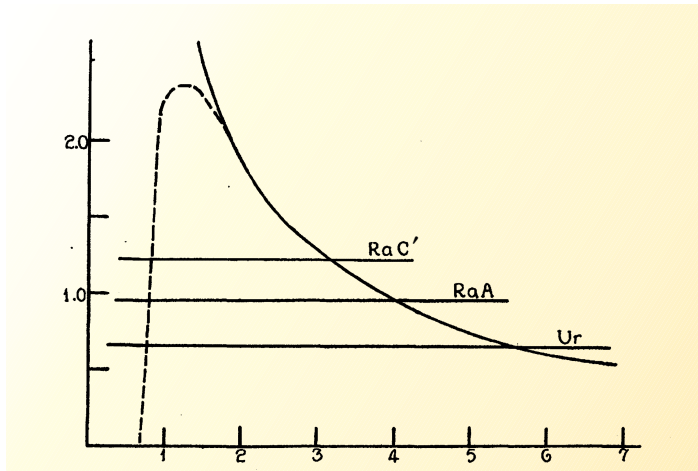


# Superaligned Alpha Decay

Rod Clark

# Wave Mechanics and Radioactive Disintegration

R. W. Gurney and E.U. Condon, Nature (London) 122, 439 (1928)

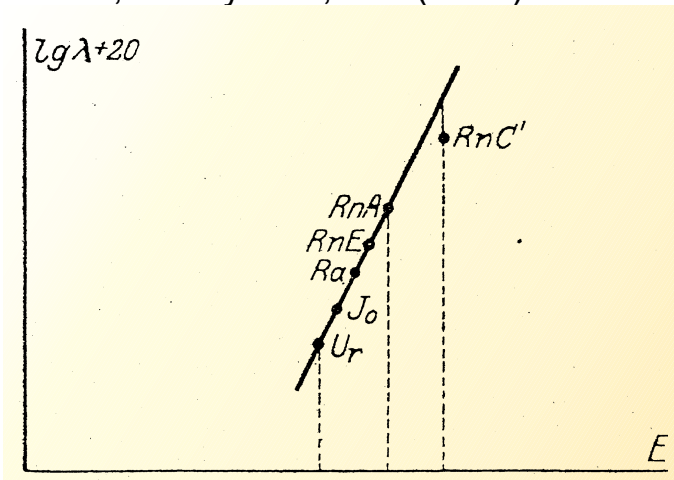


|           |                   | $E_\alpha$<br>(MeV) | $T_{1/2}$         |
|-----------|-------------------|---------------------|-------------------|
| Radium C' | $^{214}\text{Po}$ | 7.7                 | 160 $\mu\text{s}$ |
| Radium A  | $^{218}\text{Po}$ | 6.0                 | 3 mins            |
| Uranium   | $^{238}\text{U}$  | 4.2                 | 4.5 GYrs          |



## Zur Quantentheorie des Atomkernes

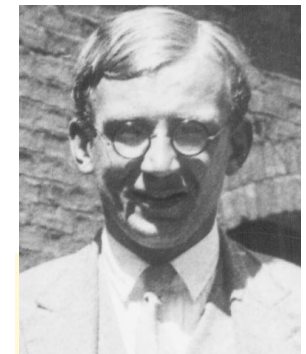
G. Gamow, Z. Phys. 51, 204 (1928)



Quantitative description of Geiger-Nuttall Rule:

$$\log(T_{1/2}) = a + bQ_\alpha^{-1/2}$$

H. Geiger and J.M. Nuttall, Phil. Mag. 22 613 (1911)



# Superfluid Tunneling Model (STM)

The Schrödinger equation describing the model is:

$$\left[ -\frac{\hbar^2}{2D} \frac{\partial^2}{\partial \xi^2} + V(\xi) \right] \psi_n(\xi) = E_n \psi_n(\xi)$$

$\xi$  = generalized deformation variable

Calculation of decay constant:

$$\lambda = P \cdot f \cdot T$$

P= preformation of decay configuration

f = frequency of hitting barrier

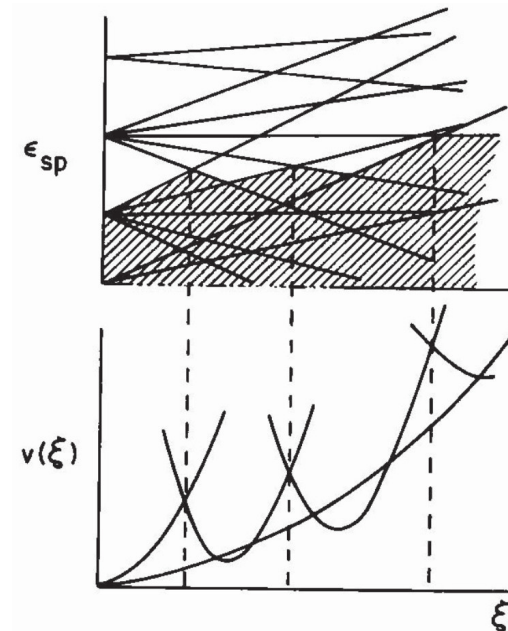
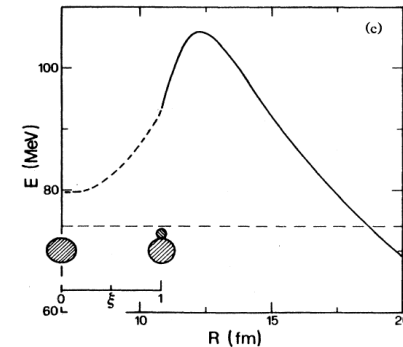
T = transmission coefficient through barrier

“Nuclear Superfluidity: Pairing in Finite Systems”

David M. Brink and Ricardo A. Broglia

Cambridge University Press, 2005

F. Barranco, G.F. Bertsch, R.A. Broglia, E. Vigezzi,  
NPA 512 253 (1990)



# Decay Constant

$$\lambda = P f T$$

Alpha-particle formation probability,

$$P = |\psi(\xi = 1)|^2 = \left(\frac{\alpha}{\sqrt{\pi}}\right) e^{-\alpha^2}, \quad \alpha^2 = \sqrt{\frac{DC}{\hbar^2}}$$

Knocking frequency,

$$f = \frac{\omega}{2\pi} = \frac{1}{2\pi} \sqrt{\frac{C}{D}}$$

Potential parameter,  $C$ ,  
is dependent on  $Q$ -value

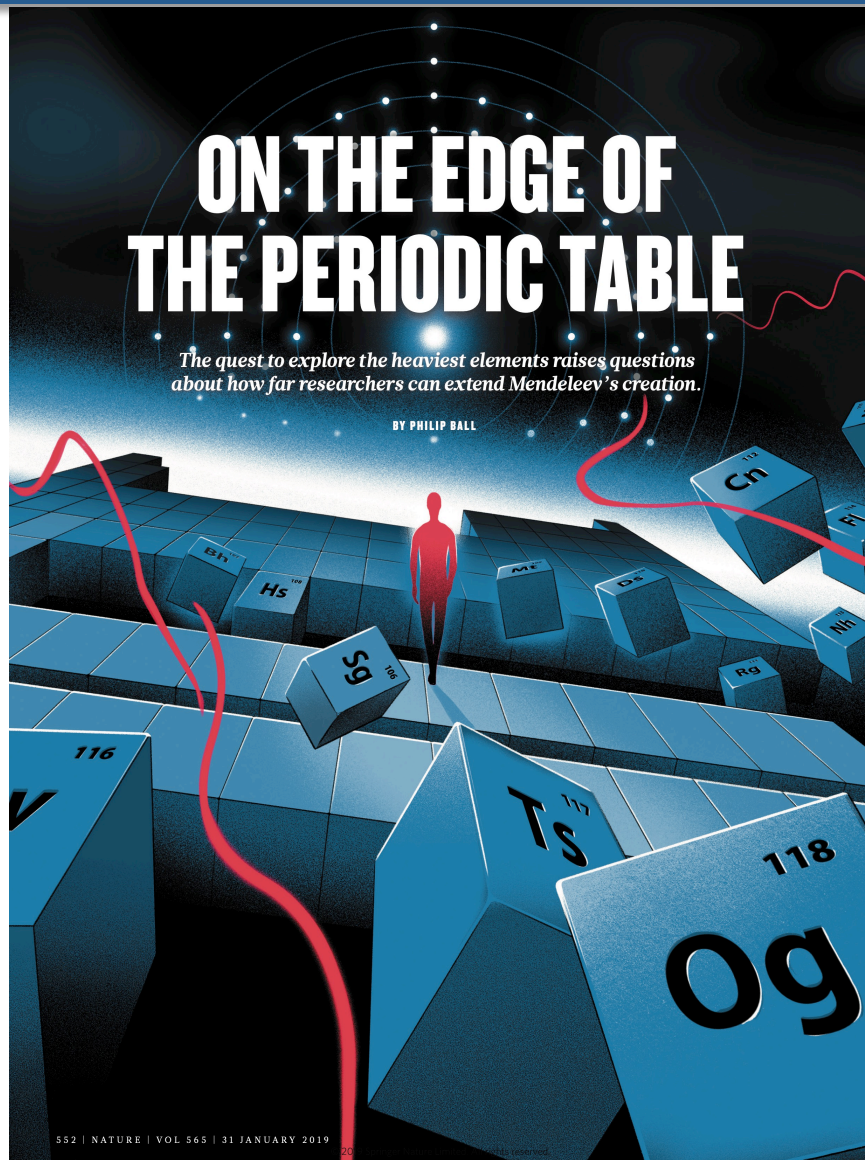
Inertial mass parameter,  $D$ ,  
depends on pairing gap,  $\Delta$ .

Transmission through barrier,

$$T_L = \frac{\rho}{F_L^2(\eta, \rho) + G_L^2(\eta, \rho)}$$

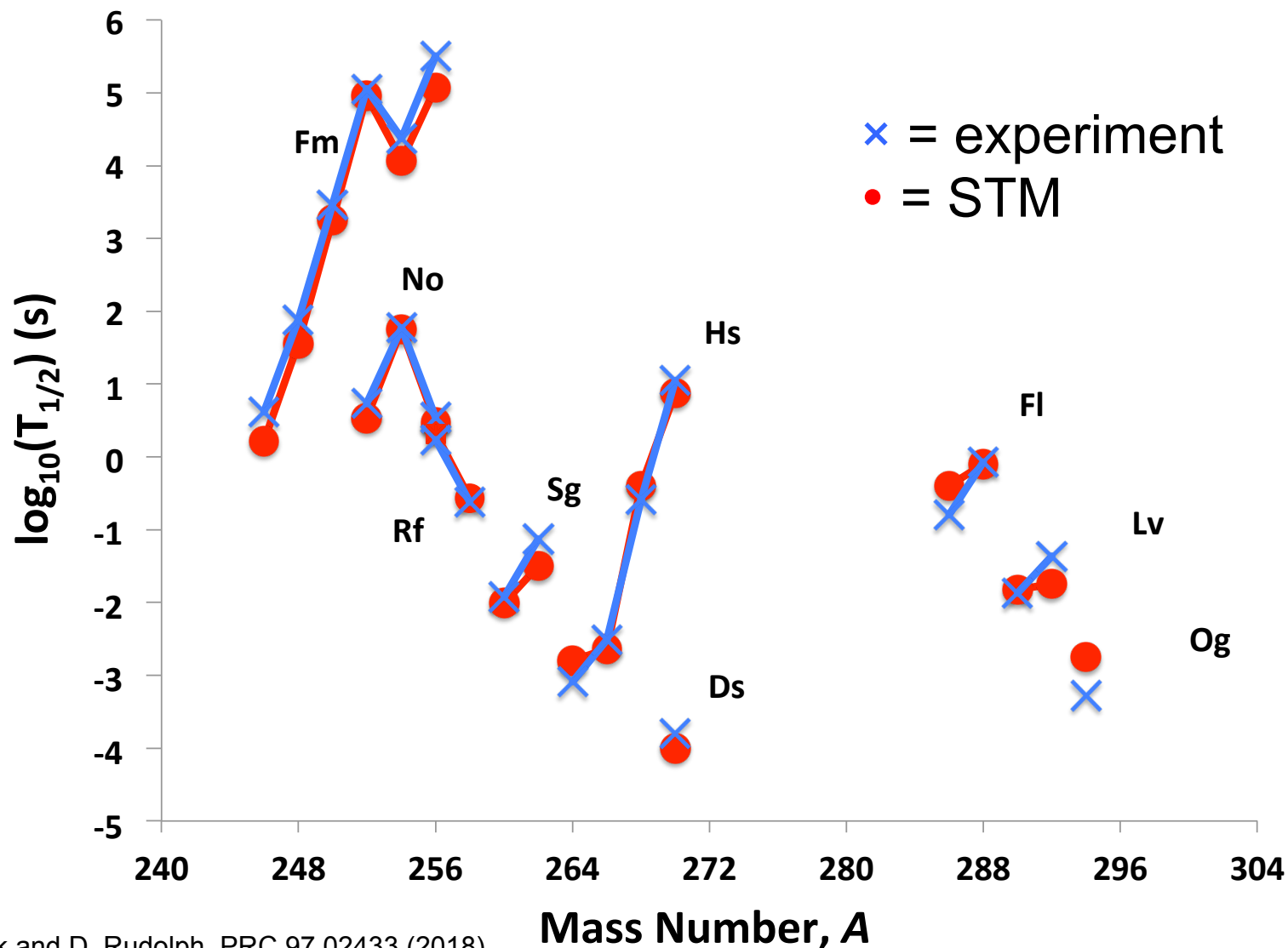
→ Investigate dependencies on  $\Delta$ ,  $Q$ ,  $L$

# Superheavy Nuclei



Nature **565**,  
553 (2019).

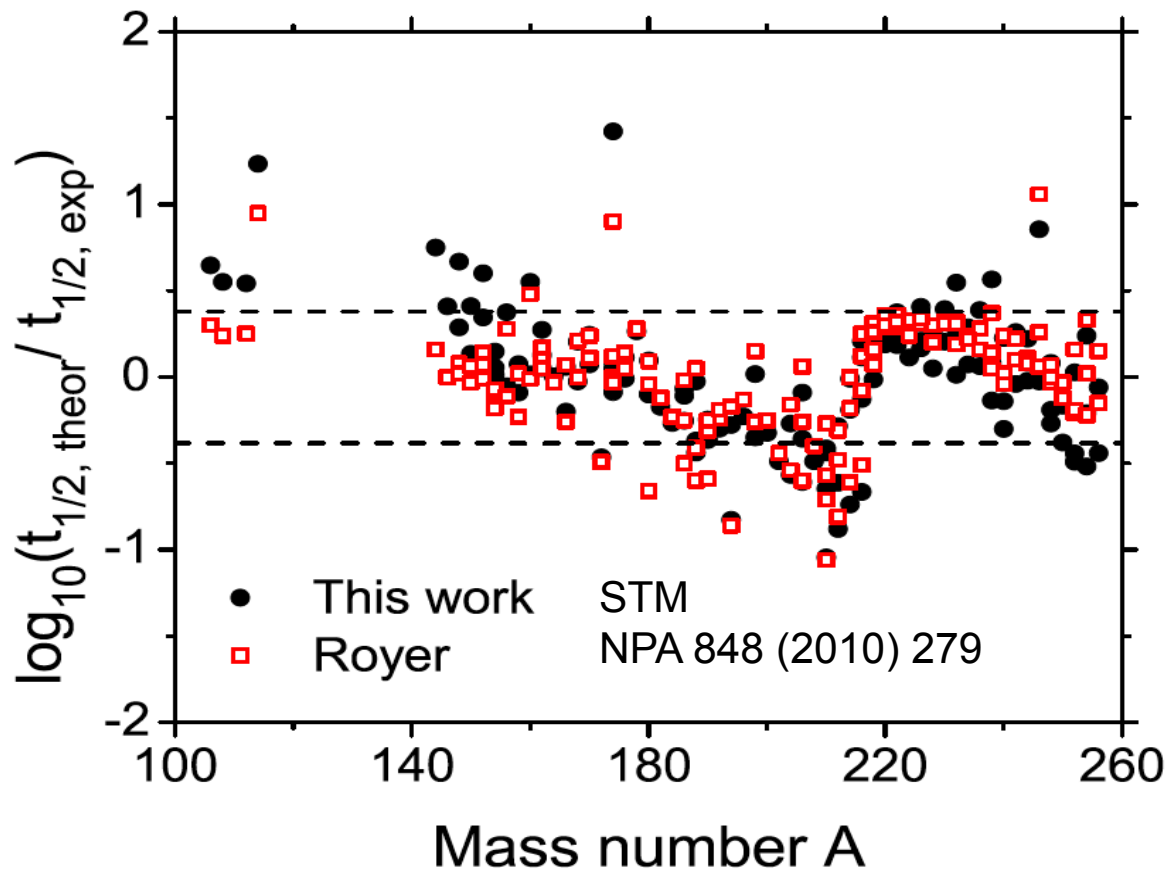
# Alpha Decay of Even-Even Isotopes: Fm to Og



R.M. Clark and D. Rudolph, PRC 97 02433 (2018)

Mass Number,  $A$

# Modern “Geiger-Nuttall” Formula



$$\log_{10}[T] = -27.750 - 1.1138A^{\frac{1}{6}}\sqrt{Z} + \frac{1.6378Z}{\sqrt{Q}} + \frac{1.7383 \times 10^{-6} ANZ[l(l+1)]^{\frac{1}{4}}}{Q} + 0.002457A[1 - (-1)^l]$$

$\left. \vphantom{\log_{10}[T]} \right\} \times 4 \left\{ \begin{array}{l} \text{even-odd} \\ \text{odd-even} \\ \text{odd-odd} \\ \text{even-even} \end{array} \right.$

# Superaligned Alpha Decay

VOLUME 14, NUMBER 4

PHYSICAL REVIEW LETTERS

25 JANUARY 1965

## NEW REGION OF ALPHA RADIOACTIVITY\*

Ronald D. Macfarlane

Department of Chemistry, McMaster University, Hamilton, Ontario, Canada

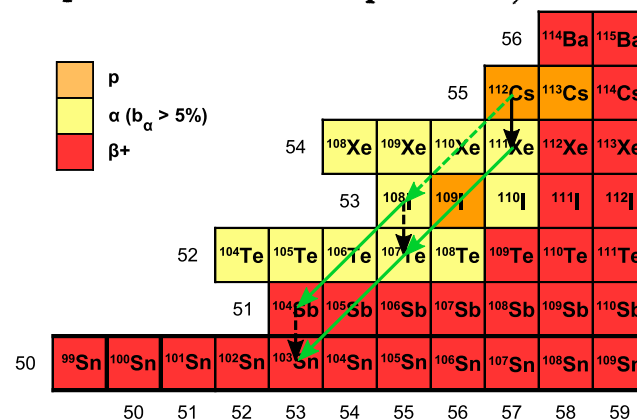
and

Antti Siivola

Lawrence Radiation Laboratory, University of California, Berkeley, California

(Received 30 November 1964)

These nuclides represent the first opportunity to study alpha decay from nuclei where the “valence” neutrons and protons are in the same single-particle level, in this case, the  $1g_{7/2}$  level. This may give rise to a kind of “super-aligned” alpha decay resulting in large reduced alpha widths. At present, we cannot



K. Auranen et al., Phys. Lett. B 792 (2019) 187

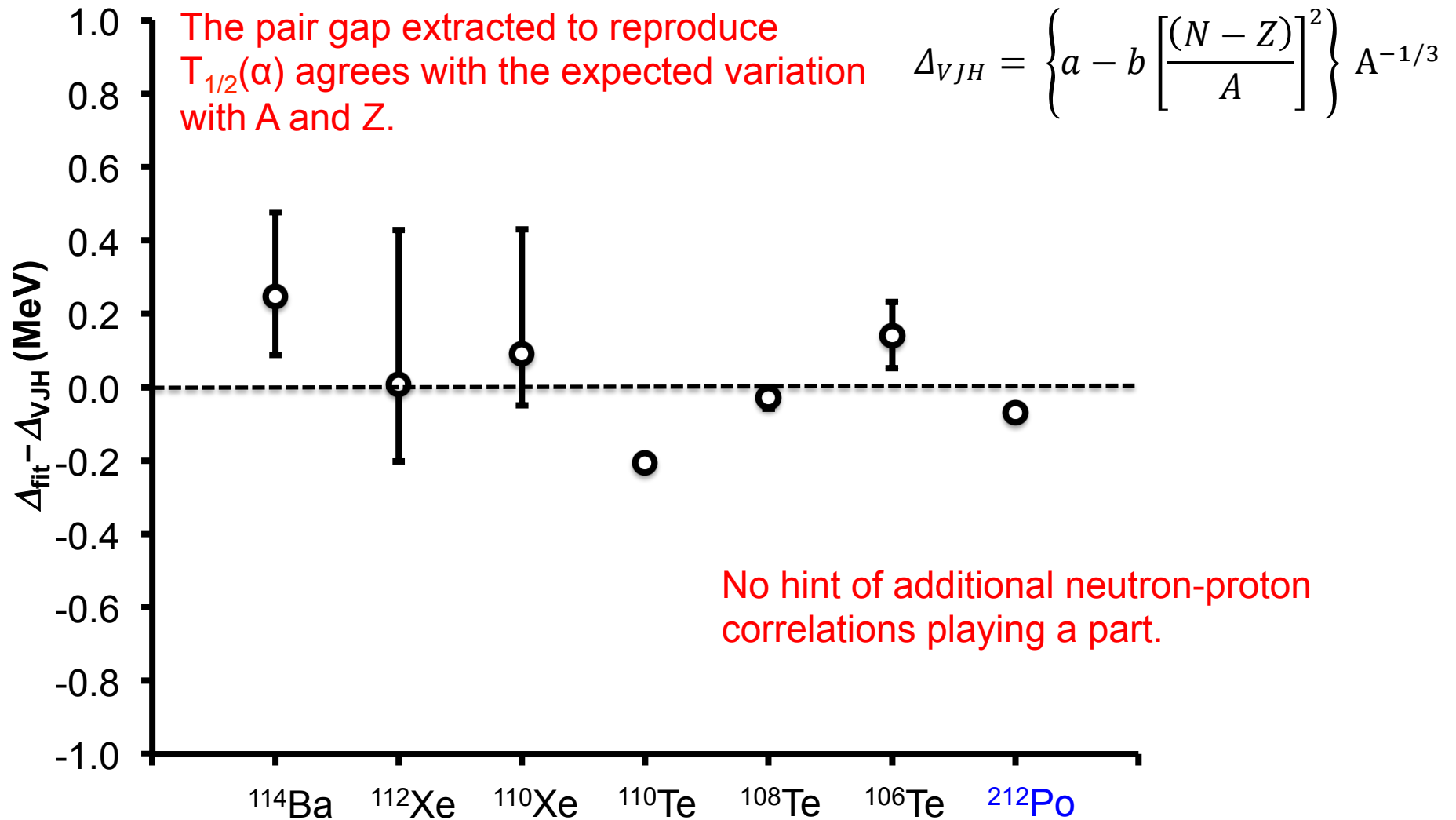


# Experiment vs. Models

| Nucleus           | $T_{1/2,\text{expt}}(\alpha)$<br>(s)        | $E_\alpha$<br>(MeV) | $\log_{10}(T_{1/2,\text{expt}})$<br>(s) | $\log_{10}(T_{1/2,\text{Royer}})$<br>(s) | $\log_{10}(T_{1/2,\text{STM}})$<br>(s) | $\Delta_{\text{fit}}$<br>(MeV) | $P$<br>( $\times 10^{-2}$ ) |
|-------------------|---|---------------------|---|--|--|--------------------------------|-----------------------------|
| $^{114}\text{Ba}$ | $42_{-18}^{+25}$<br>[6]                     | 3.480(20)<br>[6]    | $1.62_{-0.24}^{+0.21}$                  | $2.21_{-0.13}^{+0.14}$                   | $1.98_{-0.13}^{+0.14}$                 | $1.74_{-0.28}^{+0.39}$         | $1.74_{-0.89}^{+1.63}$      |
| $^{112}\text{Xe}$ | $338_{-234}^{+475}$<br>[7,8]                | 3.216(7)<br>[7]     | $2.53_{-0.51}^{+0.38}$                  | $2.74_{-0.05}^{+0.05}$                   | $2.54_{-0.05}^{+0.05}$                 | $1.49_{-0.24}^{+0.46}$         | $0.98_{-0.55}^{+1.70}$      |
| $^{110}\text{Xe}$ | $0.148_{-0.087}^{+0.090}$<br>[6]            | 3.720(20)<br>[6]    | $-0.83_{-0.38}^{+0.21}$                 | $-0.52_{-0.12}^{+0.11}$                  | $-0.71_{-0.12}^{+0.11}$                | $1.58_{-0.20}^{+0.47}$         | $1.36_{-0.60}^{+1.95}$      |
| $^{108}\text{Xe}$ | $58_{-23}^{+106} \times 10^{-6}$<br>[2]     | 4.4(2)<br>[2]       | $-4.25_{-0.21}^{+0.46}$                 | $-4.01_{-0.89}^{+0.96}$                  | $-4.14_{-0.88}^{+0.94}$                | $1.58_{-0.63}^{+1.72}$         | $1.56_{-1.43}^{+8.43}$      |
| $^{110}\text{Te}$ | $2.78(12) \times 10^6$<br>[9,10]            | 2.624(15)<br>[9]    | $6.44_{-0.02}^{+0.02}$                  | $6.23_{-0.14}^{+0.15}$                   | $6.06_{-0.14}^{+0.15}$                 | $1.27_{-0.07}^{+0.09}$         | $0.43_{-0.12}^{+0.17}$      |
| $^{108}\text{Te}$ | 4.3(4)<br>[7,8,11]                          | 3.314(4)<br>[11]    | $0.63_{-0.04}^{+0.04}$                  | $0.75_{-0.03}^{+0.03}$                   | $0.58_{-0.03}^{+0.03}$                 | $1.47_{-0.05}^{+0.04}$         | $1.00_{-0.12}^{+0.14}$      |
| $^{106}\text{Te}$ | $70_{-15}^{+20} \times 10^{-6}$<br>[6,7,12] | 4.128(9)<br>[7]     | $-4.15_{-0.11}^{+0.10}$                 | $-3.85_{-0.04}^{+0.04}$                  | $-3.97_{-0.04}^{+0.04}$                | $1.66_{-0.12}^{+0.13}$         | $1.98_{-0.46}^{+0.57}$      |
| $^{104}\text{Te}$ | $<18 \times 10^{-9}$<br>[2]                 | 4.9(2)<br>[2]       | $<-7.74$                                | $-7.10_{-0.70}^{+0.74}$                  | $-7.14_{-0.69}^{+0.76}$                | $>1.47$                        | $>1.65$                     |

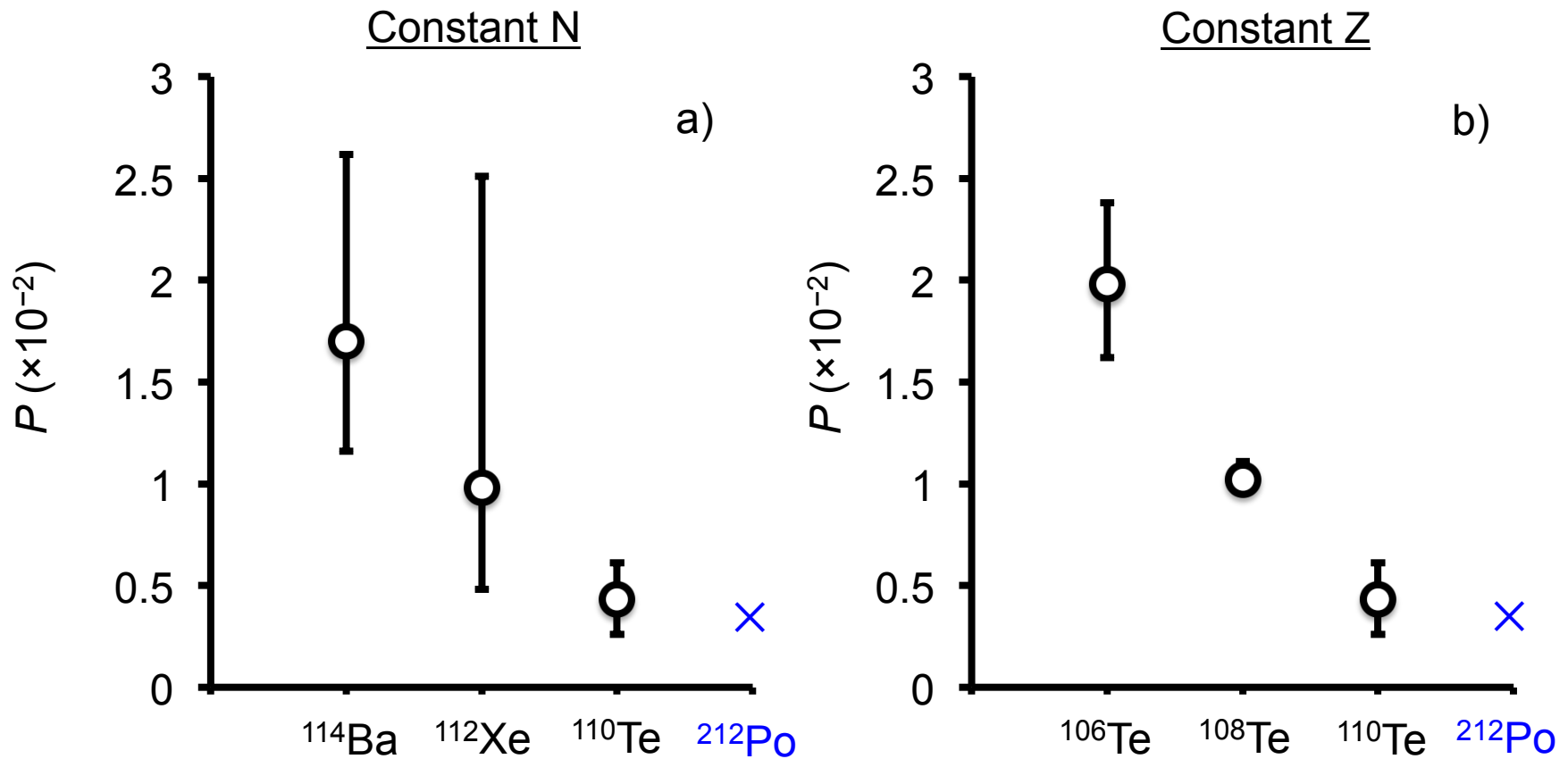
R.M. Clark et al., Phys. Rev. C 101 034313 (2020)

# STM Pair Gaps vs. Expectation



P. Vogel, B. Jonson, and P.G. Hansen, Phys. Lett. B **139**, 227 (1984).

# Extracted Alpha-Particle-Formation Probabilities, $P$



However, the “expected” variation of the pairing does lead to an increase in the formation probability of the alpha particle,  $P$ , as one approaches  $N=Z$ .

R.M. Clark et al., Phys. Rev. C 101 034313 (2020)

# Alpha Decay Chain $^{108}\text{Xe} \rightarrow ^{104}\text{Te} \rightarrow ^{100}\text{Sn}$

If extra proton-neutron correlations then the effect is likely to be biggest at  $N=Z$ . They would show up in unrealistically large values for the fitted pairing gaps.

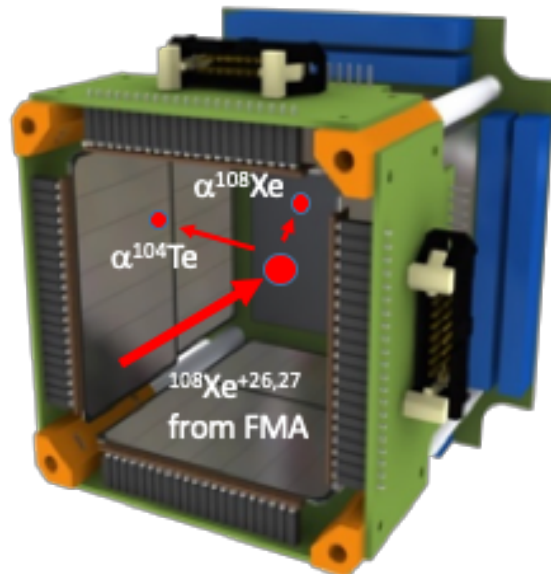
PHYSICAL REVIEW LETTERS **121**, 182501 (2018)

Editors' Suggestion

Featured in Physics

## Superallowed $\alpha$ Decay to Doubly Magic $^{100}\text{Sn}$

K. Auranen,<sup>1,\*</sup> D. Seweryniak,<sup>1</sup> M. Albers,<sup>1</sup> A. D. Ayangeakaa,<sup>1,†</sup> S. Bottoni,<sup>1,‡</sup> M. P. Carpenter,<sup>1</sup> C. J. Chiara,<sup>1,2,§</sup> P. Copp,<sup>1,3</sup> H. M. David,<sup>1,||</sup> D. T. Doherty,<sup>4,¶</sup> J. Harker,<sup>1,2</sup> C. R. Hoffman,<sup>1</sup> R. V. F. Janssens,<sup>5,6</sup> T. L. Khoo,<sup>1</sup> S. A. Kuvin,<sup>1,7</sup> T. Lauritsen,<sup>1</sup> G. Lotay,<sup>8</sup> A. M. Rogers,<sup>1,\*\*</sup> J. Sethi,<sup>1,2</sup> C. Scholey,<sup>9</sup> R. Talwar,<sup>1</sup> W. B. Walters,<sup>2</sup> P. J. Woods,<sup>4</sup> and S. Zhu<sup>1</sup>



Two events (!) with the average properties:  $^{108}\text{Xe}$

$$E_{\alpha} = 4.4(2) \text{ MeV}$$

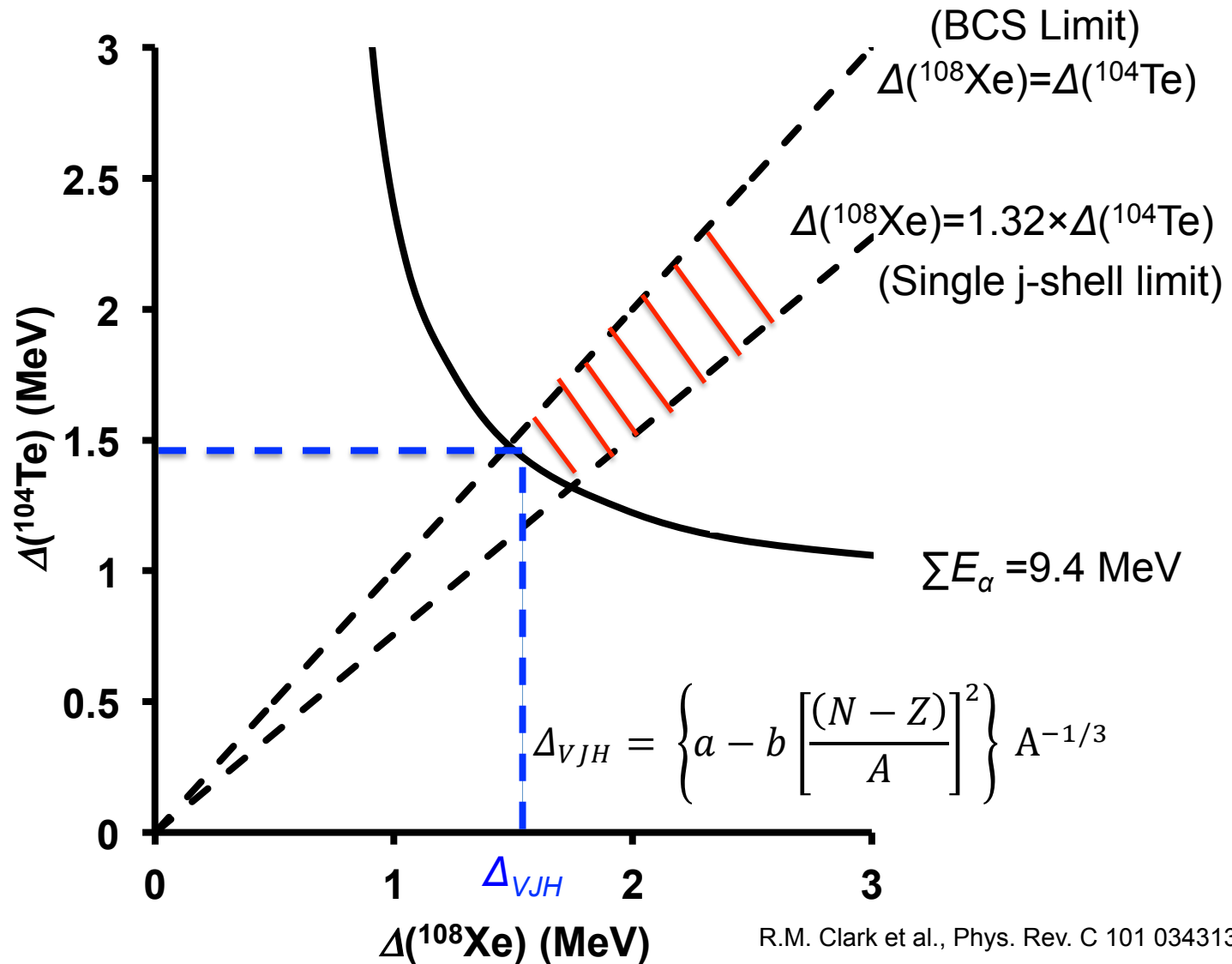
$$T_{1/2} = 58_{-23}^{+106} \mu\text{s}$$

$^{104}\text{Te}$

$$E_{\alpha} = 4.9(2) \text{ MeV}$$

$$T_{1/2} < 18 \text{ ns.}$$

# Exclusion Plot for Pairing Gaps



R.M. Clark et al., Phys. Rev. C 101 034313 (2020)

# Summary

## Superfluid Tunneling Model

- Able to reproduce known data on alpha decay
- All ingredients ( $Q_\alpha$ ,  $L$ ,  $\Delta$ ) essential to understanding

## In the $^{100}\text{Sn}$ region

- Alpha preformation is larger near  $^{100}\text{Sn}$  than  $^{208}\text{Pb}$
- Expected due to variation of pair gap
- Current data on N=Z nuclei  $^{104}\text{Te}$  and  $^{108}\text{Xe}$  still leave open possibility of superallowed alpha decay

# Future Experiments at FRIB



|              | $^{108}\text{Xe}$    |                      | $^{112}\text{Ba}$    |                      |
|--------------|----------------------|----------------------|----------------------|----------------------|
| FRIB Phase   | Day 1                | Max                  | Day 1                | Max                  |
| Primary Beam | $^{124}\text{Xe}$    | $^{124}\text{Xe}$    | $^{238}\text{U}$     | $^{238}\text{U}$     |
| Rate         | $7.6 \times 10^{-4}$ | $3.0 \times 10^{-2}$ | $1.0 \times 10^{-6}$ | $3.0 \times 10^{-3}$ |
| Per Day      | 65                   | 2600                 | 0.1                  | 260                  |

$^{108}\text{Xe}$  should be feasible on Day One,  $^{112}\text{Ba}$  comes into reach later

Thanks!



Extra Slides

# Inertia

$$\left( -\frac{\hbar^2}{2D} \frac{\partial^2}{\partial \xi^2} + V(\xi) \right) \psi_n(\xi) = E_n \psi_n(\xi)$$

Inertial mass parameter,

$$D = -\frac{\hbar^2}{2v} n^2$$

$n$  is the number of level crossings in rearrangement ( $n=4$ )

$v$  is the interaction matrix element,

$$\begin{aligned} v &\approx -G \langle \text{BCS} | P_d^\dagger | \text{BCS} \rangle \langle \text{BCS} | P_u | \text{BCS} \rangle \\ &\approx -\frac{G}{4} \langle \text{BCS} | P | \text{BCS} \rangle^2 = -\frac{1}{4} \frac{\Delta^2}{G}, \end{aligned}$$

**Inertial mass parameter,  $D$ , is dependent on pairing gap,  $\Delta$ .**

# Potential Energy

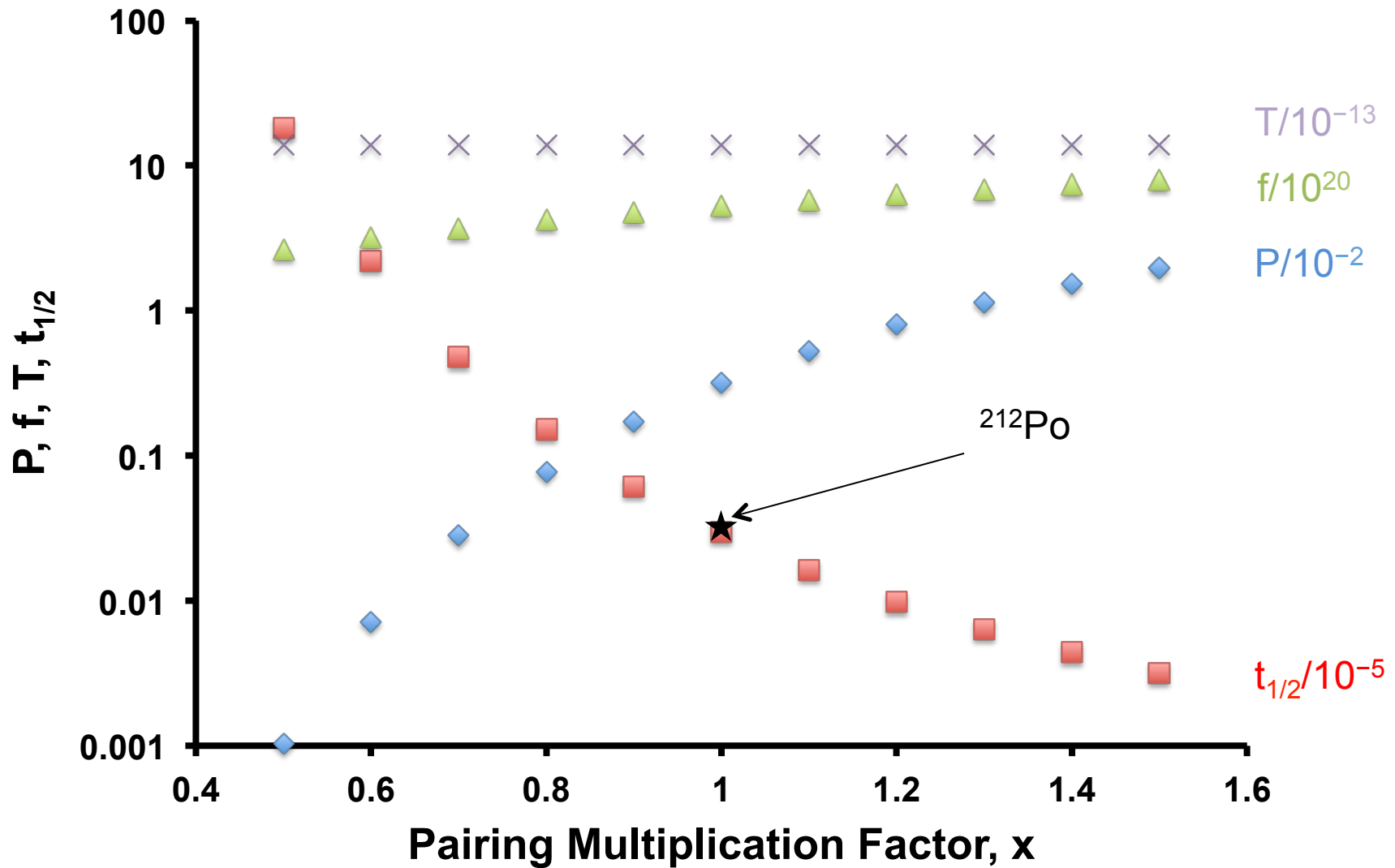
$$\left( -\frac{\hbar^2}{2D} \frac{\partial^2}{\partial \xi^2} + V(\xi) \right) \psi_n(\xi) = E_n \psi_n(\xi)$$

$$V(\xi) = \frac{1}{2} C \xi^2$$

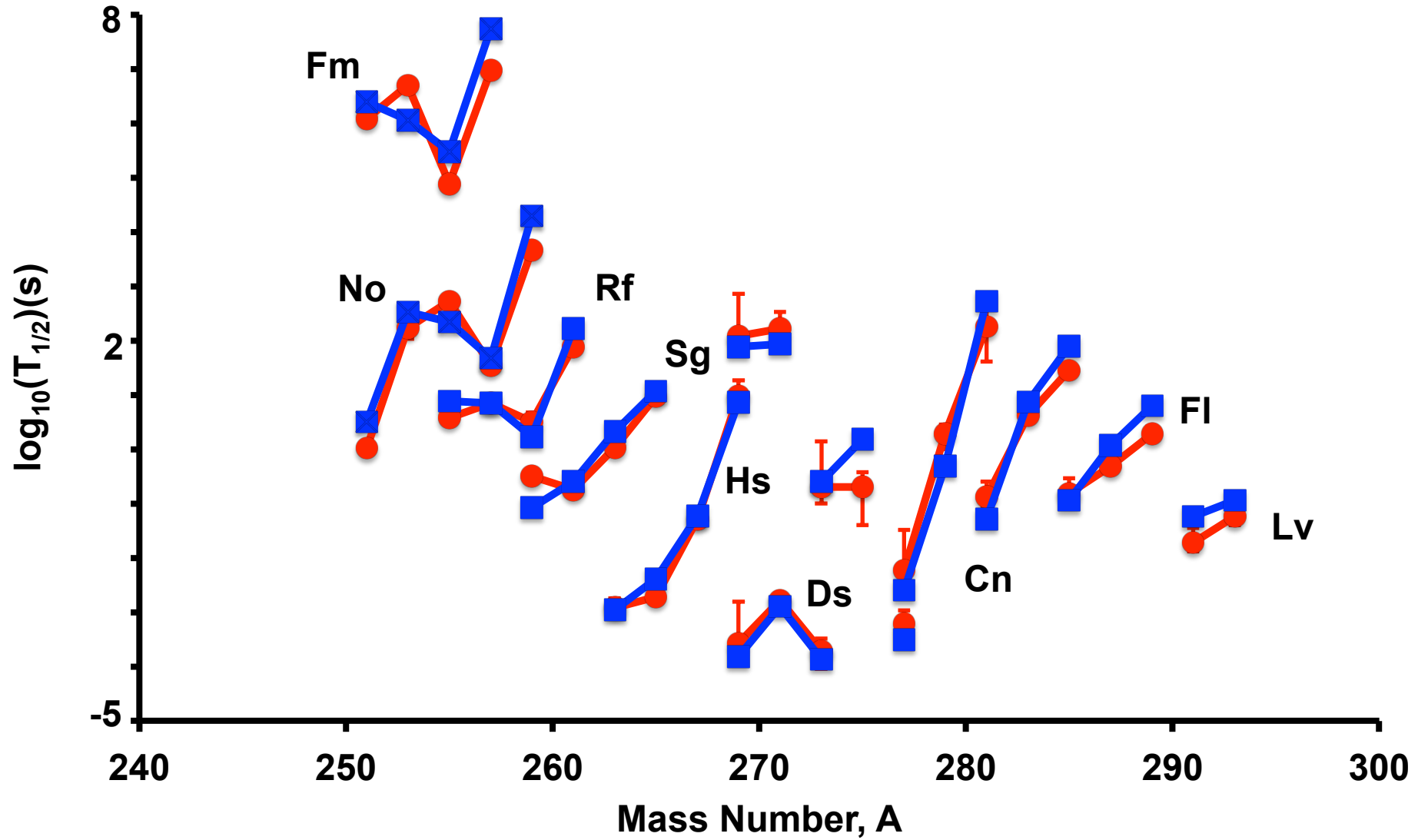
$$V(\xi=1) = V_n + V_c - Q = C/2$$

Potential energy parameter,  $C$ , is dependent on  $Q$ -value

# Dependence on Pairing Gap, $\Delta = x\Delta_0$



# Even-Z, Odd-N SHN



# Using the STM: Variation of $\Delta$ with $A$ and $Z$

Bohr and Mottelson:

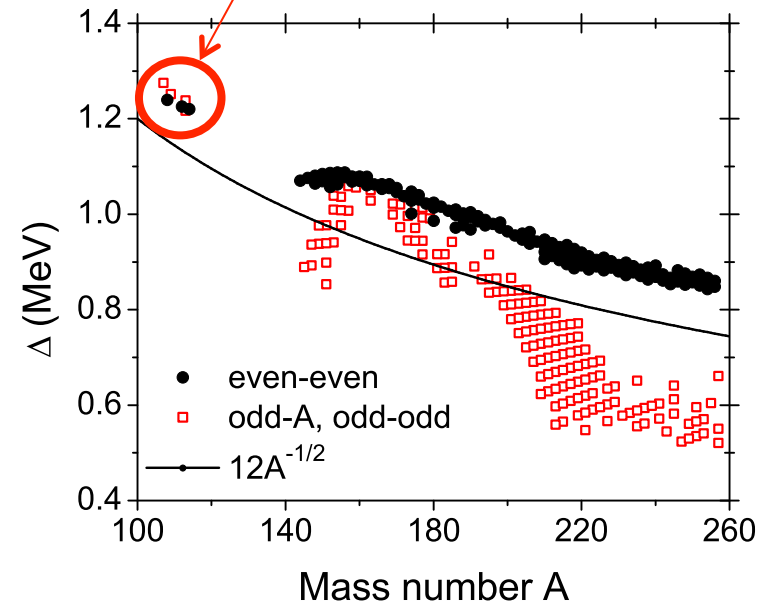
$$\Delta_{BM} = 12 A^{-1/2} \text{ MeV}$$

Vogel, Jonson, Hansen:

$$\Delta_{VJH} = \left\{ a - b \left[ \frac{(N - Z)^2}{A} \right] \right\} A^{-1/3}$$

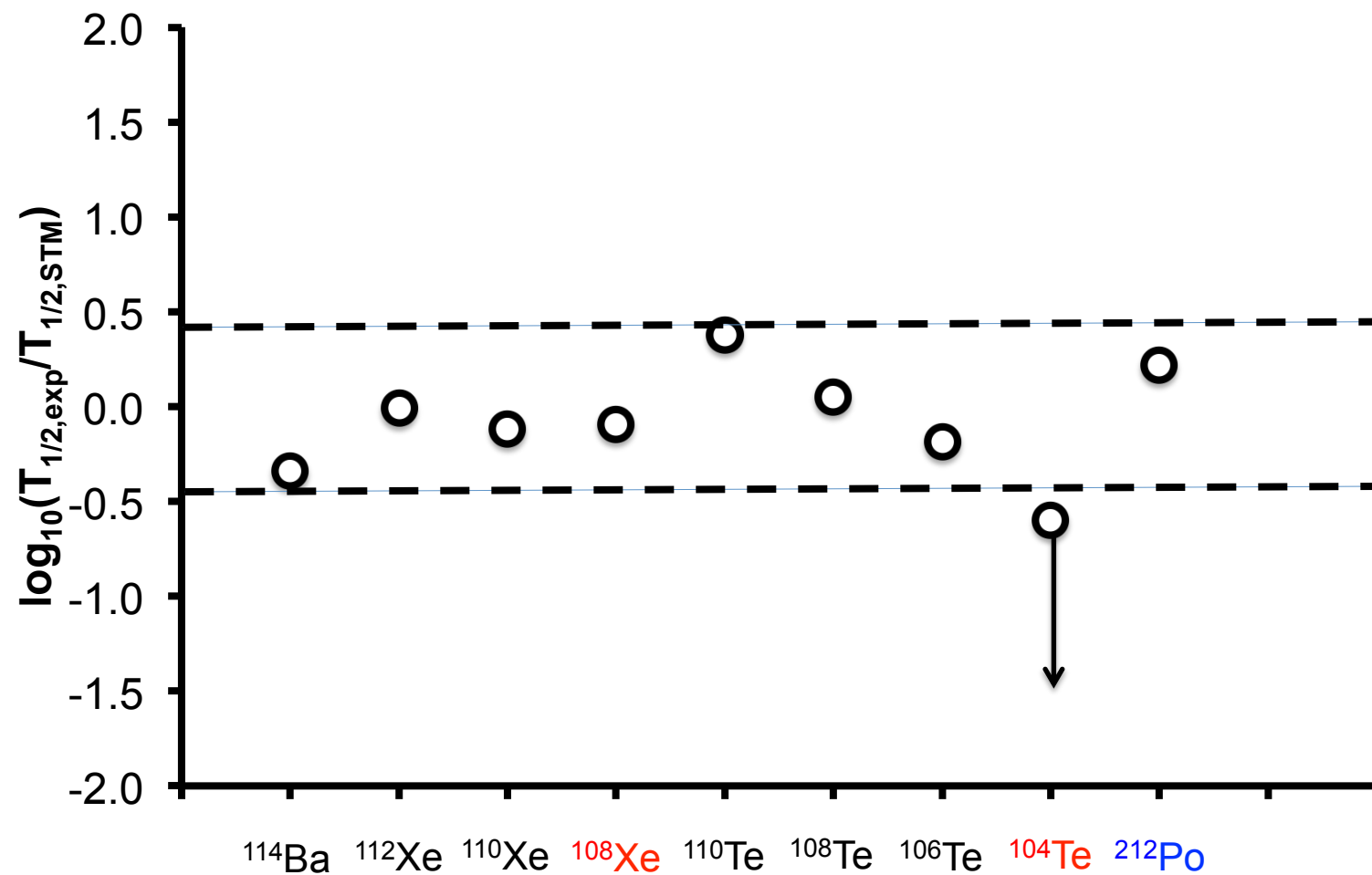
$$a = 7.2 \text{ MeV and } b = 44 \text{ MeV}$$

First hint that nothing unusual about “superallowed” decay?

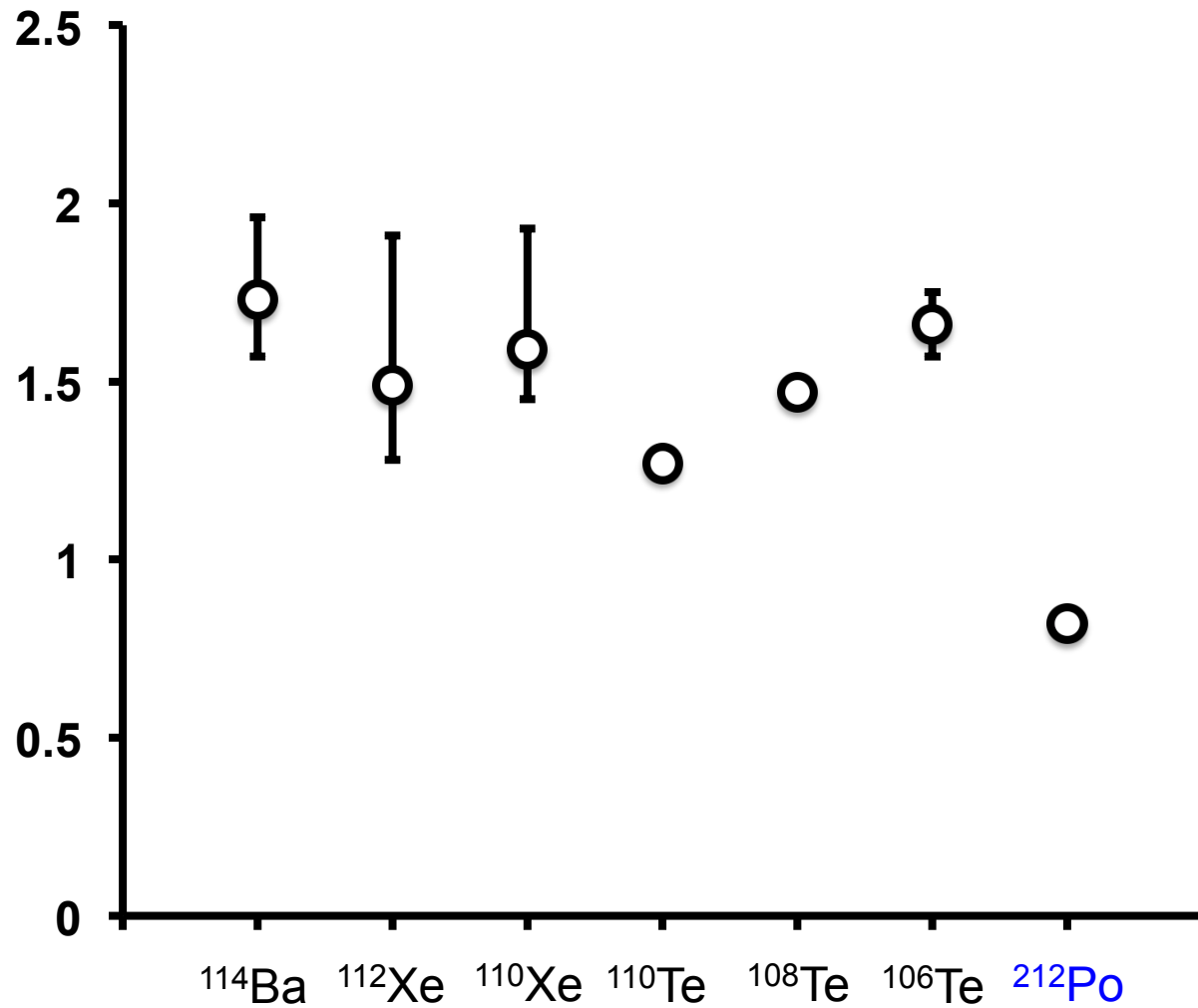


A. Bohr and B. Mottelson, *Nuclear Structure* (Benjamin, New York, 1975), Vol. 1.  
P. Vogel, B. Jonson, and P.G. Hansen, Phys. Lett. B **139**, 227 (1984).

# STM vs. Experiment



# Extracted Values of $\Delta$





# What About N=Z?

If extra proton-neutron correlations are contributing then the effect is likely to be biggest with the N=Z chain  $^{108}\text{Xe} \rightarrow ^{104}\text{Te} \rightarrow ^{100}\text{Sn}$ .

They would show up in unrealistically large values for the fitted pairing gaps.

$$^{108}\text{Xe}, E_{\alpha}=4.4(2) \text{ MeV}, T_{1/2} = 58_{-23}^{+106} \mu\text{s}$$

$$^{104}\text{Te}, E_{\alpha}=4.9(2) \text{ MeV}, T_{1/2} < 18 \text{ ns.}$$

The **sum of the alpha energy** is better constrained to value of 9.3(1)MeV

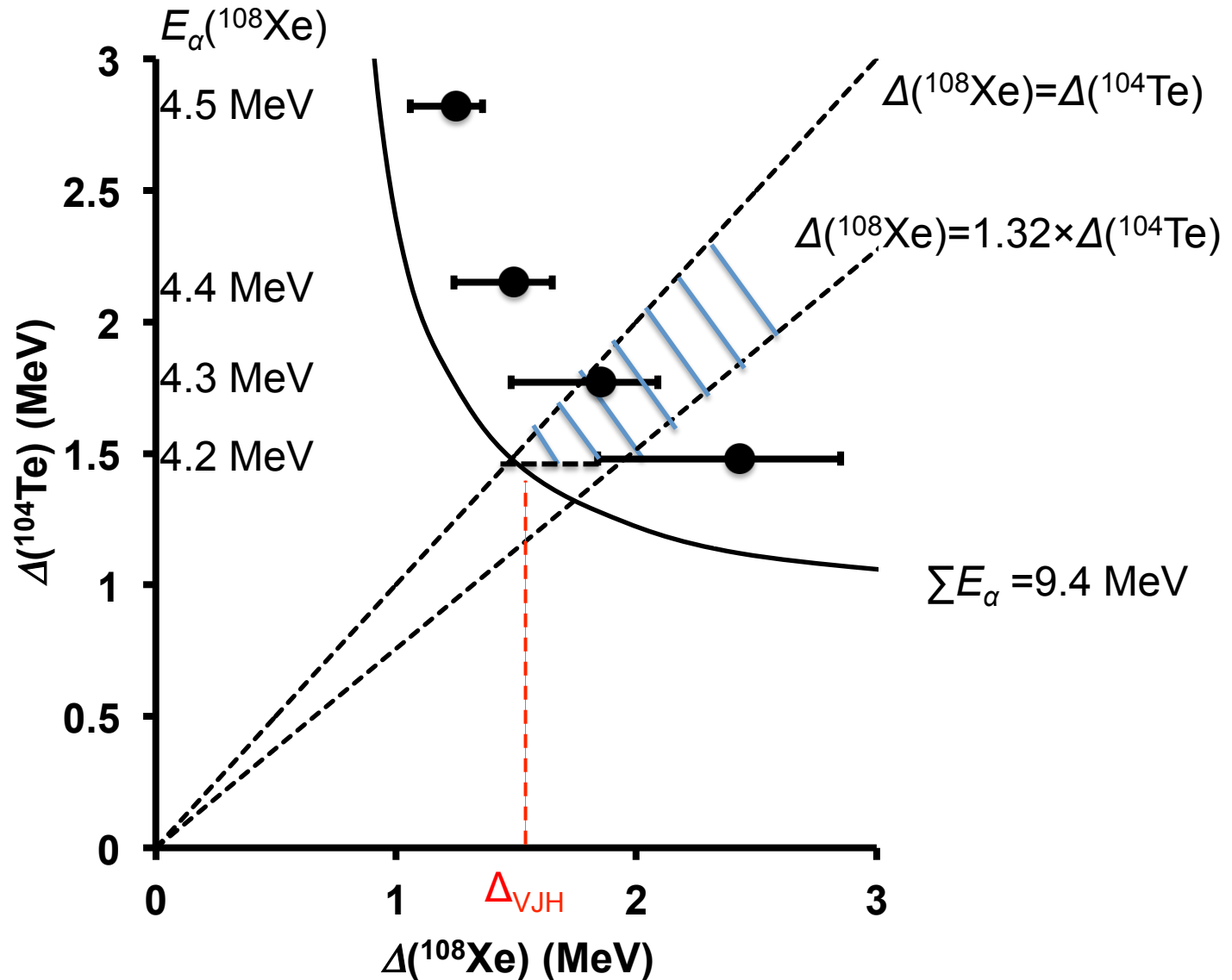
Also, assuming BCS applies,  $\Delta(^{108}\text{Xe}) > \Delta(^{104}\text{Te})$

In case of pure pairing in single-j shell:

$$\Delta = G \sqrt{\frac{n}{2} \left( \Omega - \frac{n}{2} \right)}$$

Implying maximum of  $\Delta(^{108}\text{Xe}) = 1.32 \times \Delta(^{104}\text{Te})$

# Exclusion Plot for Pairing Gaps



# New Results from JAEA

## Search for $\alpha$ decay of $^{104}\text{Te}$ with a novel recoil-decay scintillation detector

Phys. Rev. C, Y. Xiao et al., **Accepted 16 August 2019**

### ABSTRACT

A search for super-allowed  $\alpha$  decay of  $N=Z$  nuclei  $^{104}\text{Te}$  and  $^{108}\text{Xe}$  was carried out using novel recoil-decay scintillator detector at the tandem accelerator facility at Japan Atomic Energy Agency (JAEA). Inorganic crystal scintillation material of YAP:Ce (Yttrium Aluminium Perovskite) coupled to position-sensitive photo-multiplier tube (PSPMT) was implemented for the first time in radioactive decay experiment. Residues from the fusion-evaporation reaction  $^{58}\text{Ni}+^{54}\text{Fe}\rightarrow^{112}\text{Xe}^*$  were separated by the JAEA Recoil Mass Separator (RMS) and implanted into the YAP:Ce crystal.  $\alpha$  decays of neutron-deficient tellurium isotopes were identified and proton-emission of  $^{109}\text{I}$  was observed. The  $\alpha$  decay chain  $^{109}\text{Xe}\rightarrow^{105}\text{Te}\rightarrow^{101}\text{Sn}$  was recorded with time interval of 960 ns between two  $\alpha$  pulses. Position localization in the crystal for decays and ions in the energy range from hundreds keV to 60 MeV was achieved with the accuracy of 0.67 mm, proving that this detector is capable of making temporal and spatial correlations for fast decay events. No conclusive evidence was found for the decay chain  $^{108}\text{Xe}\rightarrow^{104}\text{Te}\rightarrow^{100}\text{Sn}$  within 3 days experiment. However, **two events were observed with properties consistent with the reported observation at Fragment Mass Analyzer (FMA) but with the separation between signals less than 4 ns**. The cross section limit of 130 pb was obtained for production of two events of  $^{108}\text{Xe}$ , about an order of magnitude below the expectation based on earlier cross section measurements and HIVAP fusion-evaporation code.

$$T_{1/2}(^{104}\text{Te}) < 4\text{ns?}$$

# Exclusion Plot for Pairing Gaps

