

Quantum simulation of open quantum systems in heavy-ion collisions

James Mulligan, Felix Ringer

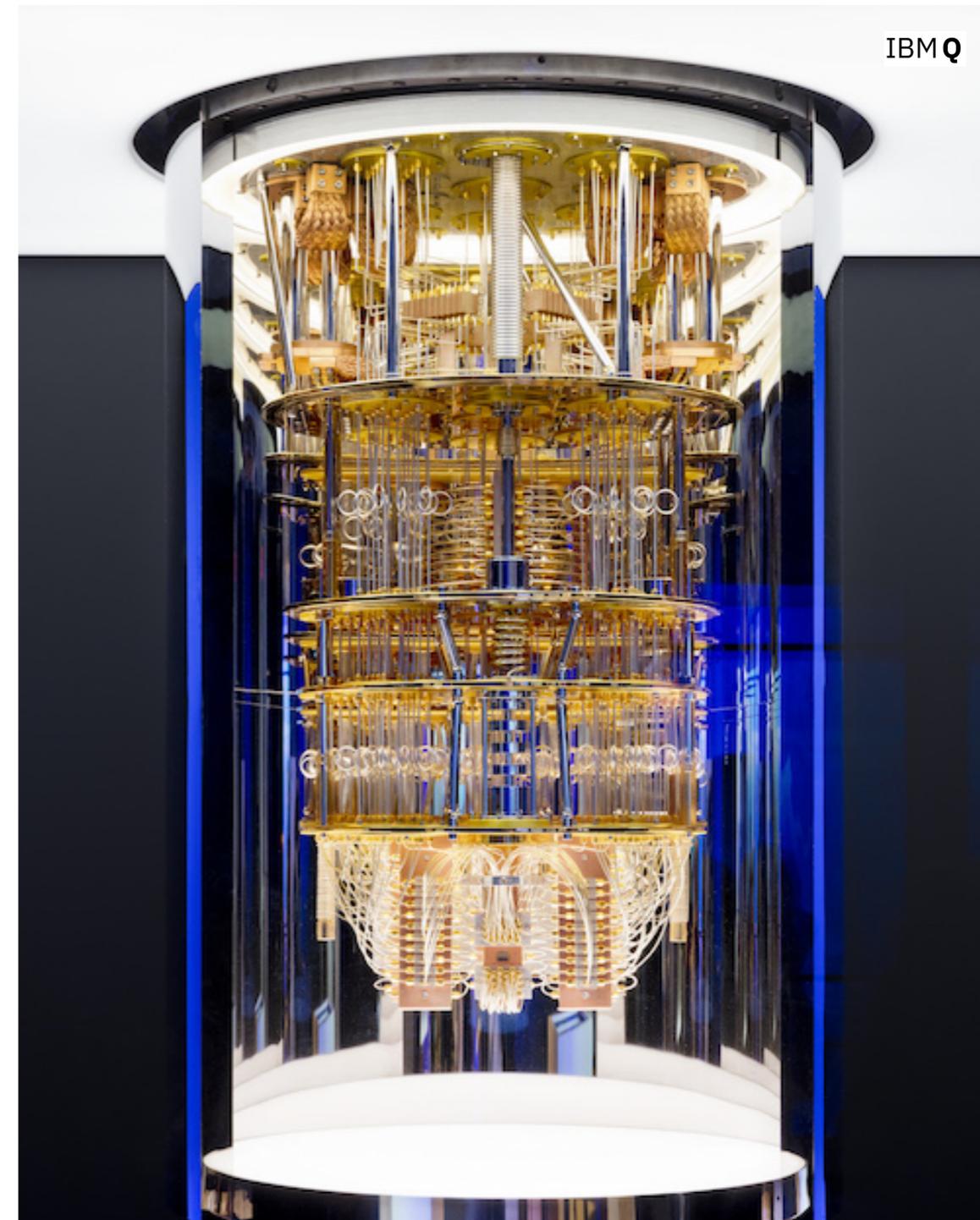
Lawrence Berkeley National Laboratory

arXiv: 2010.03571

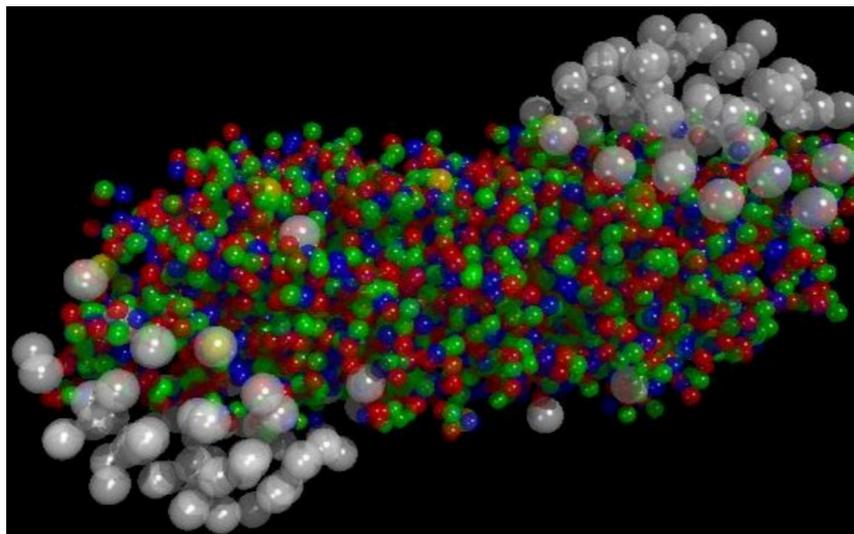
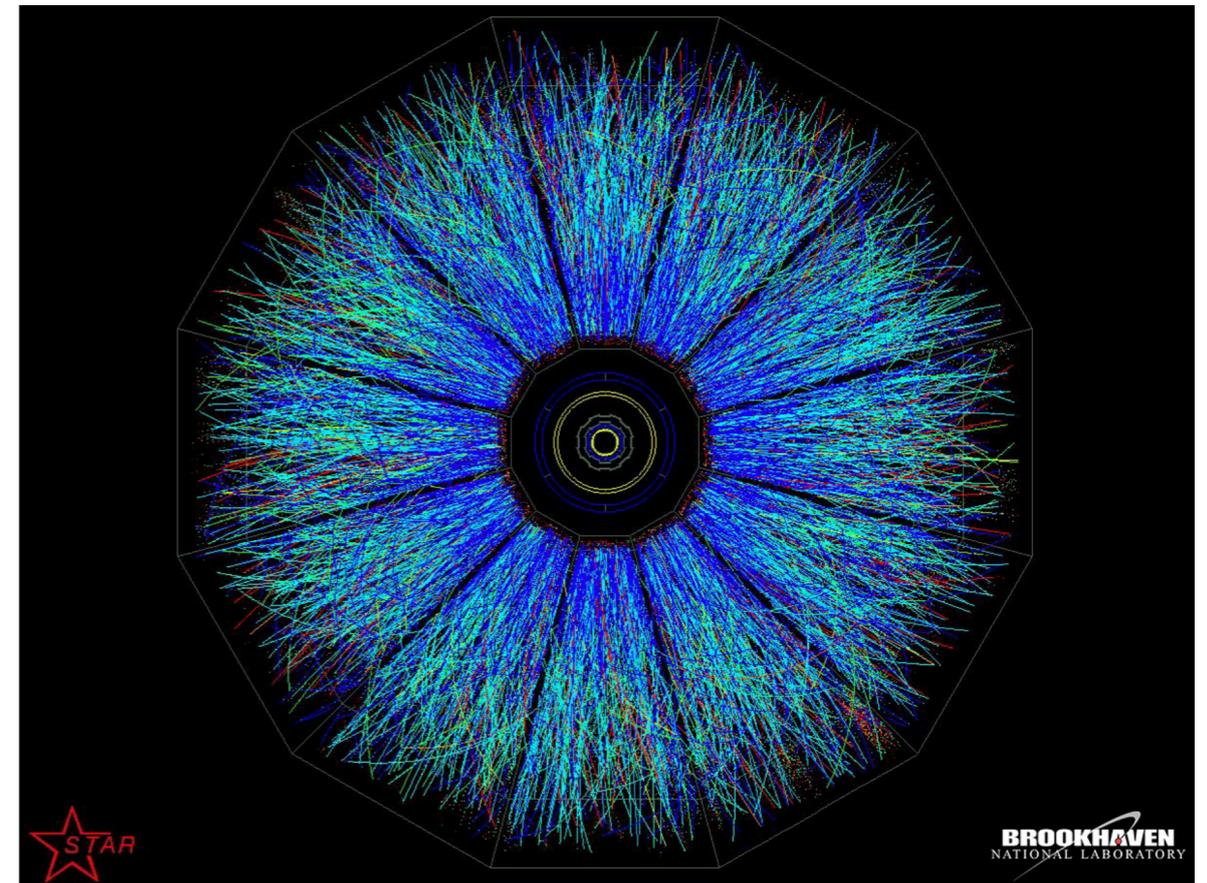
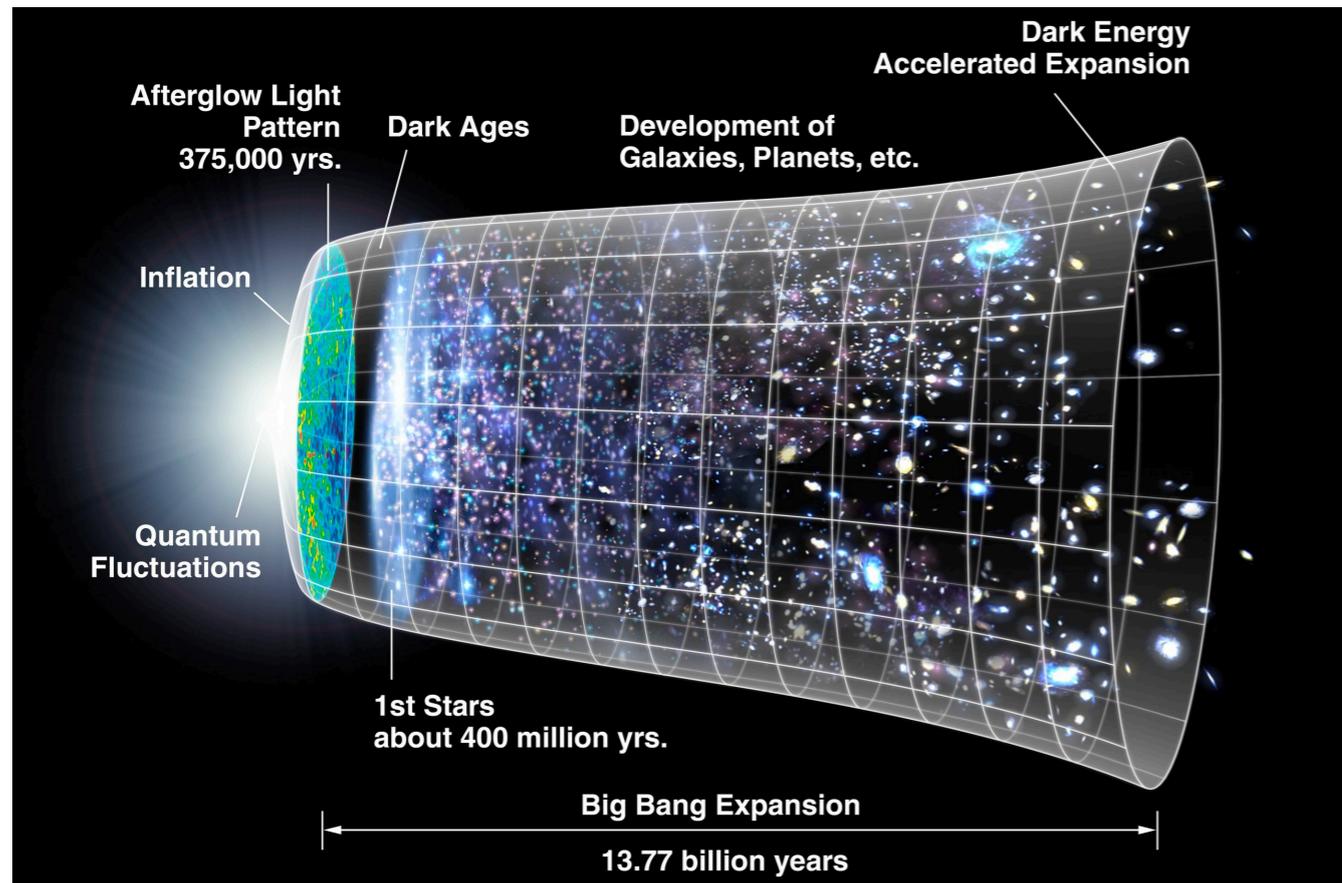
<i>Wibe de Jong</i>	}	<i>LBL CRD</i>
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<i>Mateusz Ploskon</i>		
<i>Felix Ringer</i>		
<i>Xiaojun Yao</i>	}	<i>MIT</i>



LBNL NSD staff meeting
10/20/2020



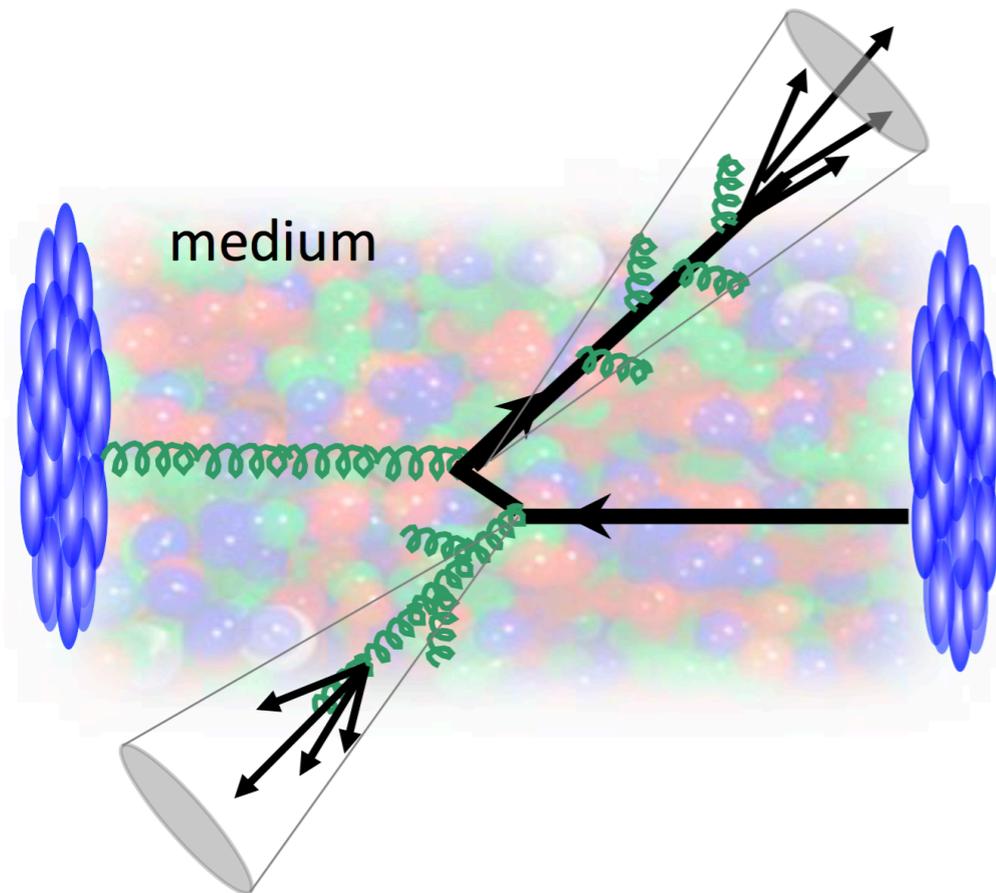
Heavy-ion collisions — the Quark Gluon Plasma



Described by Quantum chromodynamics (QCD)

Unbound quarks and gluons

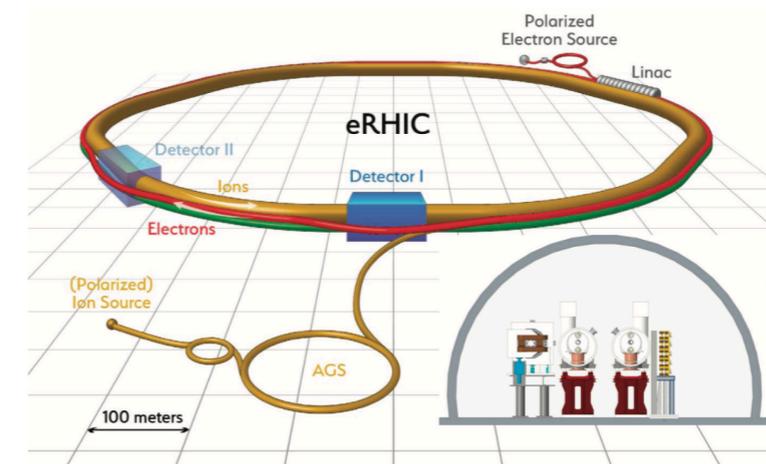
Hard probes of the Quark Gluon Plasma



LHC, RHIC



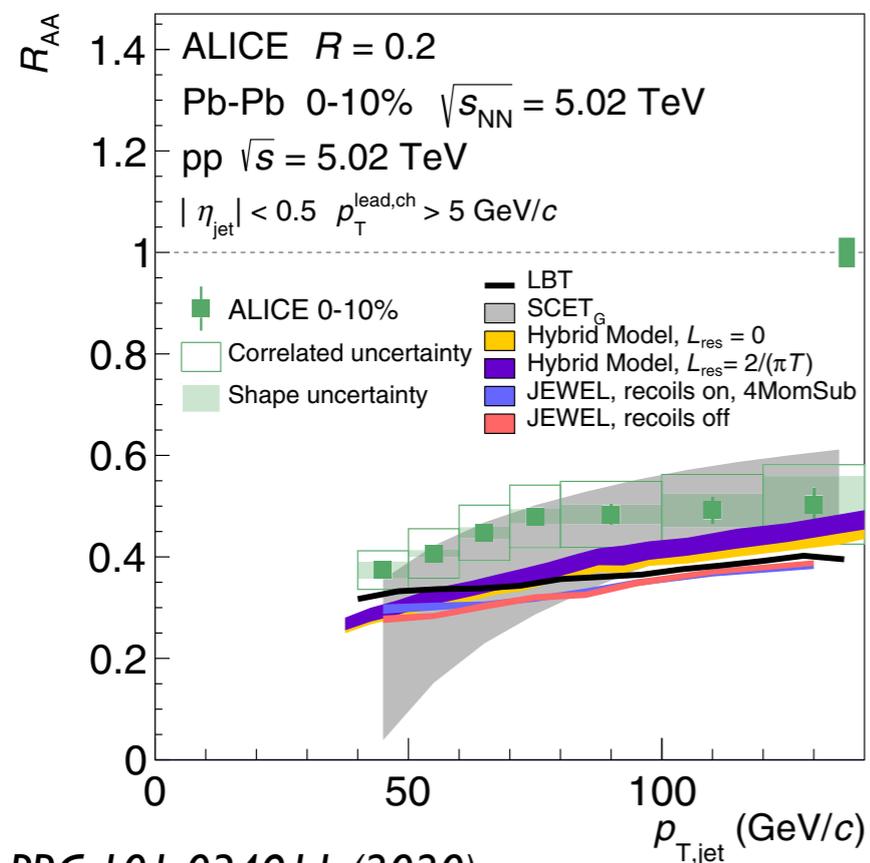
EIC - cold nuclear matter



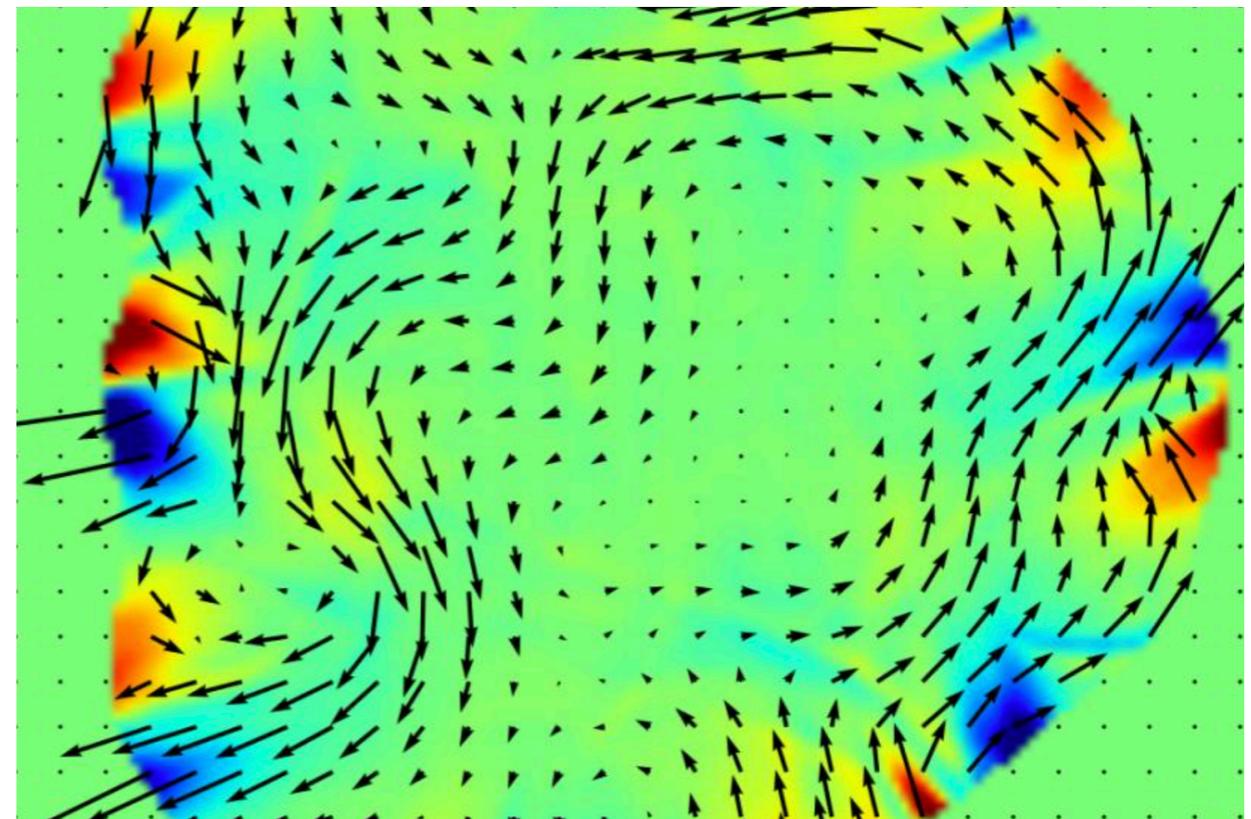
- Highly energetic particles and jets

Hard probes of the Quark Gluon Plasma

$$R_{AA} = \frac{d\sigma^{\text{PbPb}}}{\langle T_{AA} \rangle d\sigma^{pp}}$$



PRC 101 034911 (2020)



X.N.Wang et al.

- Study the **real-time** dynamics of hard probes
- Combine with hydrodynamic models of the QGP
- Study the microscopic structure of the medium

Outline

Open quantum
systems in heavy-ion
collisions

Quantum simulation
with IBM Q

Open quantum systems and the nuclear medium

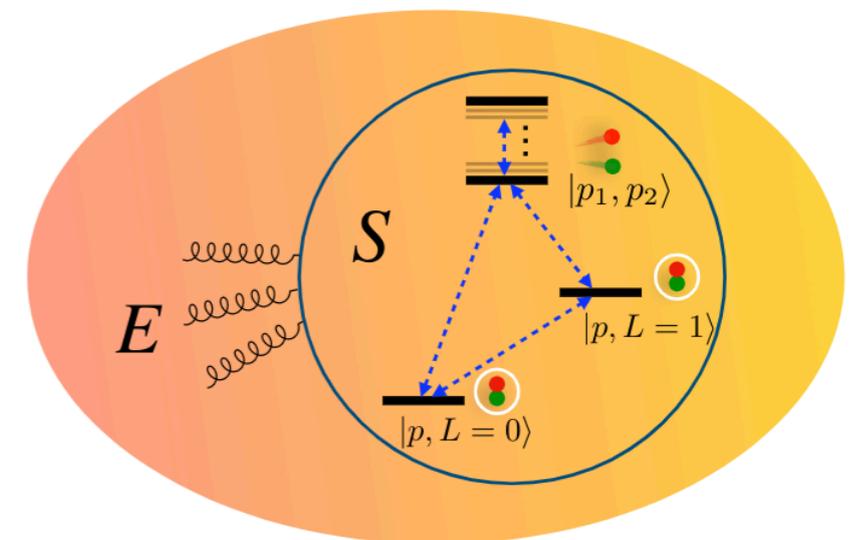
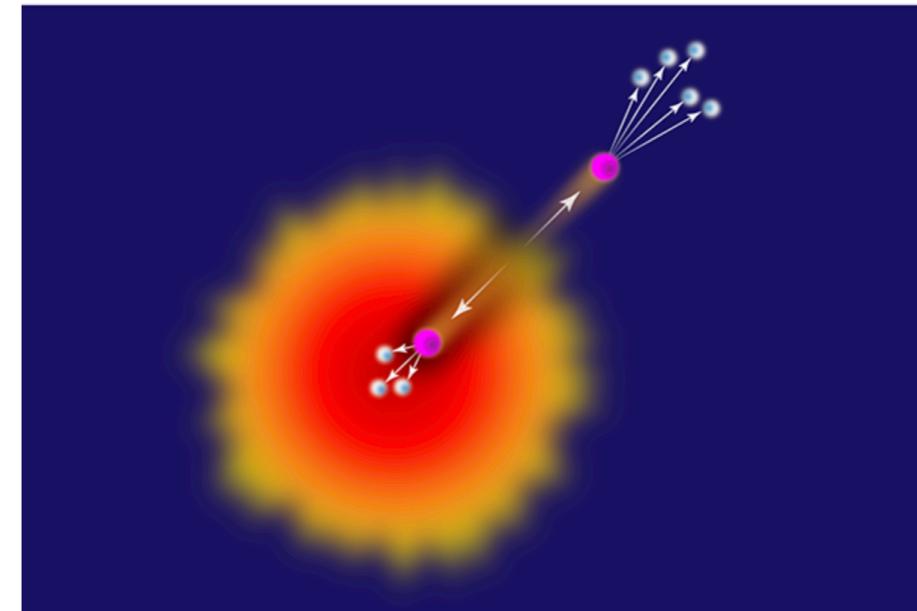
- Study the real time dynamics of the quantum evolution of probes in the nuclear medium (LHC/RHIC/EIC)
- **System** - Jet/heavy-flavor
- **Environment** - Nuclear matter

$$H(t) = H_S(t) + H_E(t) + H_I(t)$$

- The time evolution is governed by the von Neumann equation

$$\frac{d}{dt}\rho^{(\text{int})}(t) = -i \left[H_I^{(\text{int})}(t), \rho^{(\text{int})}(t) \right]$$

Hamiltonian formulation of QCD



Akamatsu, Rothkopf `12-`20, Müller et al `18, Mehen, Yao `18, Qiu, Ringer, Sato, Zurita `19, Vaidya, Yao `20

Open quantum systems and the nuclear medium

- Study the real time dynamics of the quantum evolution of probes in the nuclear medium (LHC/RHIC/EIC)
- **System** - Jet/heavy-flavor
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$$H(t) = H_S(t) + H_E(t) + H_I(t)$$

- Lindblad equation in the Markovian limit

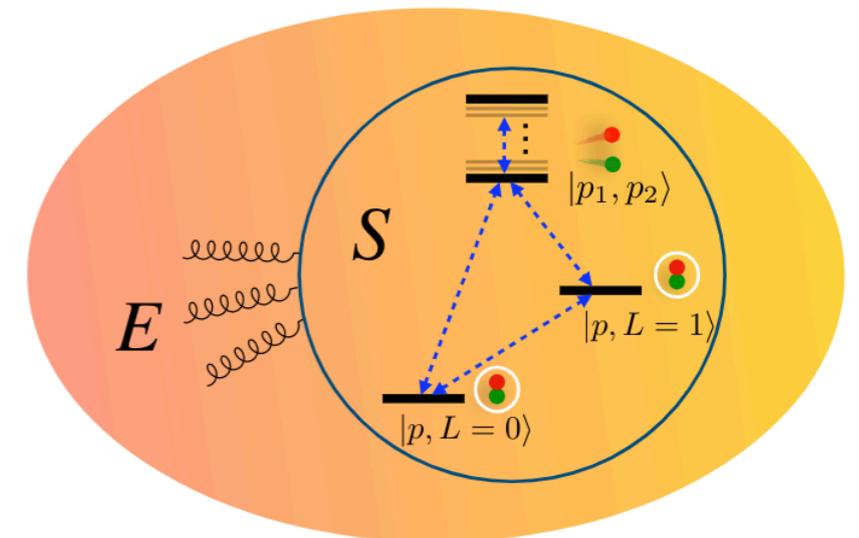
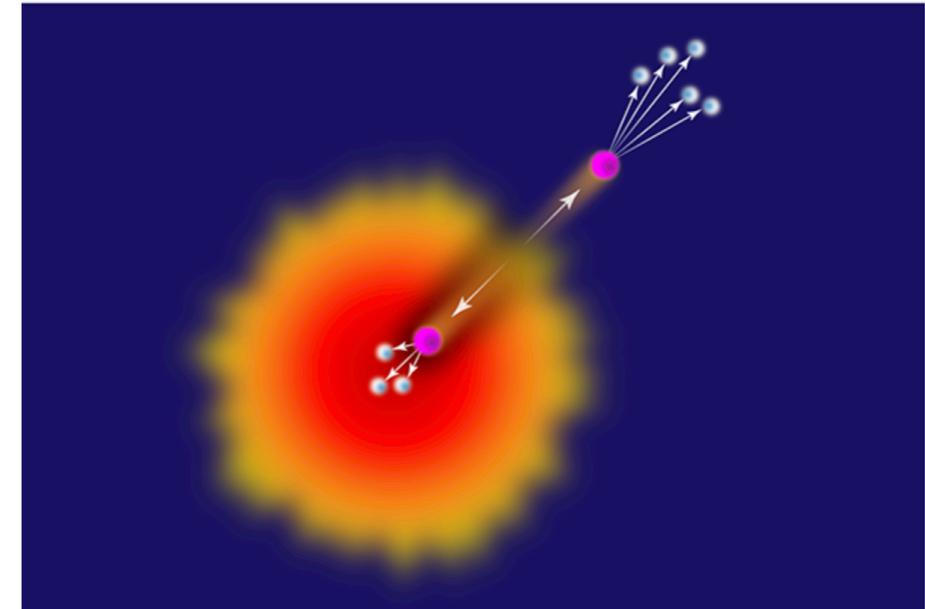
$$\rho_S = \text{tr}_E[\rho]$$

$$\frac{d}{dt}\rho_S = -i[H_S, \rho_S] + \sum_{j=1}^m \left(L_j \rho_S L_j^\dagger - \frac{1}{2} L_j^\dagger L_j \rho_S - \frac{1}{2} \rho_S L_j^\dagger L_j \right)$$

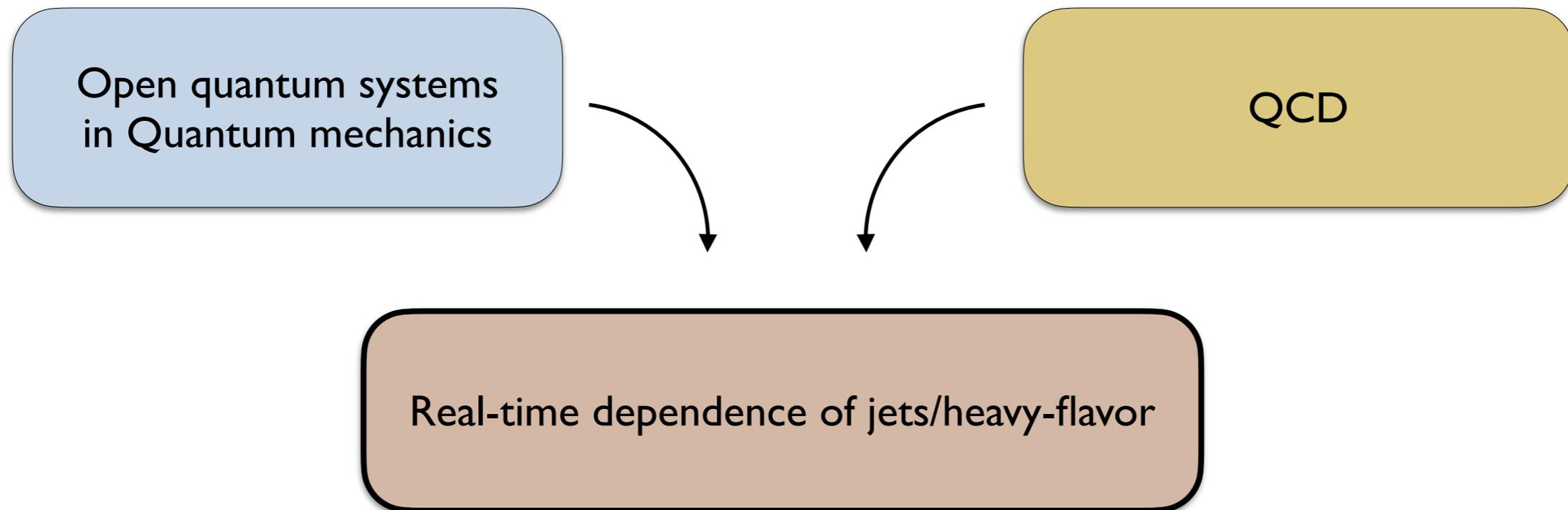
- See also e.g. non-global logs and CGC

Neill `15, Armesto et al. `19, Li, Kovner `20

Akamatsu, Rothkopf `12-`20, Müller et al `18, Mehen, Yao `18, Qiu, Ringer, Sato, Zurita `19, Vaidya, Yao `20



Open quantum systems and the nuclear medium



- Currently various approximations are considered *Blaizot, Escobedo '18, Yao, Mehen '18, '20*

- Markovian limit
- Small coupling of system and environment
- Semi-classical transport

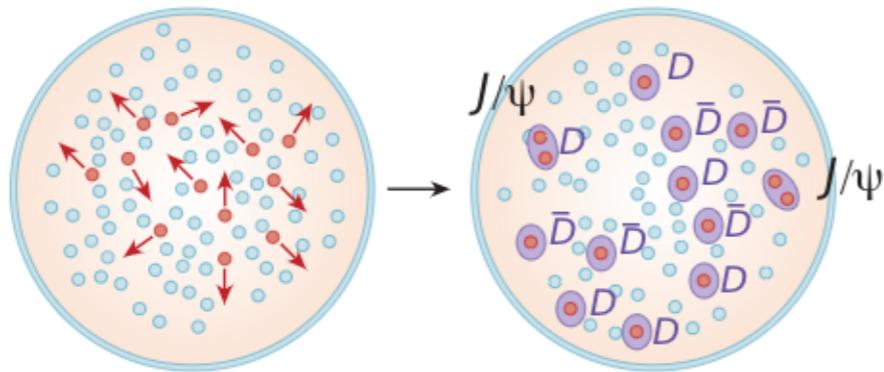
*Akamatsu, Rothkopf et al. '12-'20, Brambilla et al. '17-'20
Yao, Mueller, Mehen '18-'20, Sharma, Tiwari '20*

Yao, Vaidya '19, Vaidya '20

Quarkonium suppression

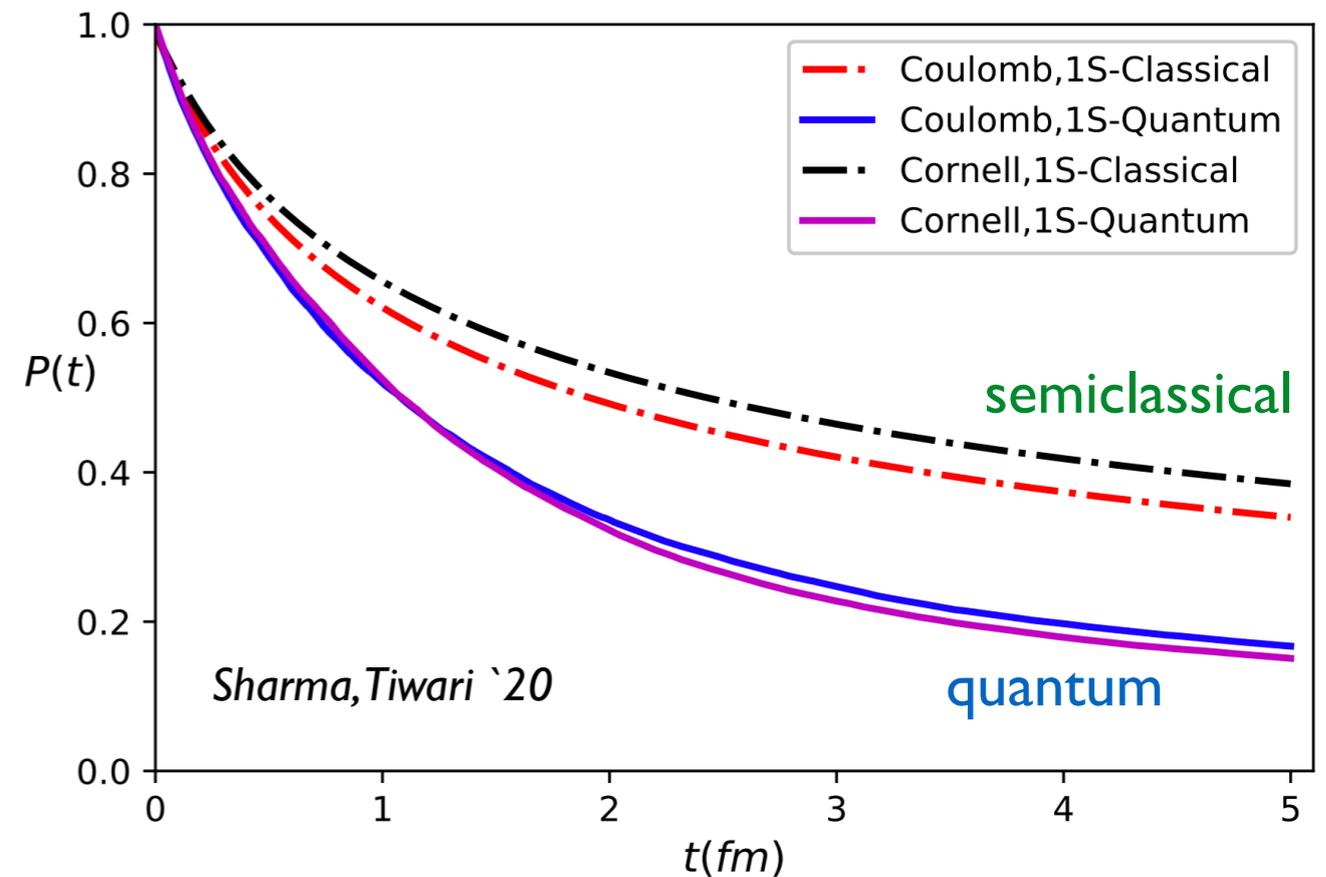
Akamatsu, Rothkopf et al. '12-'20, Brambilla et al. '17-'20
Yao, Mueller, Mehen '18-'20, Sharma, Tiwari '20

- Quarkonium production



- NRQCD + semiclassical approach compared to full quantum evolution

Survival probability of the vacuum state



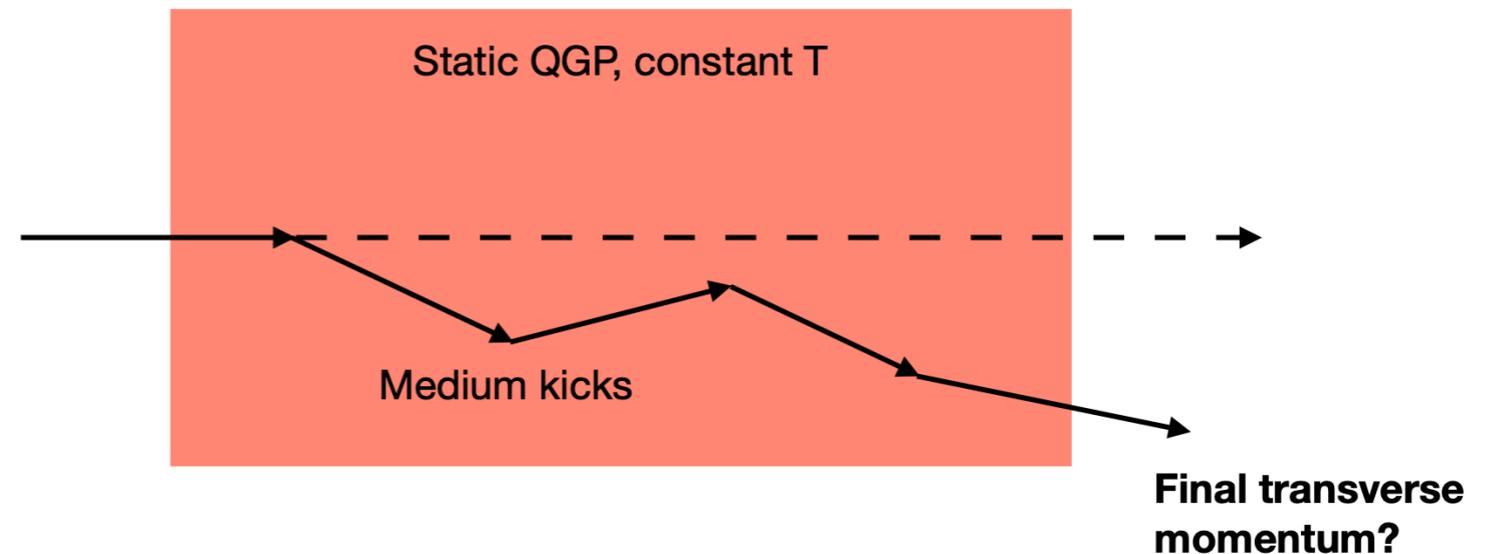
Bjorken expanding QGP $T_0 = 475$ MeV

Jet broadening

Yao, Vaidya '19

- First steps in the direction of jet physics
- Soft and collinear modes
- Forward scattering/Glauber gluons

$$H_I = \sum_{\alpha} O_{\alpha}^{(S)} \otimes O_{\alpha}^{(E)}$$



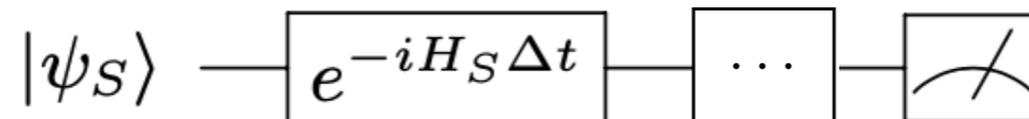
- Markovian master equation resums large logarithms $P(Q, t) = \langle Q | \rho_S(t) | Q \rangle$

Schematically
$$\partial_t P(Q, t) = -R(Q)P(Q, t) + \int \widetilde{d}q K(Q, q)P(q, t)$$

Closed quantum systems

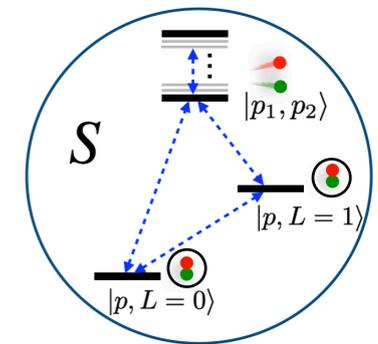
- **Time evolution of closed systems**

- Quantum simulation of the Schrödinger equation



Evolution in time steps $\Delta t = t/N_{\text{cycle}}$

- The evolution is unitary and time reversible

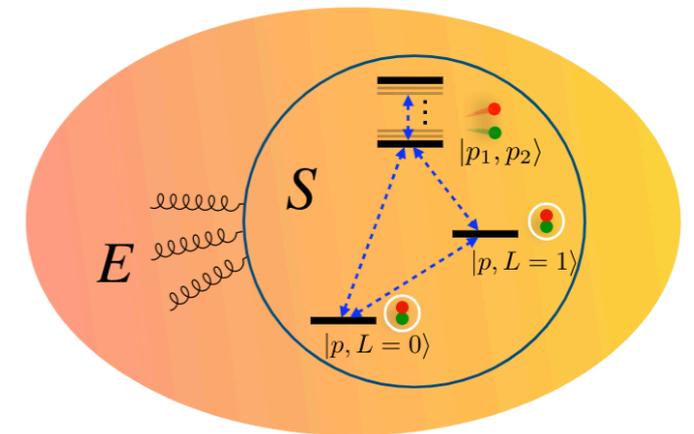
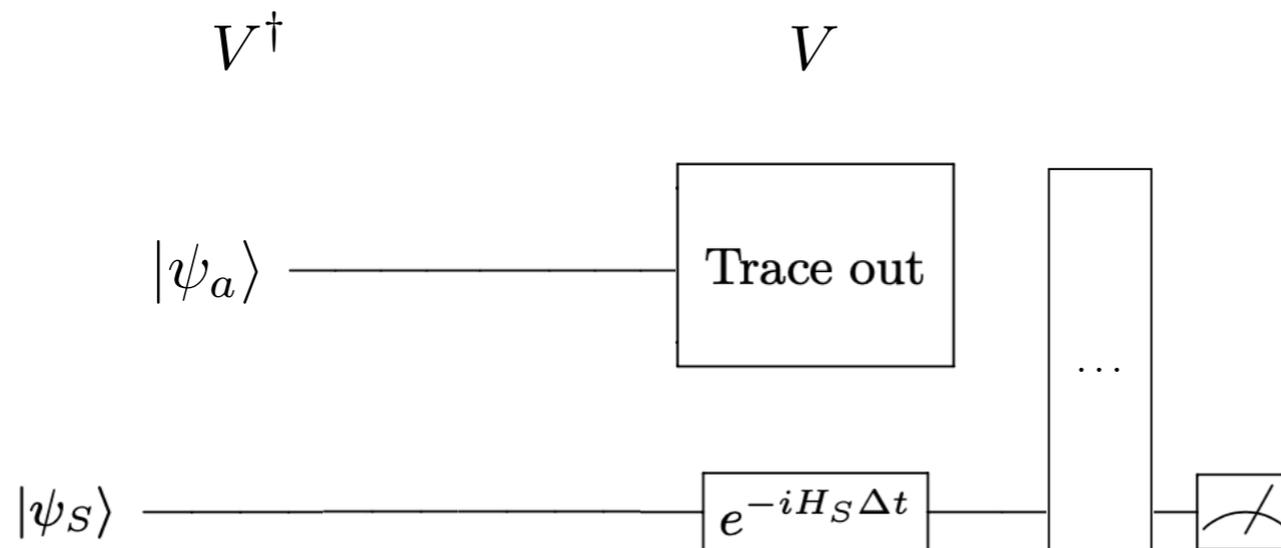


→ For open quantum systems we need to introduce a non-unitarity part

Non-unitarity and time irreversible evolution

$$\frac{d}{dt}\rho_S = -i[H_S, \rho_S] + \sum_{j=1}^m \left(L_j \rho_S L_j^\dagger - \frac{1}{2} L_j^\dagger L_j \rho_S - \frac{1}{2} \rho_S L_j^\dagger L_j \right)$$

- The Stinespring dilation theorem**



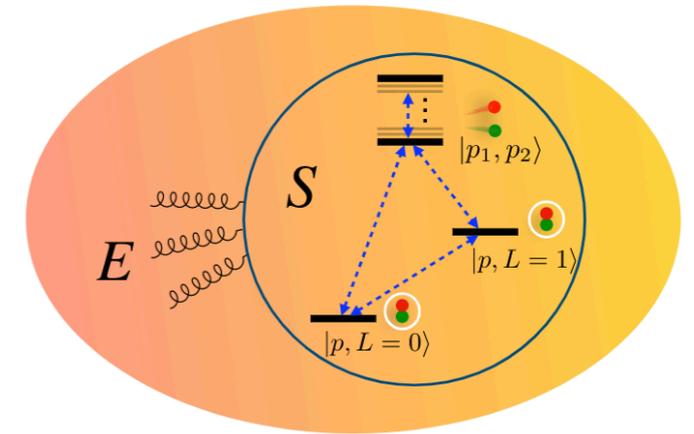
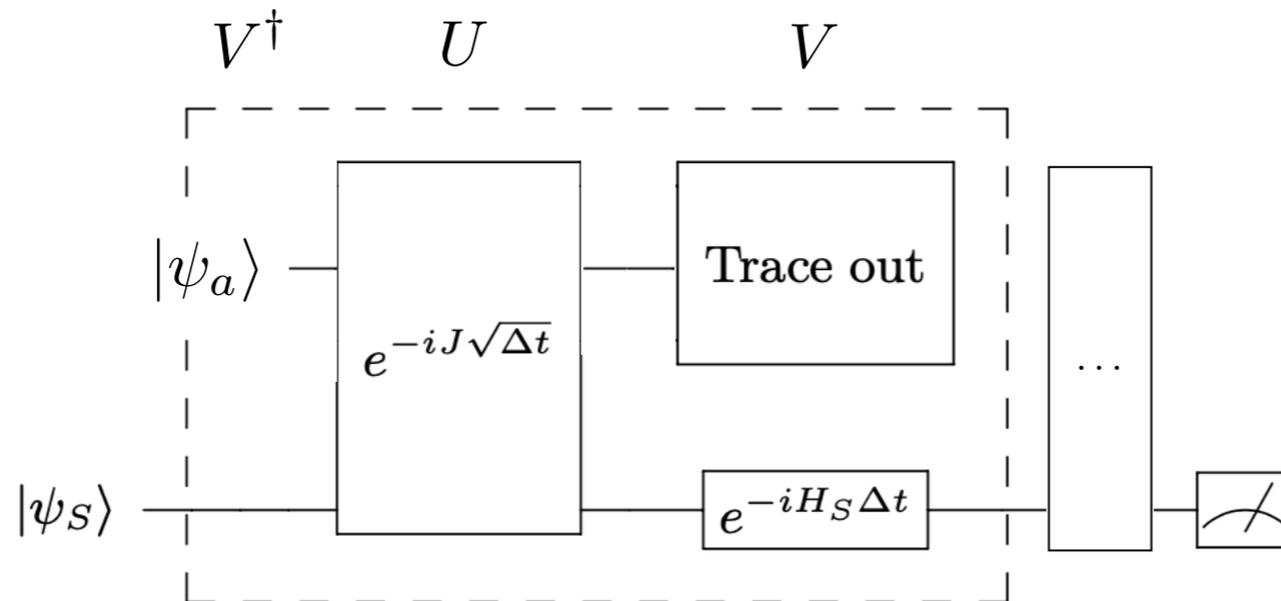
- Introducing and tracing out an ancillary system is not a unitary operation (operator theory)

$$V^\dagger V = 1 \quad VV^\dagger \neq 1$$

Non-unitarity and time irreversible evolution

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• The Stinespring dilation theorem



- Introducing and tracing out an ancillary system is not a unitary operation (operator theory)

$$V^\dagger V = 1 \quad VV^\dagger \neq 1$$

- Sandwich in between a unitary evolution step
- Evolve in time steps $\Delta t = t/N_{\text{cycle}}$

$$J = \begin{pmatrix} 0 & L_1^\dagger & \dots & L_m^\dagger \\ L_1 & 0 & \dots & 0 \\ \vdots & \vdots & \ddots & 0 \\ L_m & 0 & \dots & 0 \end{pmatrix}$$

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with IBM Q

Quantum computing

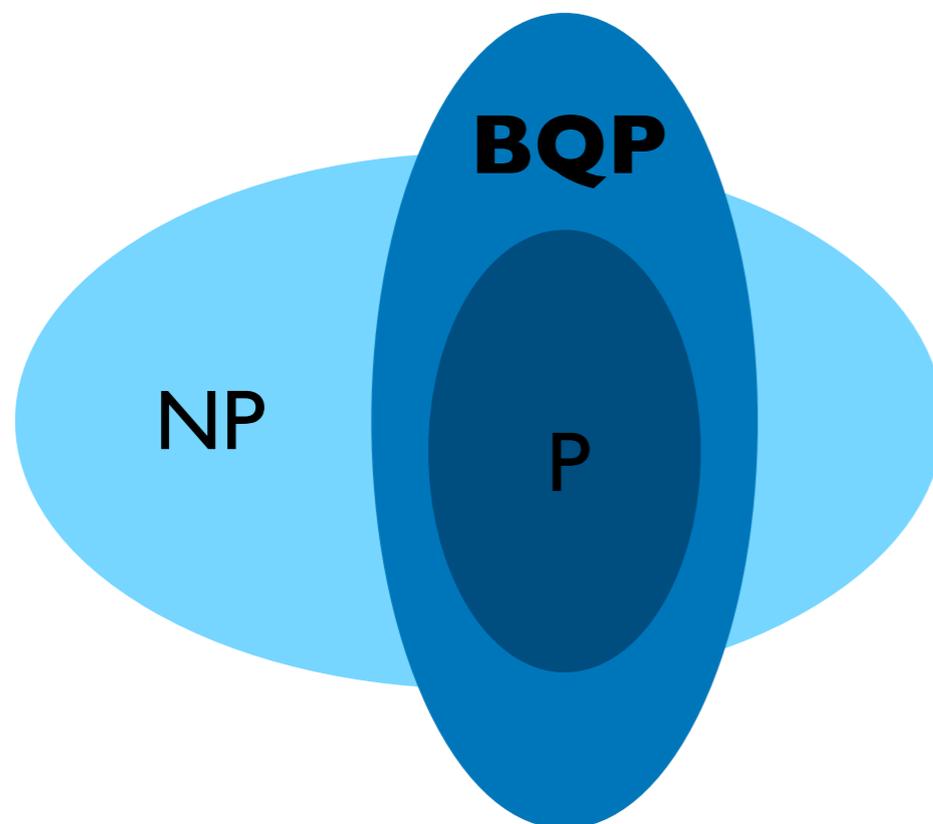
Superposition and entanglement

$$|\psi\rangle = \sum_{i=1}^{2^N} a_i |\psi_i\rangle$$

For N qubits, there are 2^N amplitudes

e.g. $|\psi\rangle = a_1|000\rangle + a_2|001\rangle + a_3|010\rangle + a_4|011\rangle + a_5|100\rangle + a_6|101\rangle + a_7|110\rangle + a_8|111\rangle$

If one can control this high-dimensional space, e.g. with appropriate interference of amplitudes, then one can potentially achieve **exponential speedup** of certain computations

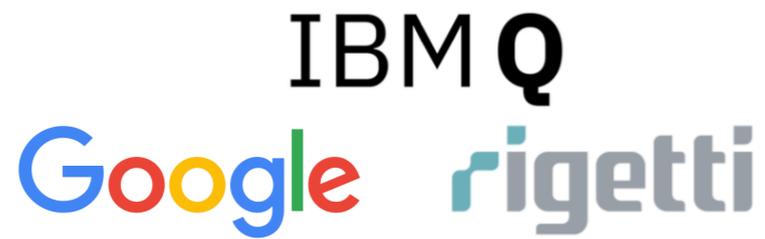


It is expected that quantum computers can solve *some* classically hard problems with exponential speedup

These include a number of highly impactful problems such as quantum simulation

Quantum devices

Superconducting circuits



...

And a variety of others...

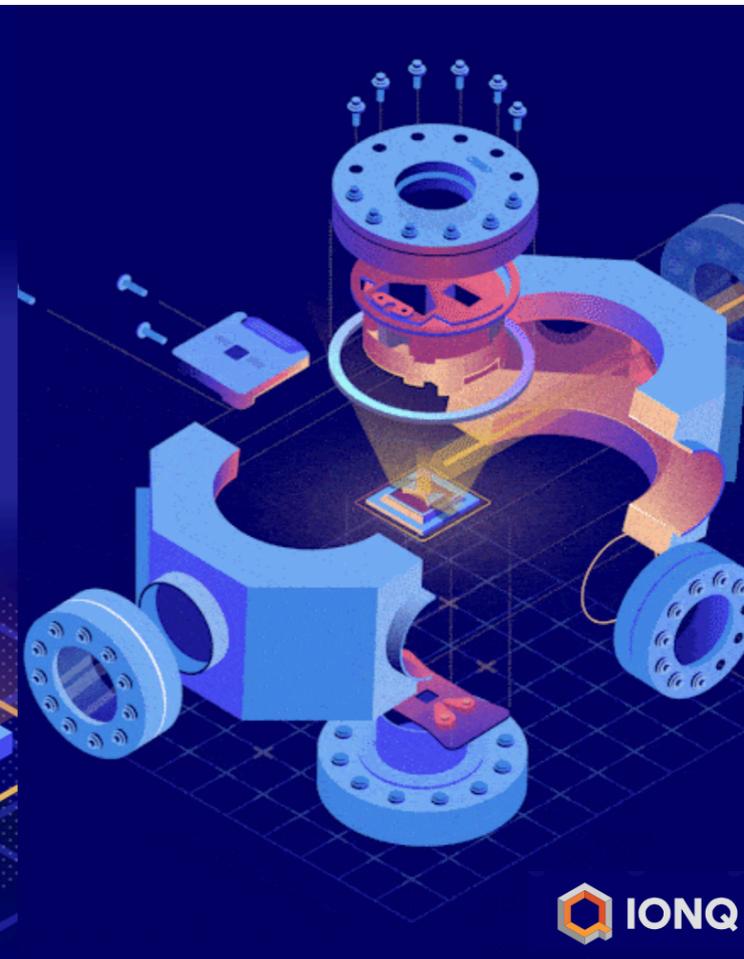
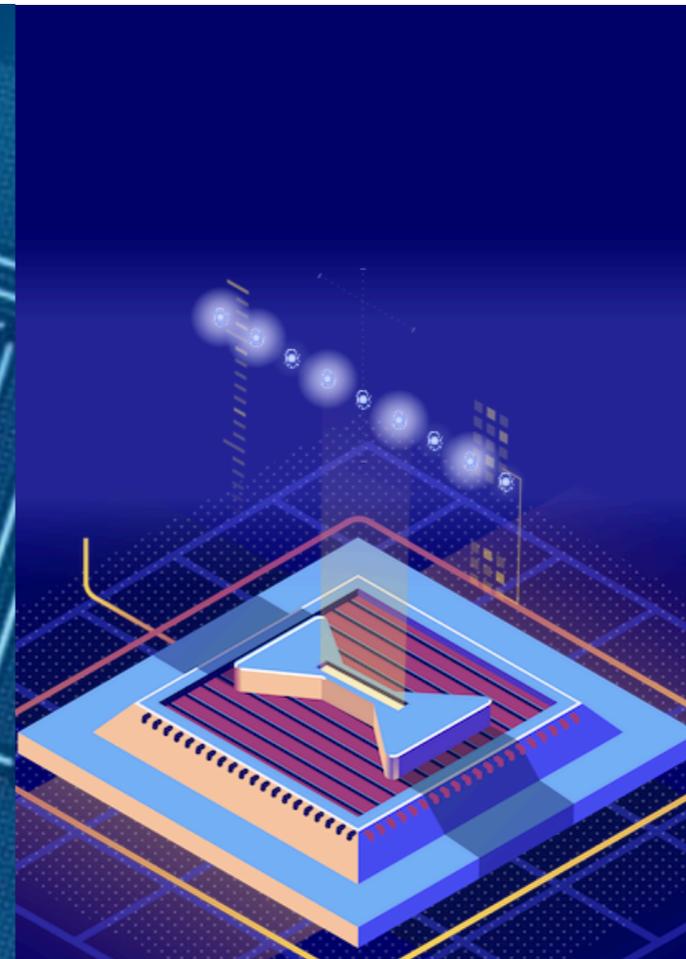
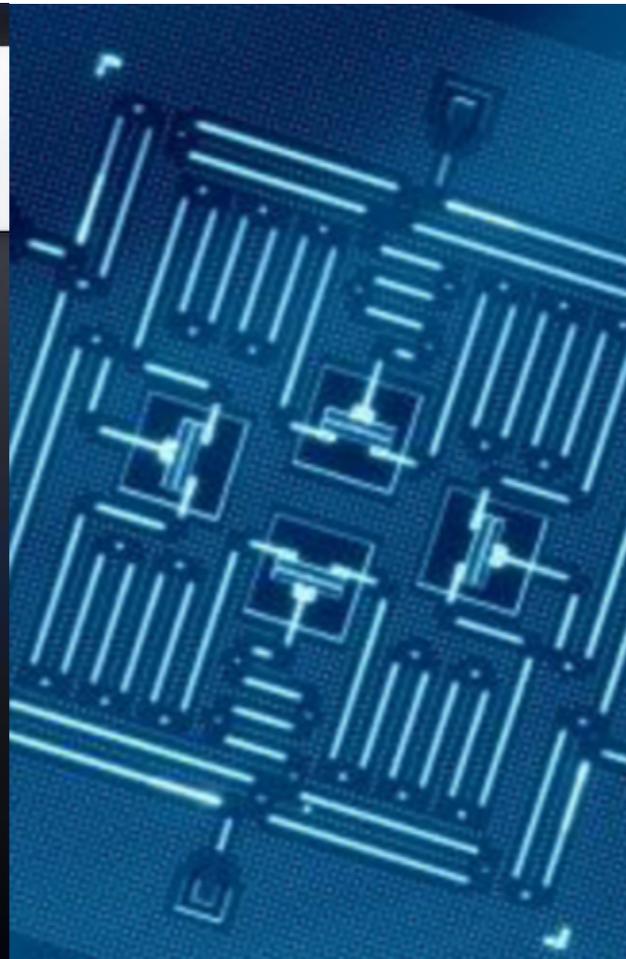
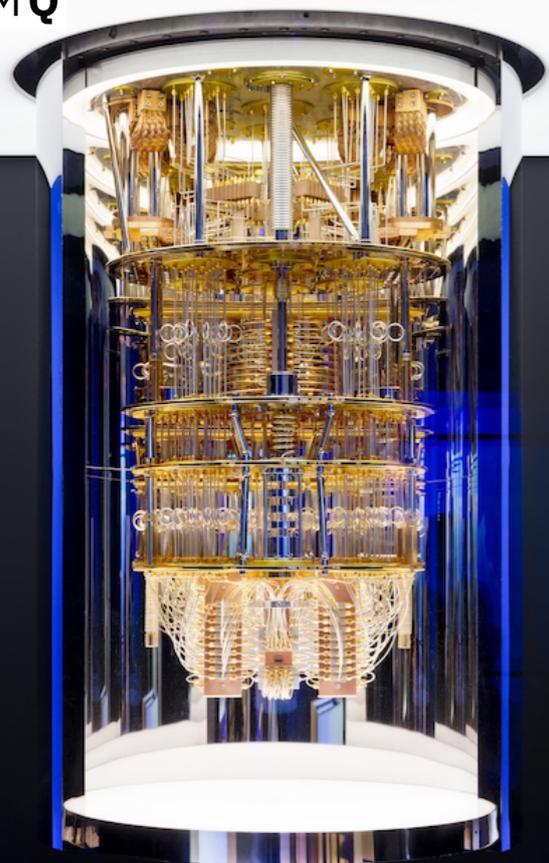
Trapped ions
 Optical lattice
 Photonics
 Topological

...



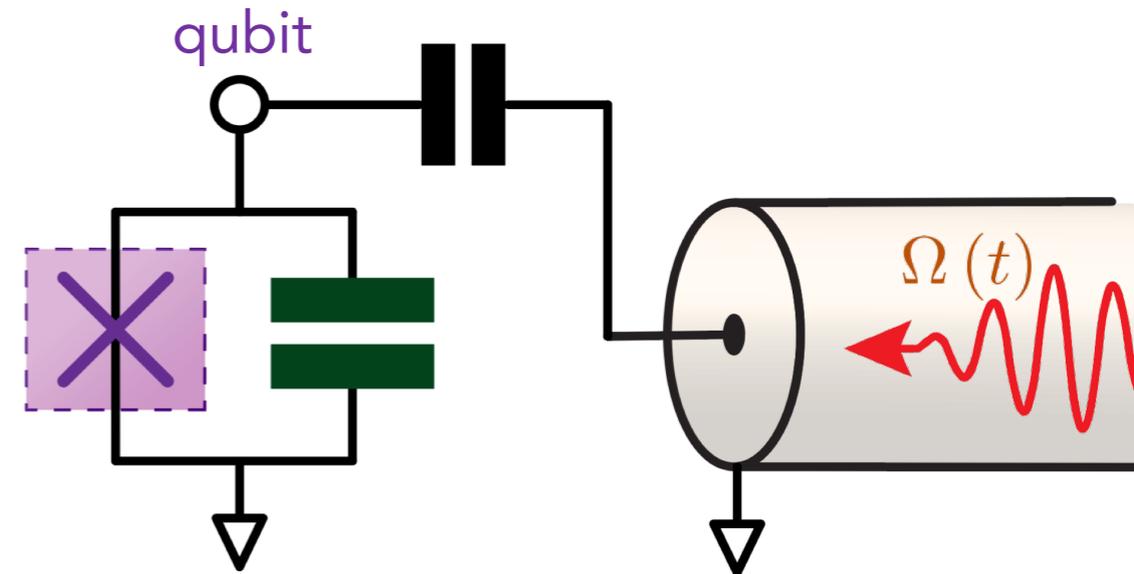
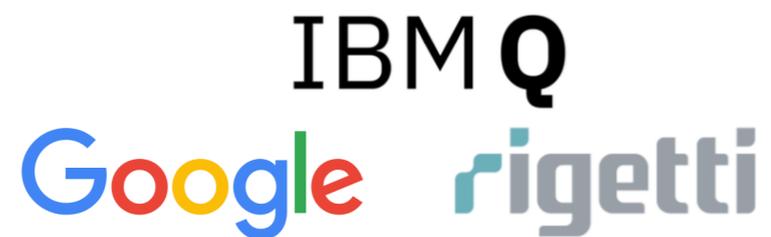
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IBM Q

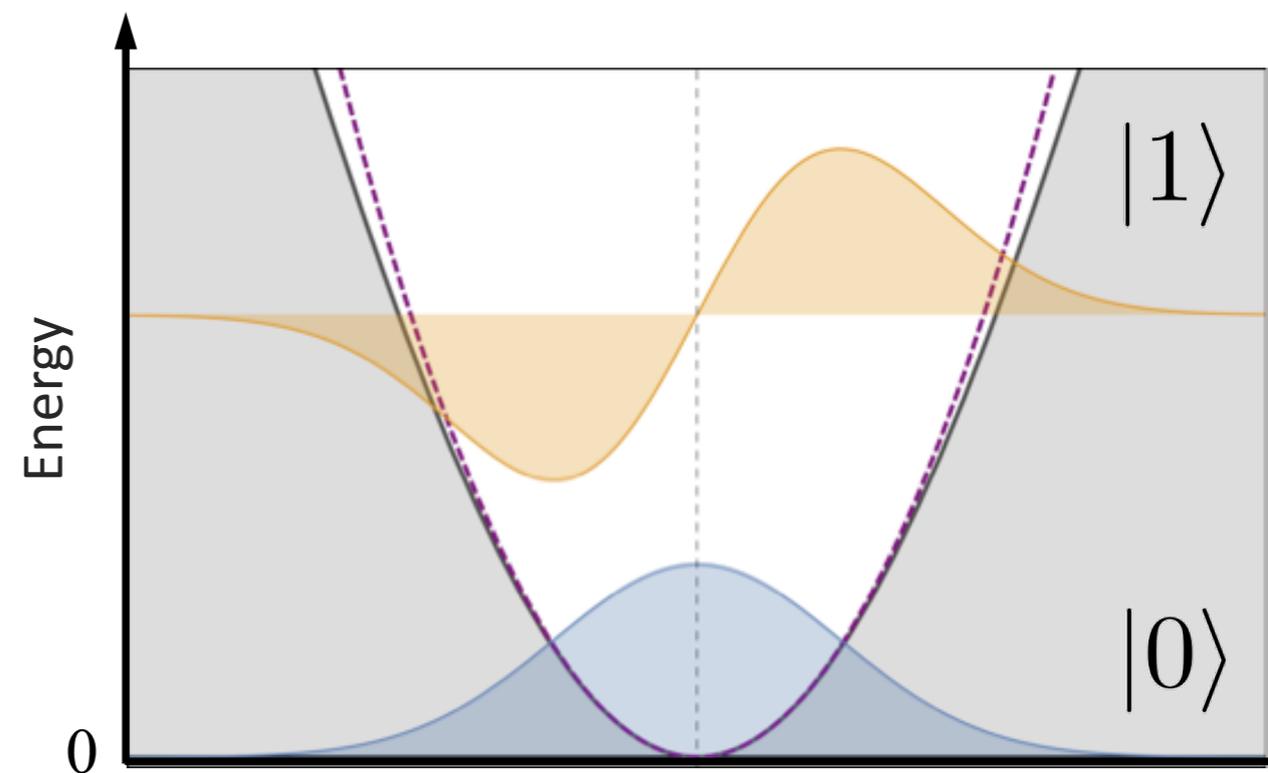


Quantum devices

Superconducting circuits

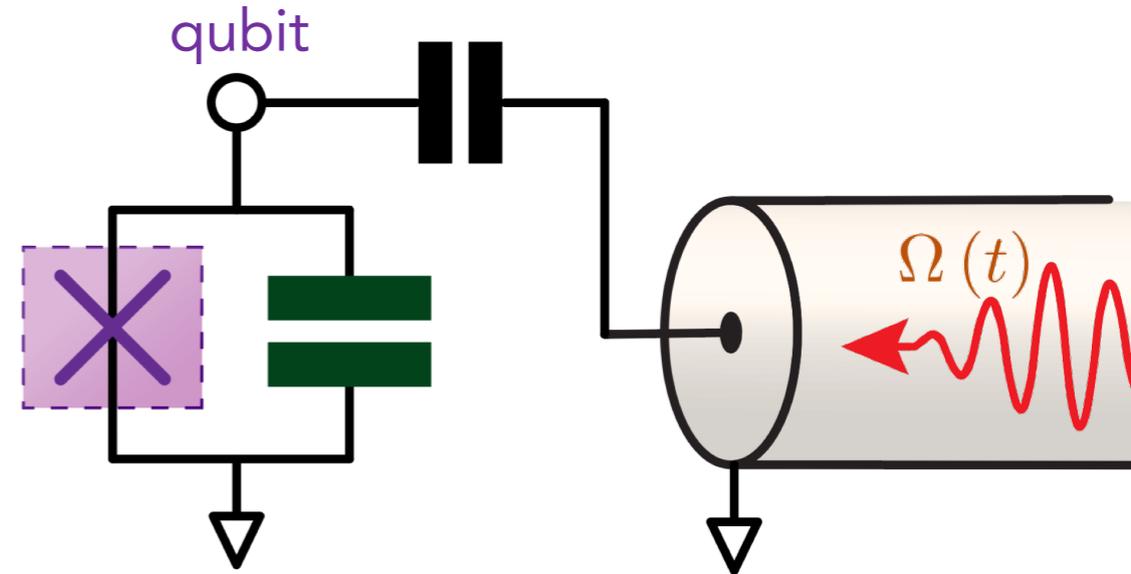
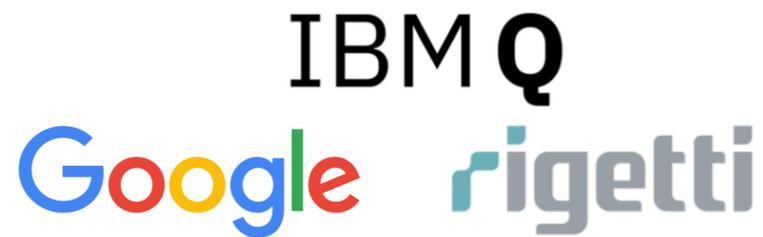


Qubits: Nonlinear quantum oscillator
Gates: coupled microwave pulses



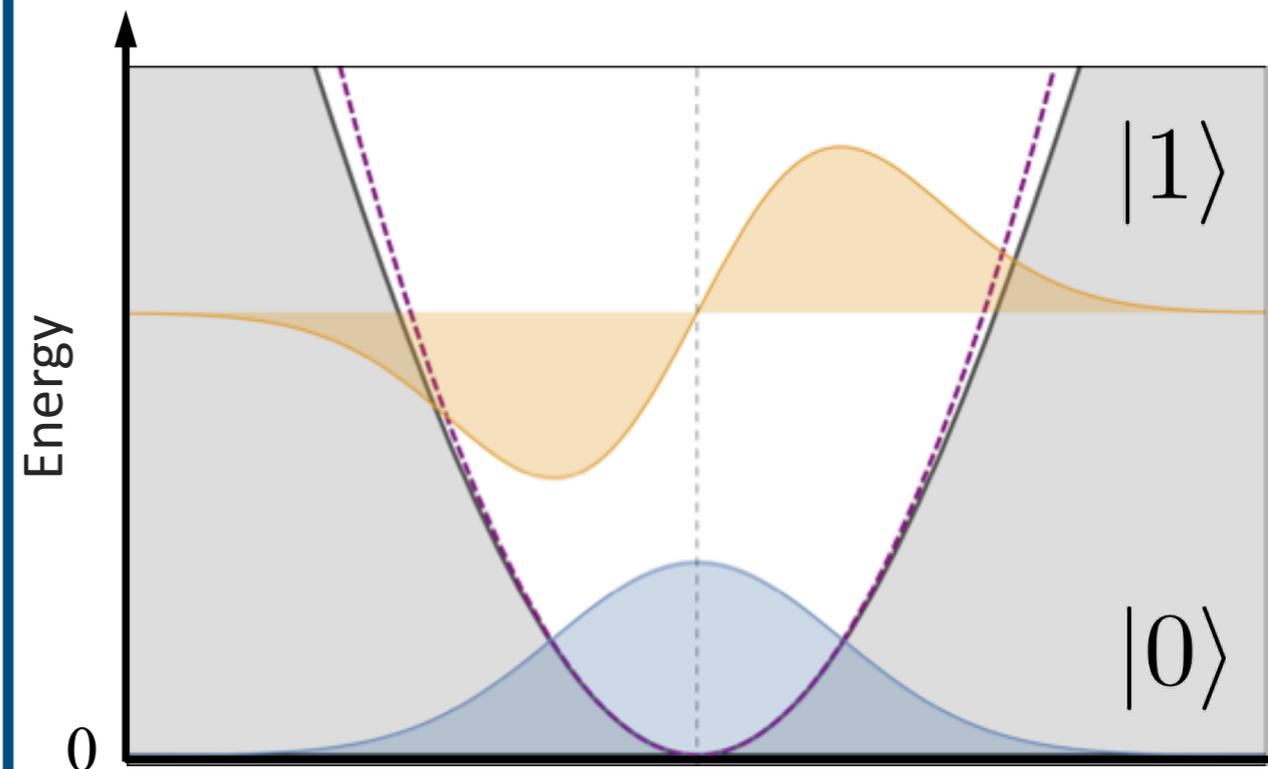
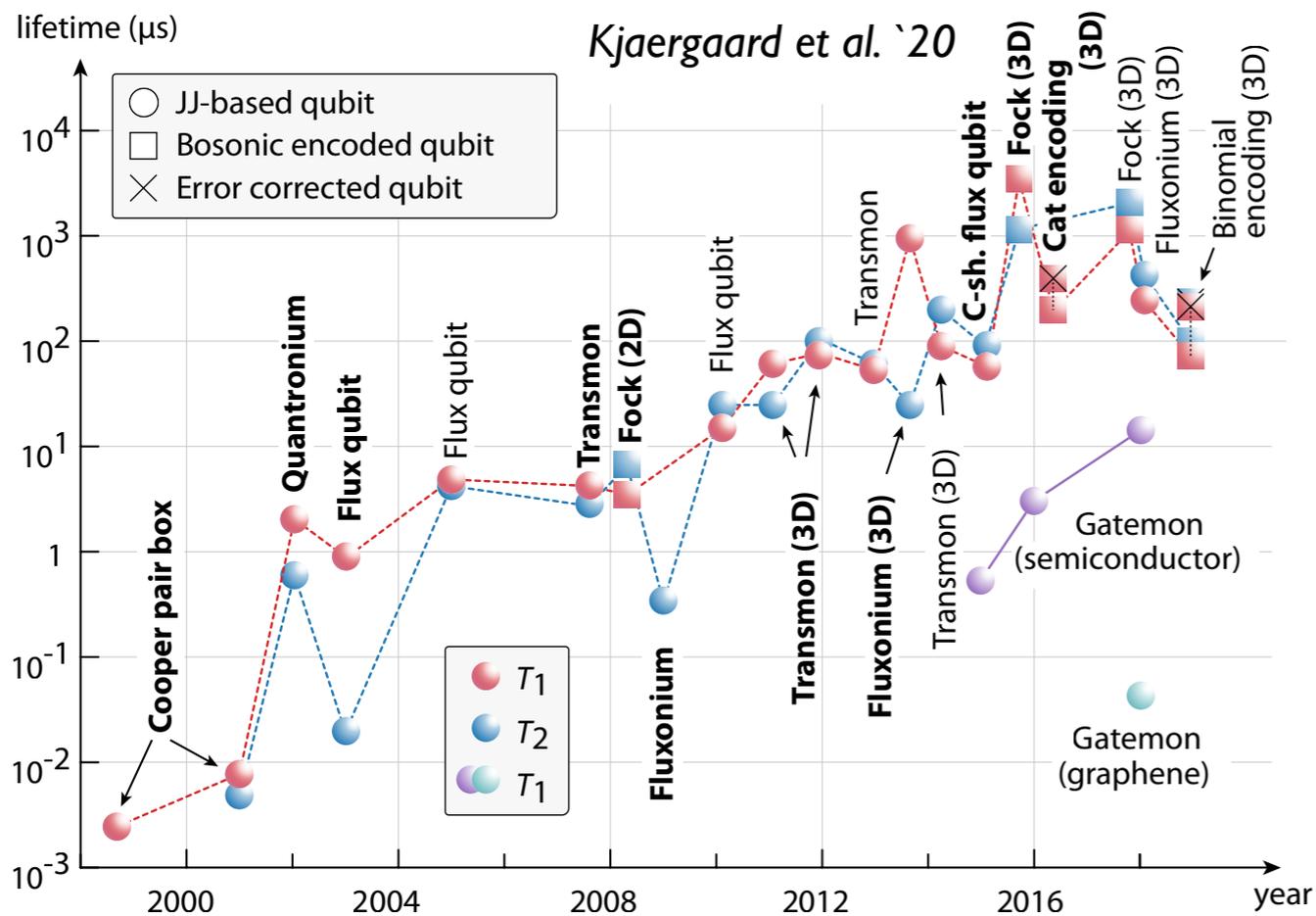
Quantum devices

Superconducting circuits



Qubits: Nonlinear quantum oscillator
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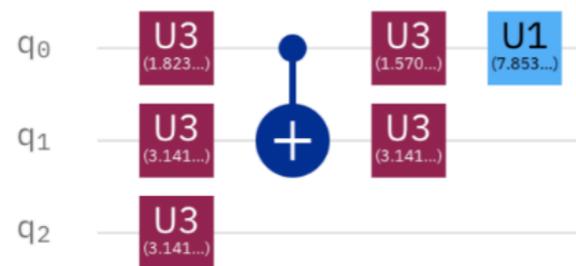
Qubit coherence times have become $\mathcal{O}(100\mu s)$, long enough to perform $\mathcal{O}(10 - 100)$ two-qubit operations



Quantum computing

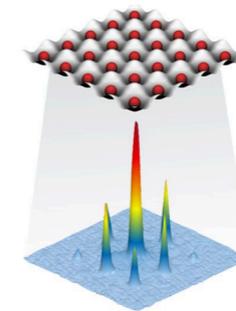
Digital quantum computers

Universal



Analog quantum computers

Application-specific



Both will likely be useful in the “near”-term

The dream: universal, fault-tolerant digital quantum computer

Shor, Preskill, Kitaev, Zoller ...

Noisy Intermediate Scale Quantum (NISQ) era

Decoherence, limited number of qubits, imperfect gates

Aim: achieve quantum advantage without full quantum error correction

Experimentation and data analysis



QUANTUM SYSTEMS ACCELERATOR

Catalyzing the Quantum Ecosystem

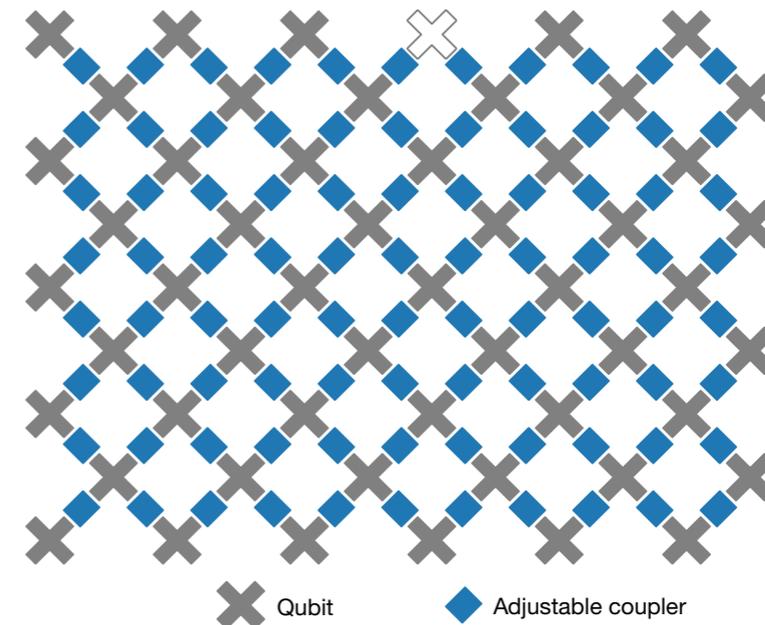
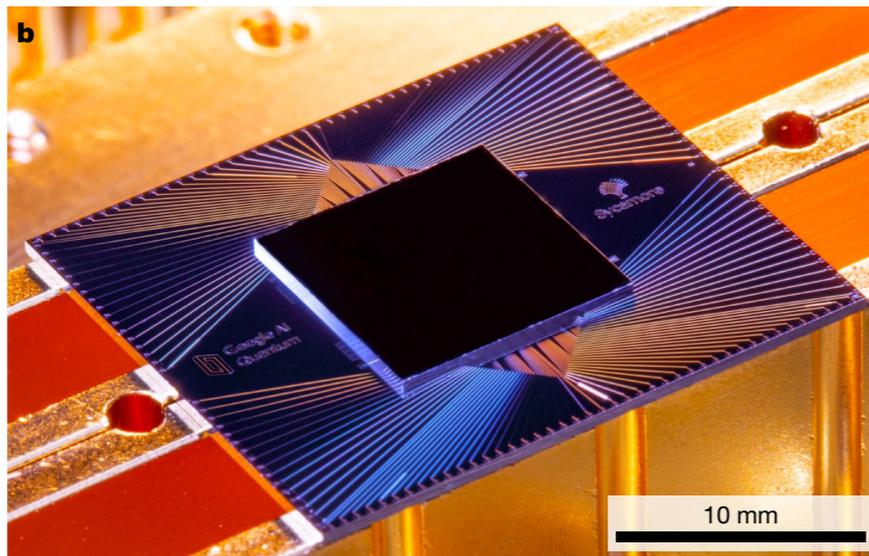
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Quantum supremacy

Google
Martinis et al. '19

Article

Quantum supremacy using a programmable superconducting processor



53-qubit sycamore device
99%+ gate fidelities

Algorithm: sampling of random circuits

$\mathcal{O}(10^3)$ **times faster than best classical supercomputers**

Quantum simulation

Feynman '81

It is exponentially expensive to simulate an N -body quantum system on a classical computer
 2^N amplitudes!

But a quantum computer can naturally simulate a quantum system

State preparation
Time evolution
Measurement



Holds great promise for particle physics

- Solve the **real-time dynamics** of QCD

Go beyond lattice QCD limitations (static quantities — sign problem)

see e.g. Jordan, Lee, Preskill '11, Preskill '18,
 Klco, Savage et al. '18-'20, Cloet, Dietrich et al. '19

Quantum simulation of open quantum systems

Toy model setup

Two-level system in a thermal environment

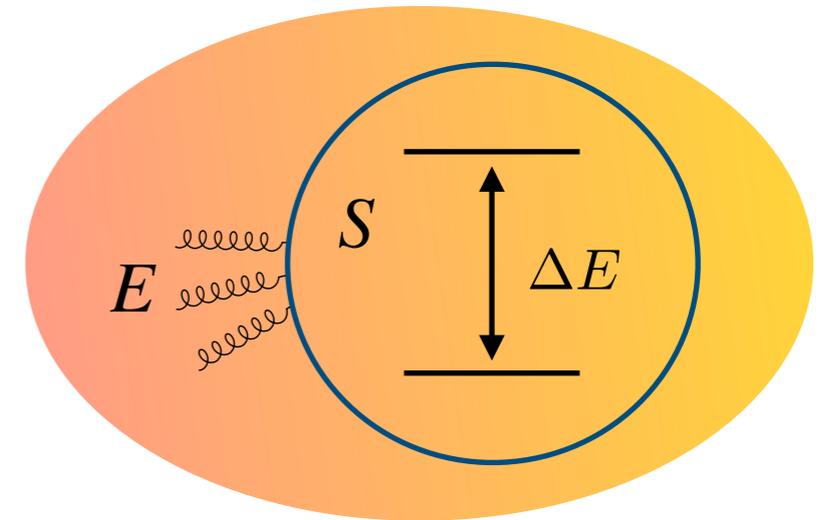
e.g. bound/unbound $J/\psi, c\bar{c}$

$$H_S = -\frac{\Delta E}{2} Z$$

$$H_E = \int d^3x \left[\frac{1}{2} \Pi^2 + \frac{1}{2} (\nabla \phi)^2 + \frac{1}{2} m^2 \phi^2 + \frac{1}{4!} \lambda \phi^4 \right]$$

$$H_I = g X \otimes \phi(x=0)$$

Pauli matrices X, Y, Z , interaction strength g



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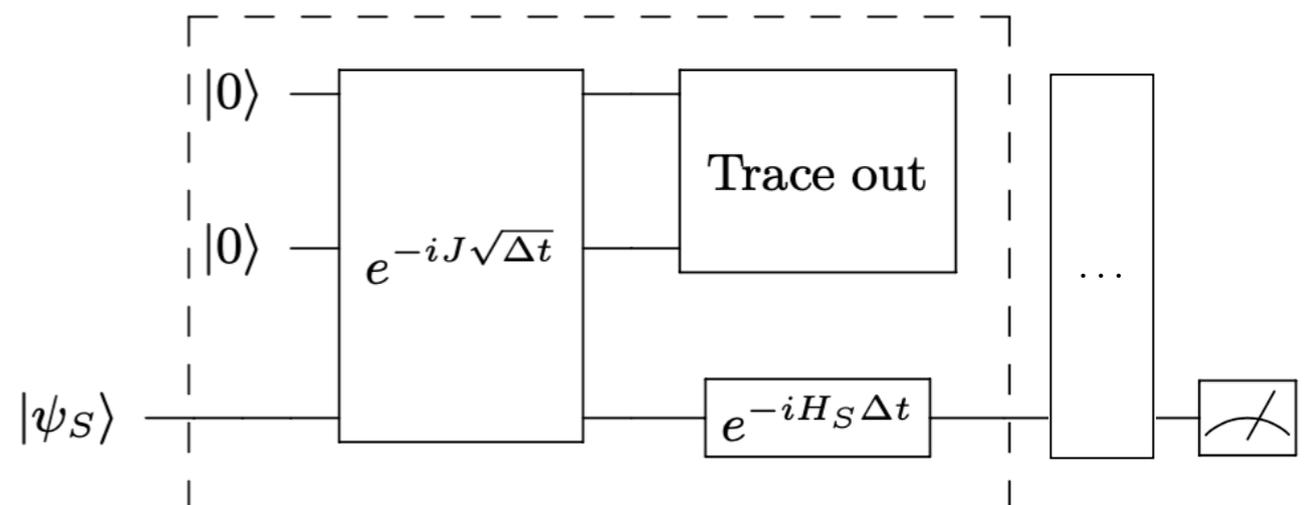
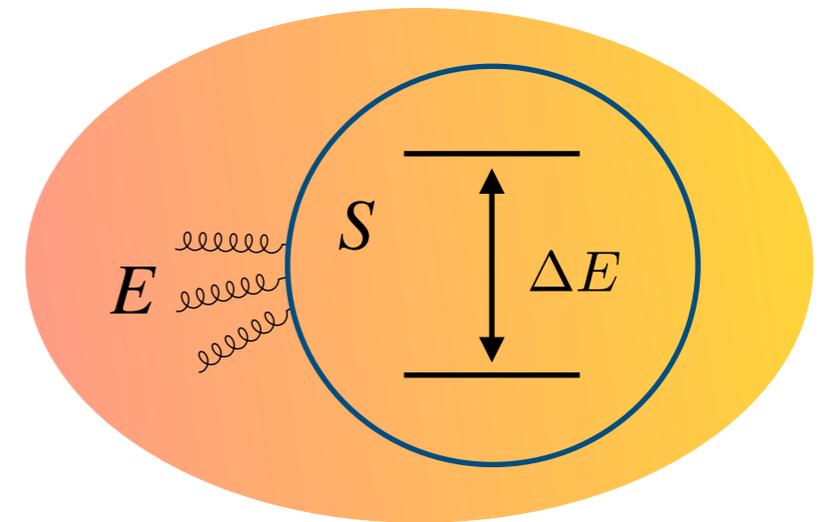
$$H_I = gX \otimes \phi(x=0)$$

Pauli matrices X, Y, Z , interaction strength g

Lindblad operators

$$L_j \sim g(X \mp iY) \quad j = 0, 1$$

$$J = \begin{pmatrix} 0 & L_0^\dagger & L_1^\dagger & 0 \\ L_0 & 0 & 0 & 0 \\ L_1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

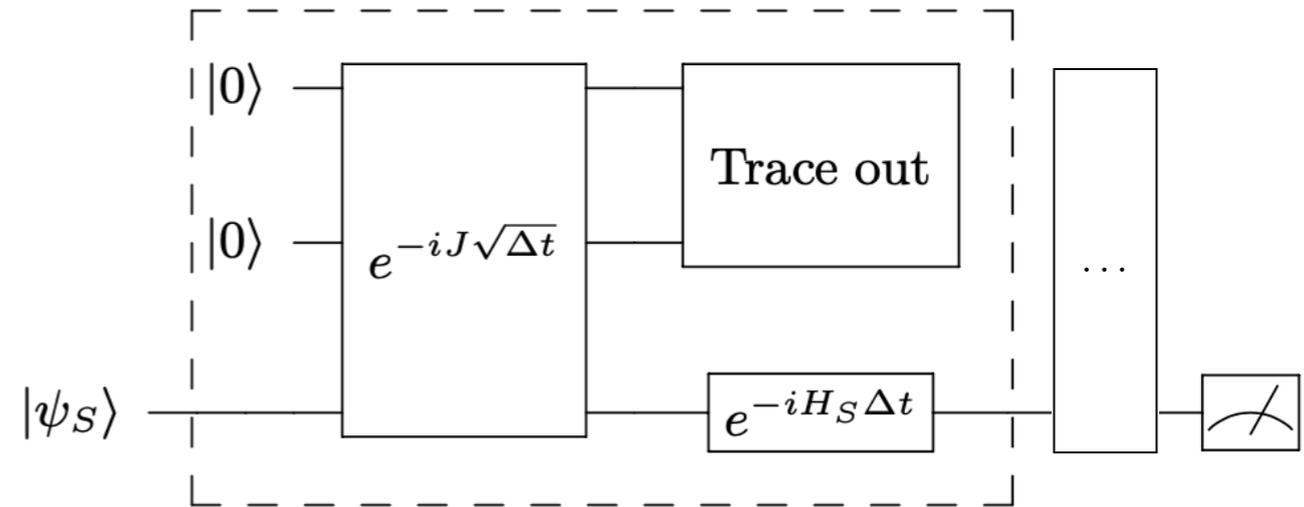


Quantum circuit synthesis

Approximate unitary operations with a compiled circuit of one- and two-qubit gates

Optimization problem w/unitary loss function

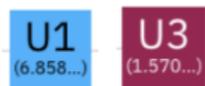
qsearch Siddiqi et al. '20 @LBNL



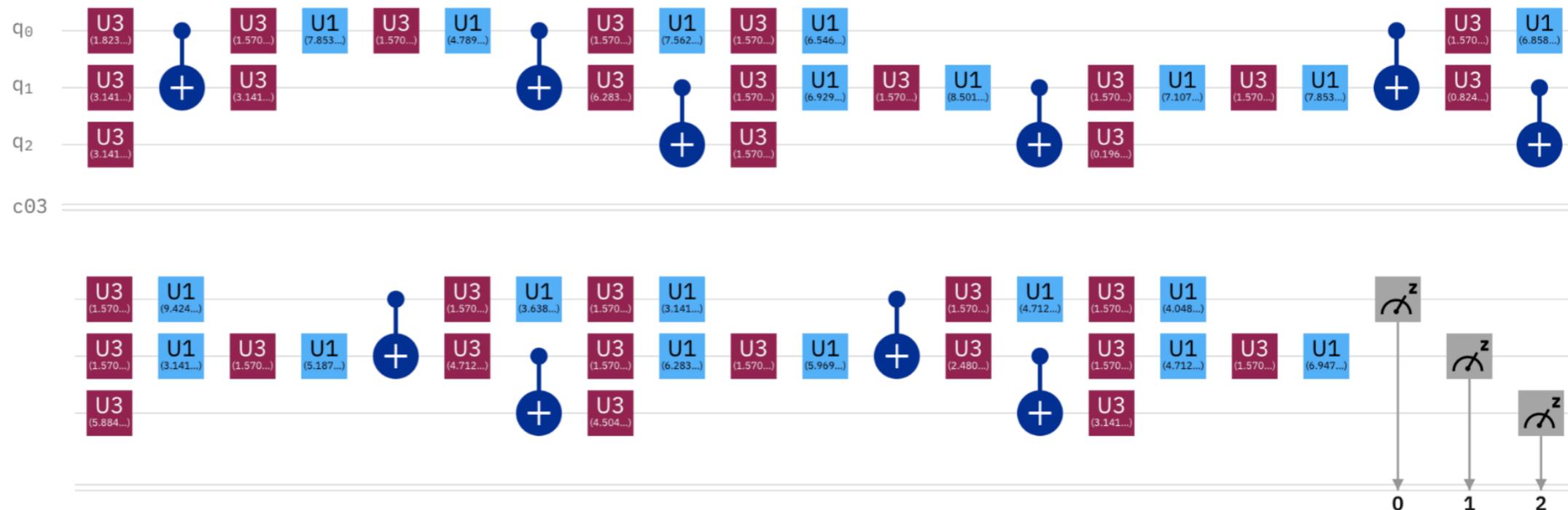
10 CNOT gates/cycle

IBM Q

Single qubit



CNOT



Error mitigation

Readout error

Constrained matrix inversion

IBM Q qiskit-ignis

Unfolding

Nachman, Urbanek, de Jong, Bauer '19

@LBNL

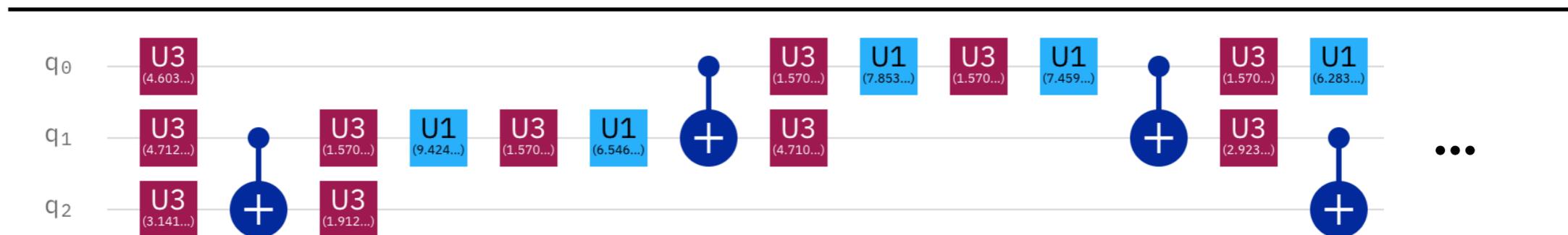
Gate error

Zero-noise extrapolation of CNOT noise using Random Identity Insertions

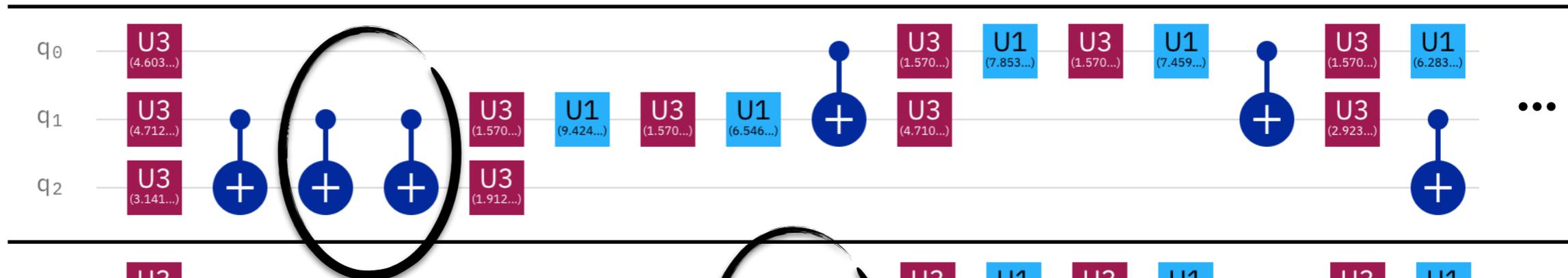
He, Nachman, de Jong, Bauer '20

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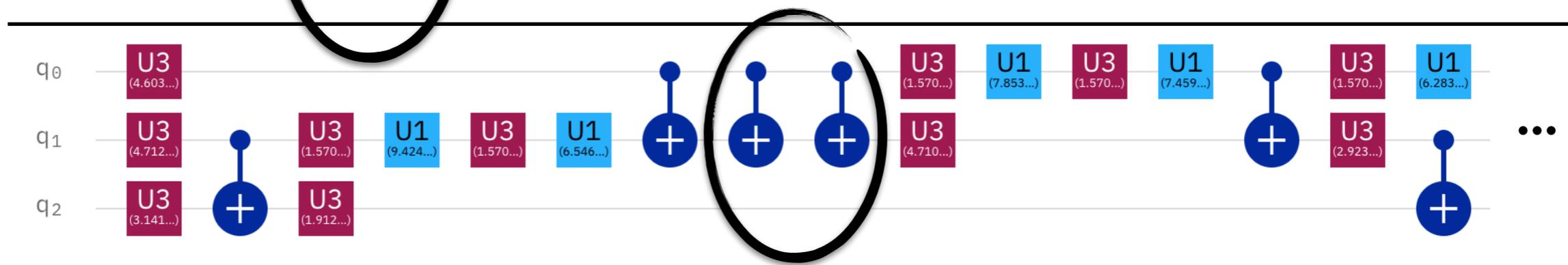
Circuit 1



Circuit 2



Circuit 3



...

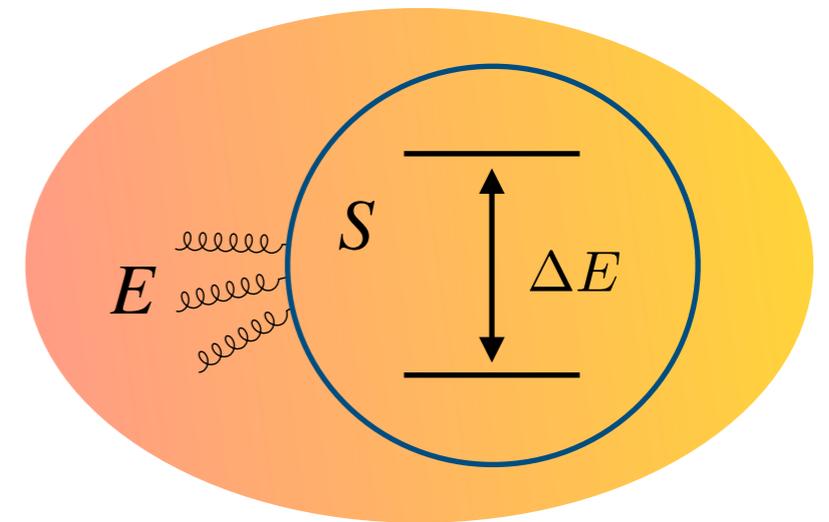
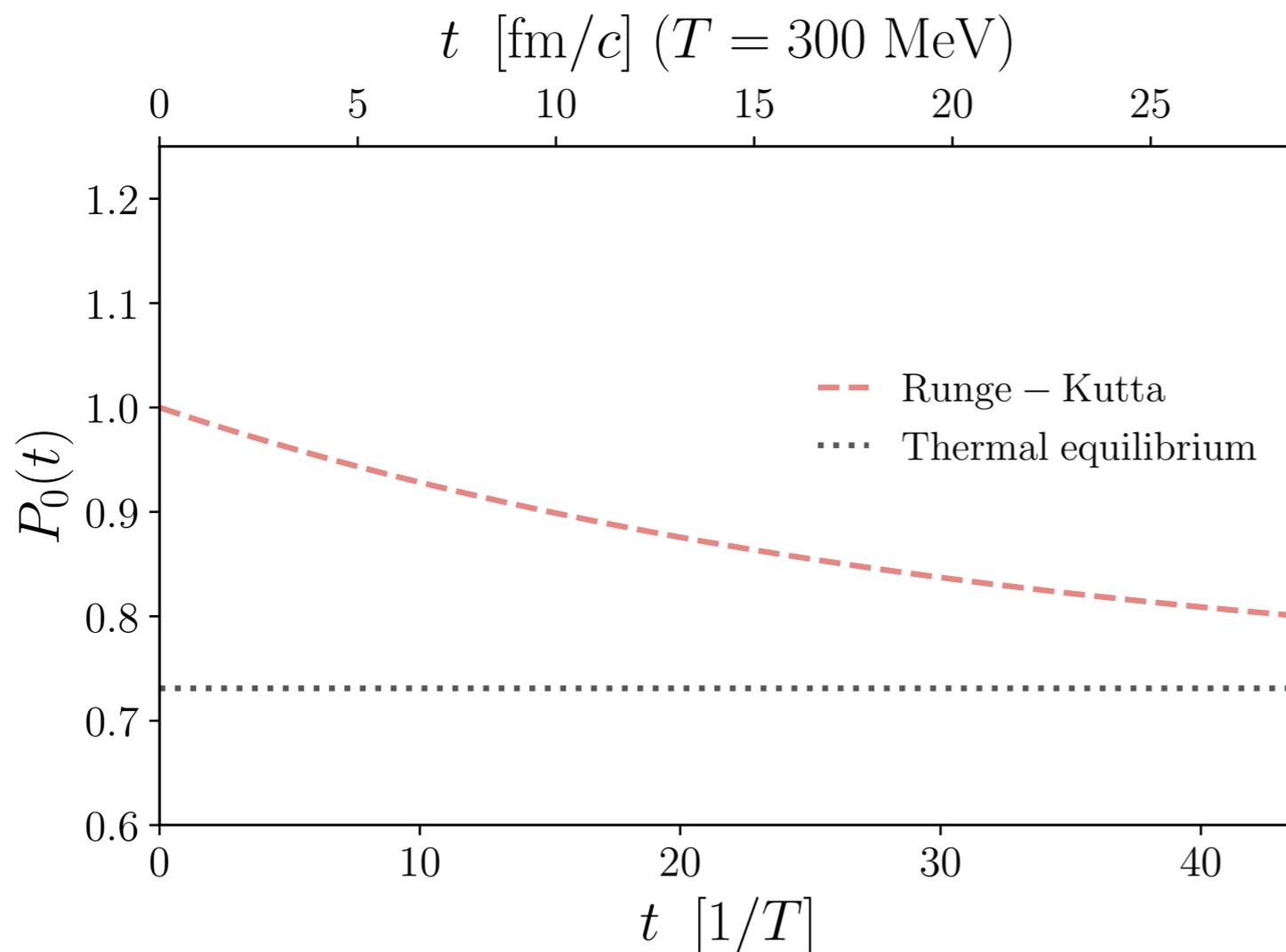
Quantum simulation of open quantum systems

arXiv: 2010.03571

Real-time evolution

$P_0(t)$ describes fraction that remains in “bound state”

$$\text{Similar to } t\text{-dependent } R_{AA} = \frac{d\sigma_{AA}}{\langle N_{\text{coll}} \rangle d\sigma_{pp}}$$



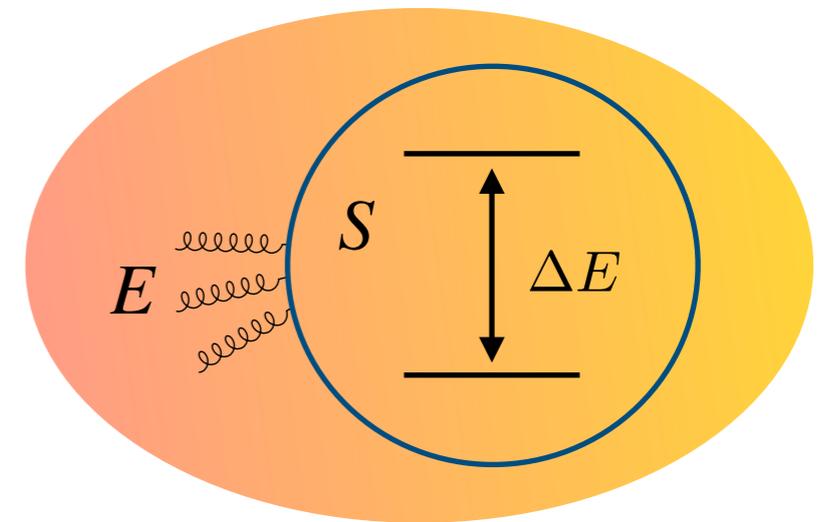
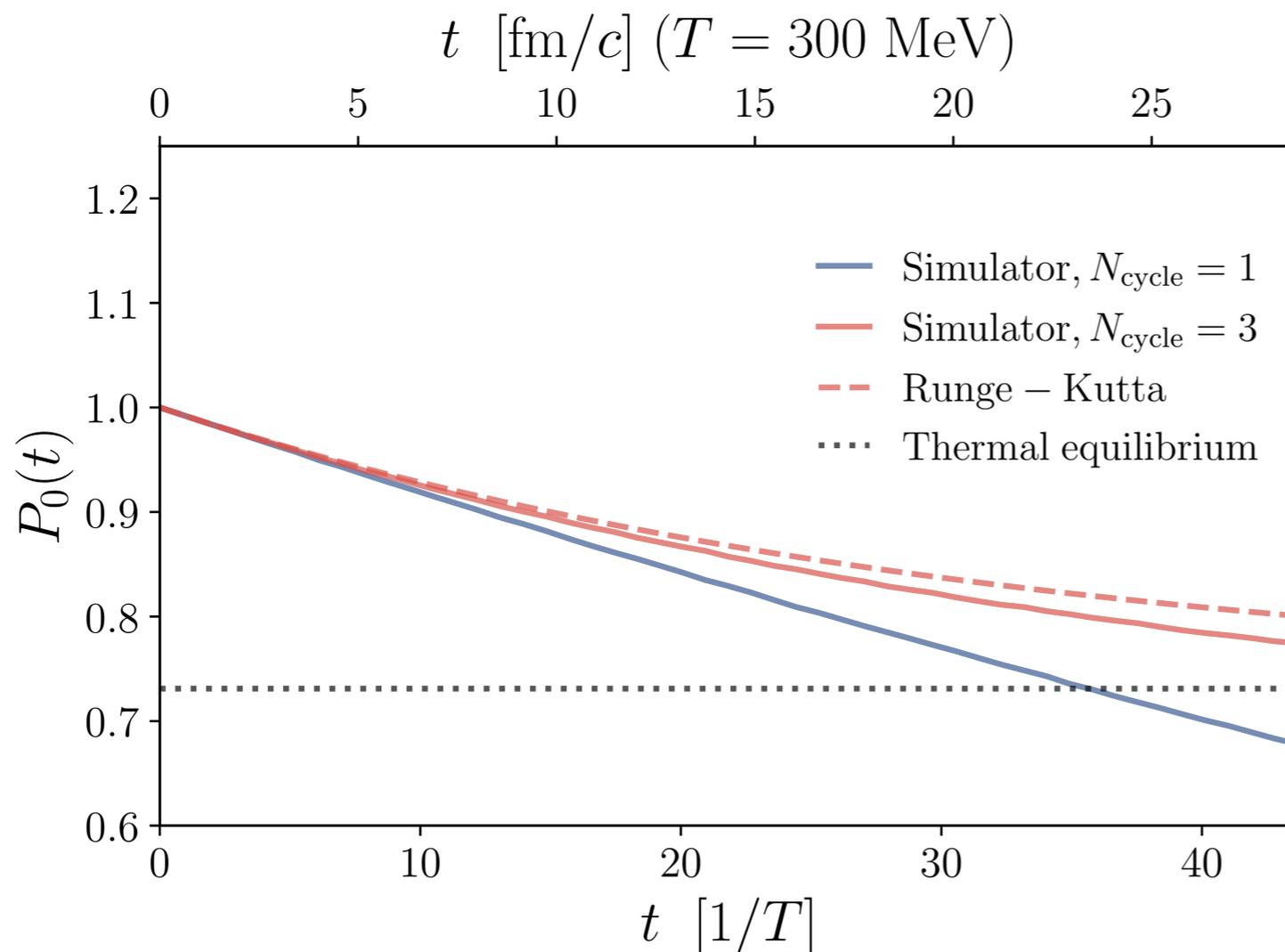
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The algorithm converges to Lindblad evolution with a small number of cycles

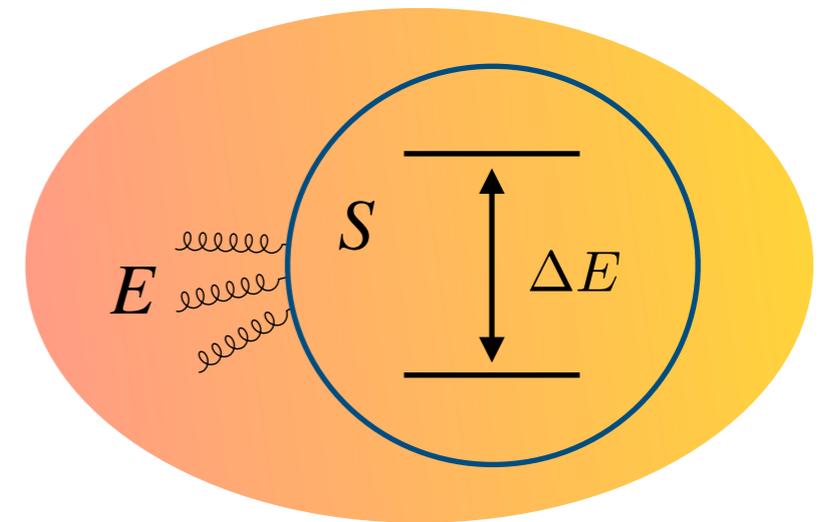
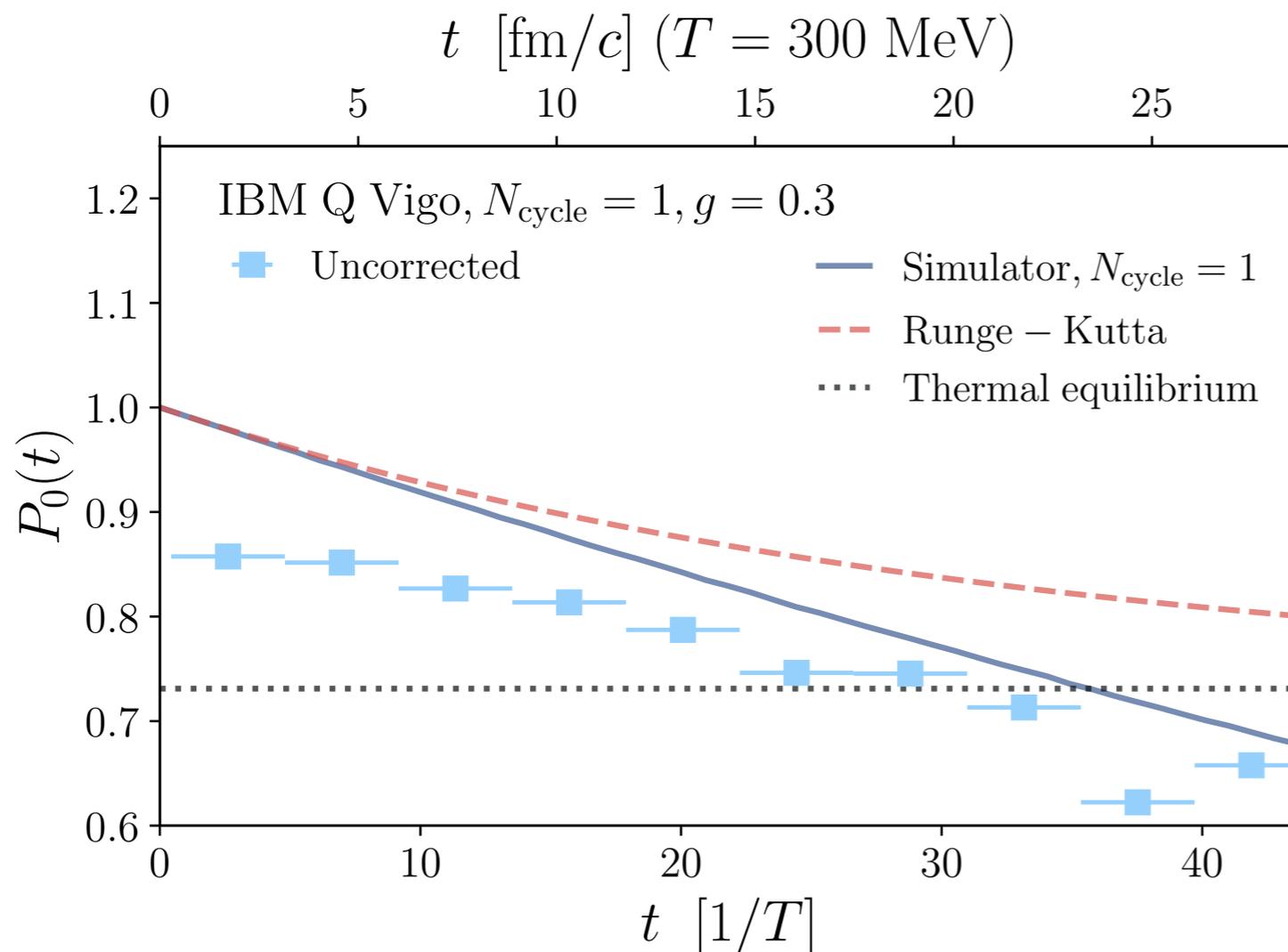
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IBM Q Vigo device

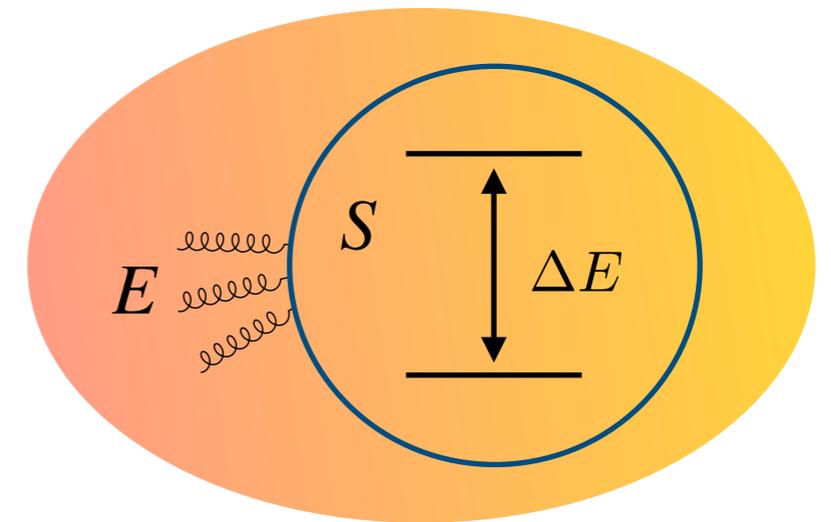
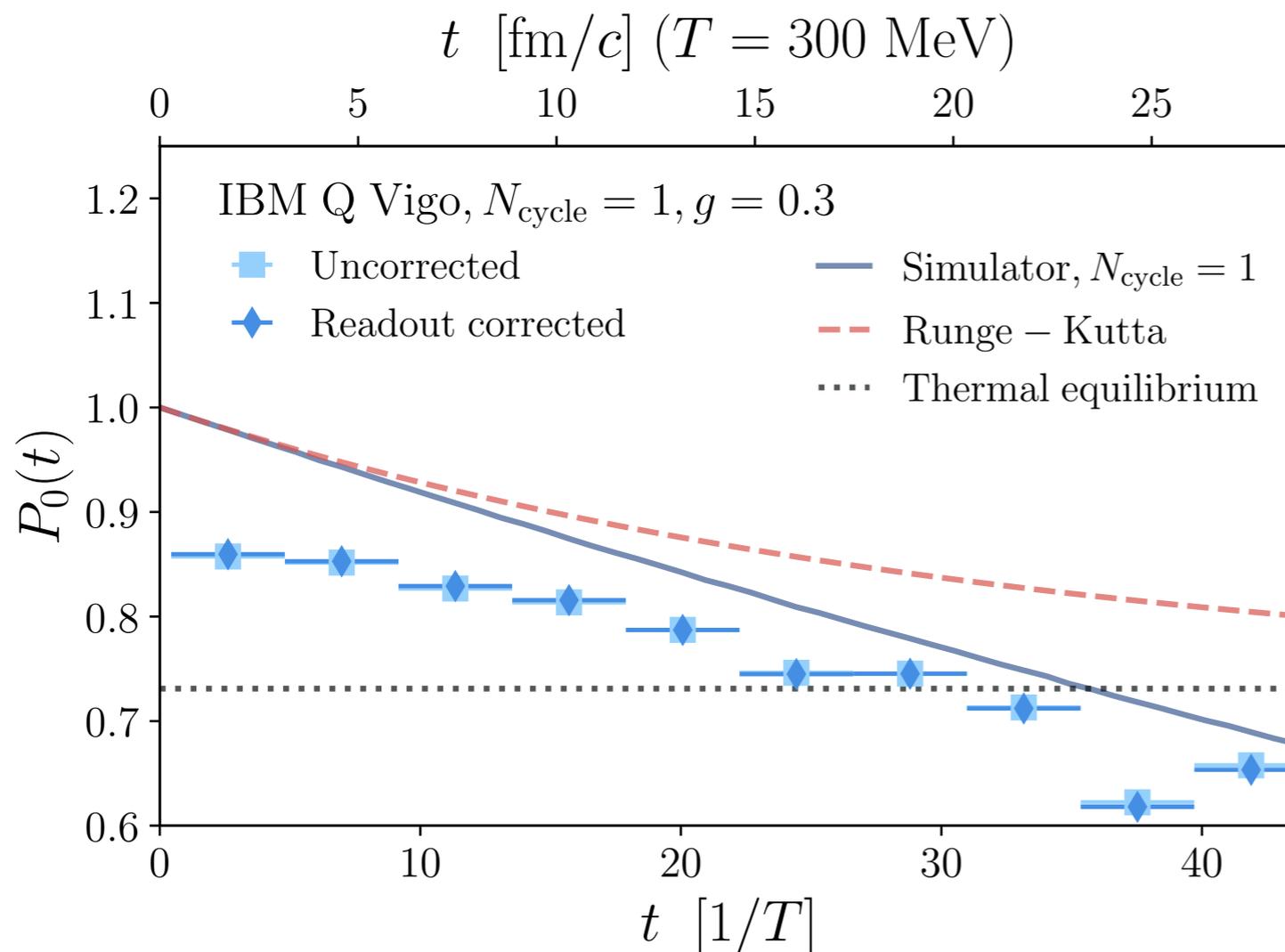
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IBM Q Vigo device

Readout correction small

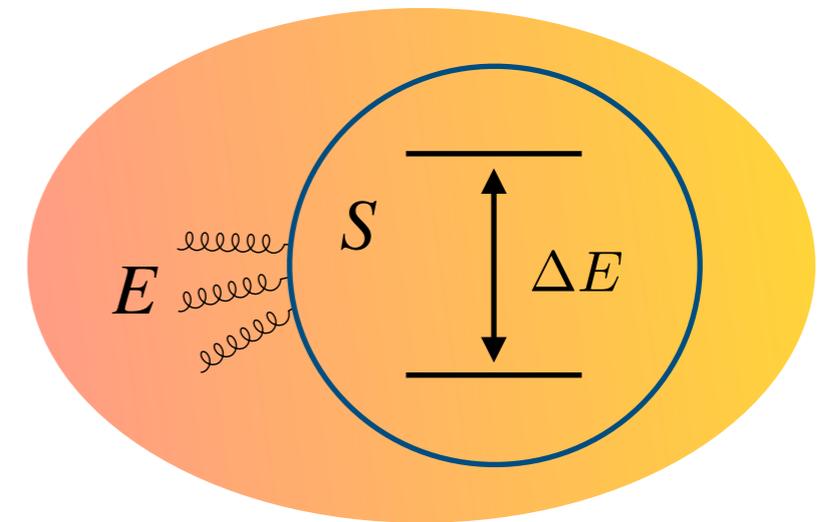
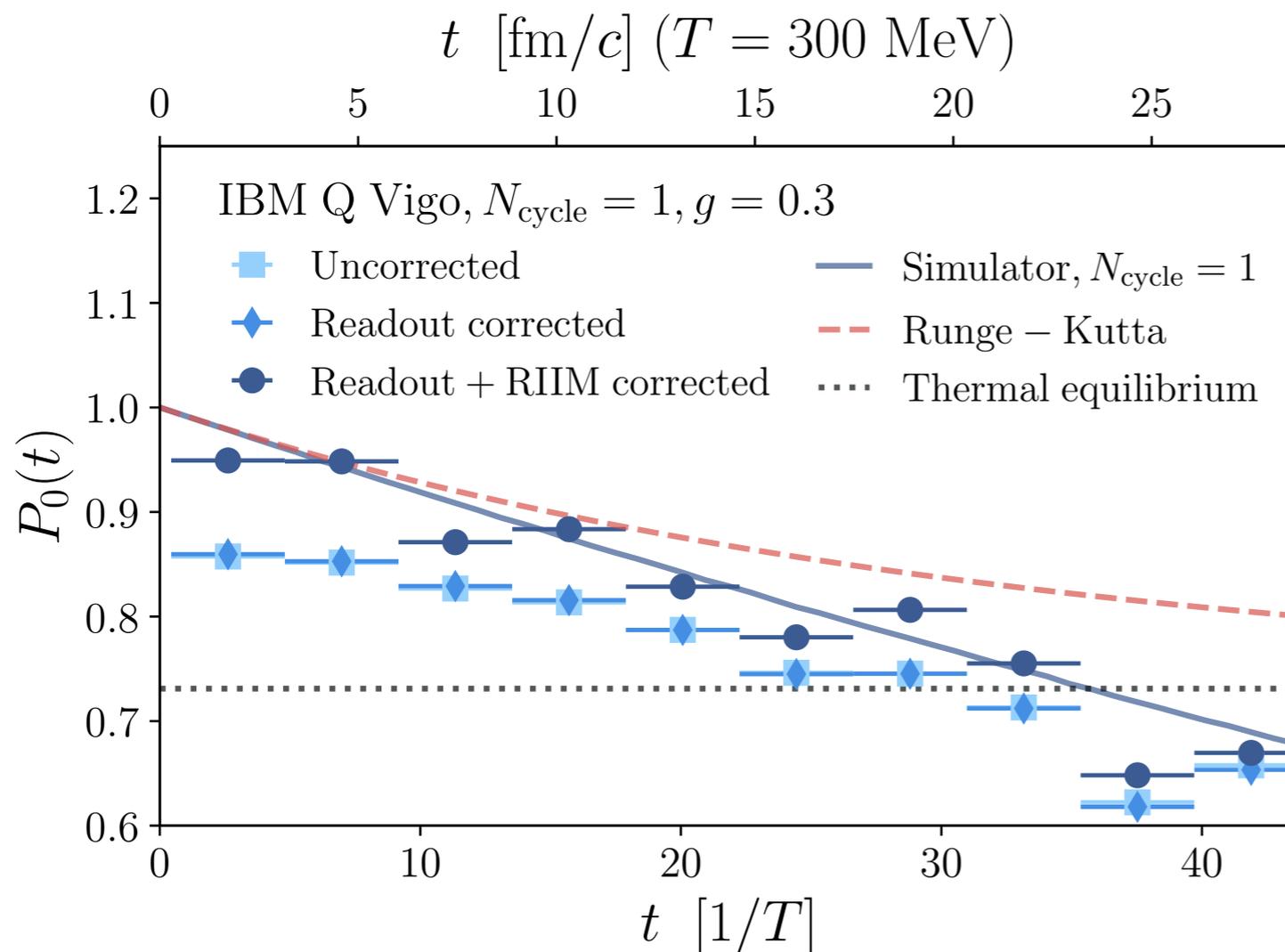
Quantum simulation of open quantum systems

arXiv: 2010.03571

Real-time evolution

$P_0(t)$ describes fraction that remains in “bound state”

$$\text{Similar to } t\text{-dependent } R_{AA} = \frac{d\sigma_{AA}}{\langle N_{\text{coll}} \rangle d\sigma_{pp}}$$



IBM Q Vigo device

Readout correction small

CNOT gate error correction
gives good agreement

Random Identity Insertion Method (RIIM)

Bauer, He, de Jong, Nachman `20

Proof of concept

Outline

Open quantum
systems in heavy-ion
collisions

Quantum simulation
with IBM Q

Conclusions and outlook

- **Open quantum system formalism describes the real-time evolution of hard probes in heavy-ion collisions**
 - Allows to go beyond semiclassical approximations in current models
- **Proof of concept that these systems can be simulated on current and near-term quantum computers (IBM Q)**
 - NISQ era digital quantum computing
 - Recently developed error mitigation techniques
- **Future steps**
 - Extension toward QCD (jets & heavy-flavor)
 - Explore different digital/analog devices
 - Cold nuclear matter at the EIC