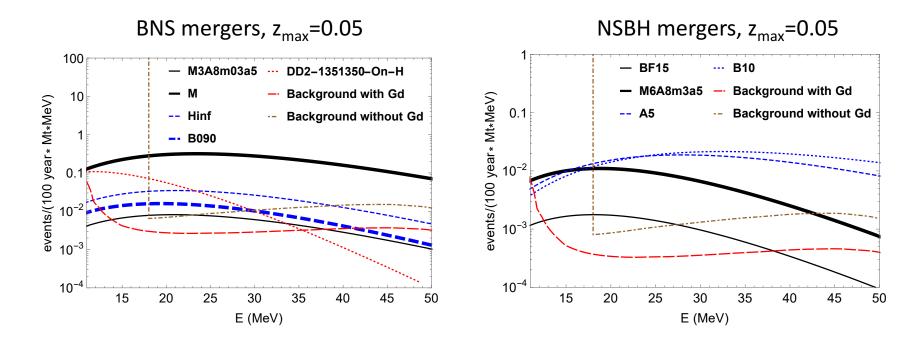
$$\Phi_{\bar{\nu}_e}(E) = f_{osc} \frac{c}{H_0} \int_0^{z_{max}} R(z) F_{\bar{\nu}_e}(E') \frac{dz}{\sqrt{\Omega_m (1+z)^3 + \Omega_\Lambda}}, \qquad E' = E(1+z)$$
$$f_{osc} = 1$$

(phenomenological factor to account for oscillations)

$$N_s \simeq N_t \Delta T \int_{E_{th}}^{E_{max}} \eta(E) \Phi_{\bar{\nu}_e}(E) \sigma(E) dE$$

$$(\bar{\nu}_e + p \rightarrow n + e^+)$$

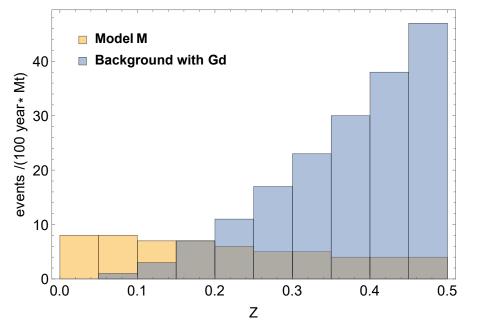
Results: event rates at Mt-scale detector



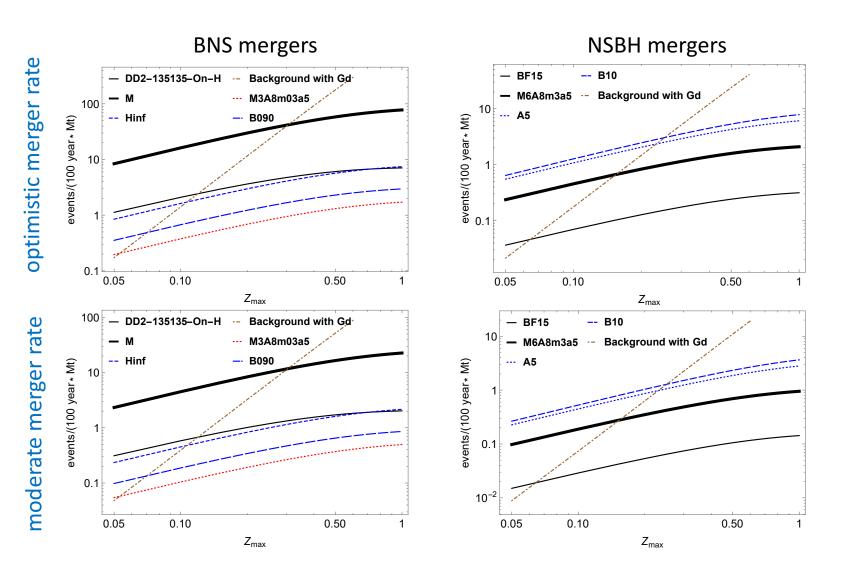
- z_{max}=0.05 ; optimistic merger rates used
- Signal/background ≥ 1 in suitable energy window

redshift dependence

BNS mergers, event rates in redshift bins



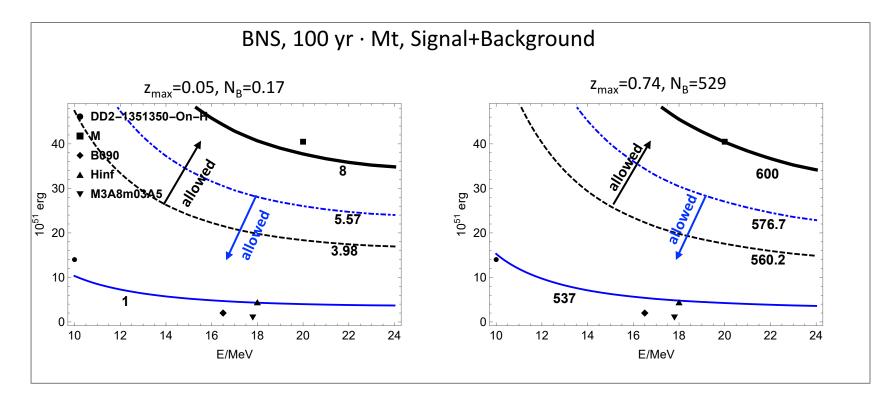
- for most optimistic parameters:
 - Up to ~ 100 signal events from mergers at z<0.5.
 - Signal/background $\gtrsim 1$ for z_{max} < 0.15



• Number of events for $0 < z < z_{max}$. 3 σ significance achieved for z_{max} up to ~0.7

Constraining the parameter space

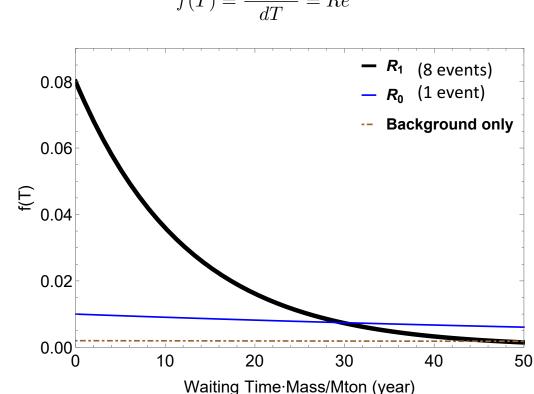
• Exclusion of extreme parameters possible



Solid: central value; non-solid: 90% CL contours. Merger rate assumed known.

Constraints from time to first detection

- Even zero or one event observed can give constraints!
 - Probability that first detection occurs in interval [T, T + dT]



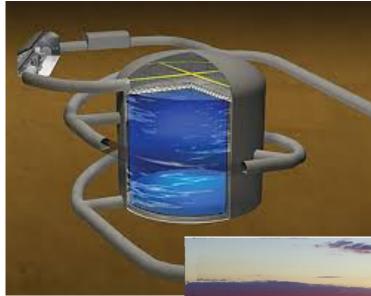
$$f(T) = \frac{dF(T)}{dT} = Re^{-RT}$$

CONCLUSIONS

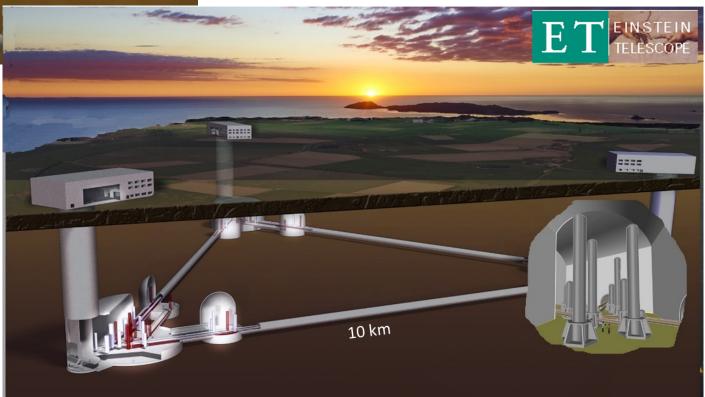
A truly multi-messenger scenario

- Detecting eutrinos from binary mergers is realistic
 - *Only* with synergy between next generation neutrino detectors and GW detectors (time coincidence)
- O(0.1 100) events expected for 100 yr · Mt exposure
 - Decades might pass before first detection
 - Single detections might be statistically significant (if low z)
- Even low statistics (or non-detection) can constrain the parameter space
 - Implications on post-merger physics (type of remnant, r-process, etc.)

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BACKUP

Merger rate

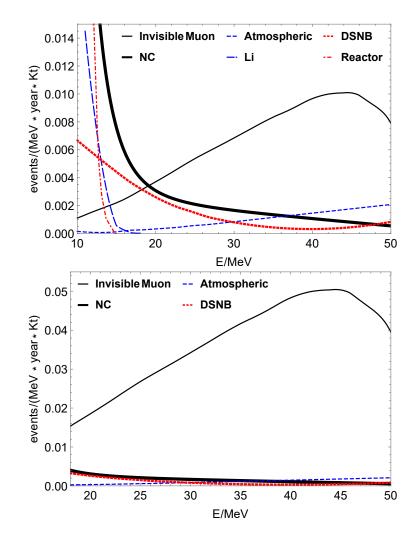
$$\frac{dV}{dz} = \frac{4\pi D_c^2 c}{H(z)},$$

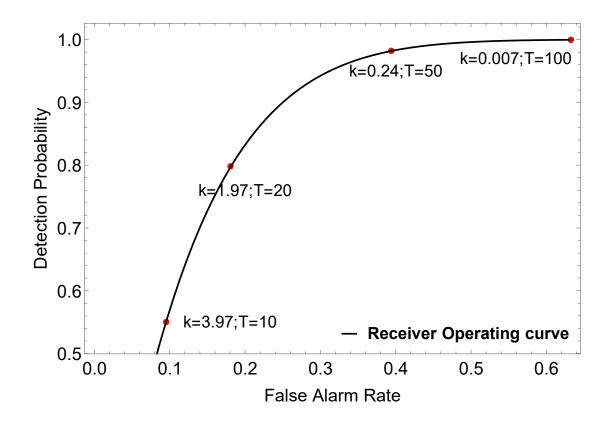
$$D_c = \int_0^z \frac{c}{H_0 \sqrt{\Omega_m (1+\tilde{z})^3 + \Omega_\Lambda}} d\tilde{z}$$

$$N(z_{max}) = \int_0^{z_{max}} \frac{R(z) dV}{(1+z) dz} dz$$

$$I = \int_0^{z_{max}} \frac{R(z) dV}{(1+z) dz} dz$$

Background composition





H. T. Janka, T. Eberl, M. Ruffert, and C. L. Fryer, Astrophys. J. 527, L39 (1999).

Y. Sekiguchi, K. Kiuchi, K. Kyutoku, and M. Shibata, Phys. Rev. Lett. 107, 051102 (2011).

O. Just et al., Mon. Not. Roy. Astron. Soc. 448, 541 (2015).

S. Fujibayashi, Y. Sekiguchi, K. Kiuchi, and M. Shibata, Astrophys. J. 846, 114 (2017).

J. Lippuner et al., Mon. Not. Roy. Astron. Soc. 472, 904 (2017).